

MECHANICAL SCIENCE-II



B. Bhattacharyya S.C. Bera





MECHANICAL SCIENCE-II

Dr. Basudeb Bhattacharyya

Department of Applied Mechanics Bengal Engineering and Science University, Shibpur, Howrah – 711103

Prof. Subal Chandra Bera

Narula Institute of Technology Agarpara, Kolkata – 700109

Previously Department of Applied Mechanics Bengal Engineering and Science University, Shibpur Howrah – 711103



NEW AGE INTERNATIONAL (P) LIMITED, PUBLISHERS

New Delhi • Bangalore • Chennai • Cochin • Guwahati • Hyderabad Jalandhar • Kolkata • Lucknow • Mumbai • Ranchi Visit us at www.newagepublishers.com Copyright © 2009, New Age International (P) Ltd., Publishers Published by New Age International (P) Ltd., Publishers

All rights reserved.

No part of this ebook may be reproduced in any form, by photostat, microfilm, xerography, or any other means, or incorporated into any information retrieval system, electronic or mechanical, without the written permission of the publisher. *All inquiries should be emailed to rights@newagepublishers.com*

ISBN (13): 978-81-224-2871-1

PUBLISHING FOR ONE WORLD

NEW AGE INTERNATIONAL (P) LIMITED, PUBLISHERS 4835/24, Ansari Road, Daryaganj, New Delhi - 110002 Visit us at **www.newagepublishers.com** Dedicated to the student community who will use this book

PREFACE

Mechanical Sciences-II, within the purview of WBUT syllabus, deals the fundamental aspects of thermodynamics and fluid mechanics. Both of thermodynamics and fluid mechanics have separate identity as full-fledged subjects. Hence to study both of them, the students of WBUT may need to separate books. Here in a single book these two subjects' matters are complied for the good of students. To achieve an in-depth knowledge of these two basic subjects, it needs very lucid elaboration of the text and also presentation of numerous worked out examples for complete understanding of the topic. Keeping these basic requirements in mind, we have tried to exert our best to come up with a most useful book.

In this newly revised edition, we have tried to rectify the errors and omissions detected earlier. The present version mainly features

- * Thorough discussion of the topic in basic level, avoiding much detailing and critical discourse.
- Inclusion of multiple choice questions relevant to the topic.
- Provision of number of useful numerical examples for holistic understanding of the topic.

The forbearance as well as the active support of our family members and the useful suggestions by our colleagues and the commendable support from the editorial team of the publishing house consorted successfully to bloom this project. Though proper care has been adopted to make the book error free, any unintended mistake or error or omission noticed may be directly communicated to the publisher or the author(s).

> Dr. Basudeb Bhattacharya Prof. Subal Chandra Bera

CONTENTS

	Preface Syllabus	vii ix
	Group–A: Thermodynamics	
Сна	PTER 1: BASIC CONCEPT AND SOME DEFINITIONS	3–31
1.1	Introduction	3
1.2	Macroscopic and Microscopic Approach	3
1.3	Thermodynamic System	4
1.4	Classification of Thermodynamic Systems	4
1.5	Control Mass and Control Volume	6
1.6	Thermodynamic Co-ordinates	7
1.7	State of a System	7
1.8	Properties of a System	7
1.9	Classification of Properties of a System	7
1.10	Phase	8
1.11	Thermodynamic Equilibrium	8
1.12	Path	9
1.13	Process	9
1.14	Cyclic Process or Thermodynamic Cycle	10
1.15	Point Function and Path Function	12
1.16	Unit	13
1.17	Systems of Units	15
1.18	Mass (M)	15
1.19	Weight (W)	16
1.20	Force (F)	16
1.21	Specific Weight (W_s)	17
1.22	Specific volume (v_s)	17

1.23	Pressure	1
1.24	Absolute, Gauge and Vacuum Pressure	1'
1.25	Pressure Measurement by Manometer	1
1.26	Normal Temperature and Pressure (N.T.P)	2
1.27	Standard Temperature and Pressure (S.T.P)	2
1.28	Energy	2
1.29	Types of Stored Energy	2
1.30	Law of Conservation of Energy	2
1.31	Power	2
Снаі	PTER 2: ZEROTH LAW AND TEMPERATURE	33–57
2.1	Temperature	3.
2.2	Zeroth Law of Thermodynamics	3
2.3	Measurement of Temperature	3
2.4	Constant Volume Gas Thermometer and Temperature Scale	3
2.5	Heat and Heat Transfer	3
2.6	Specific Heat	3
2.7	Thermal or Heat Capacity of a Substance	3
2.8	Water Equivalent of a Substance	3
2.9	Mechanical Equivalent of Heat	3
2.10	Work	3
2.11	Sign Convention of Work	4
2.12	Work Done During a Quasi-Static or Quasi-Equilibrium Process	4
2.13	Work and Heat Transfer—A Path Function	4
2.14	Comparison of Heat and Work	4
2.15	Example of Work	4
2.16	Work in Non-flow Process Versus Flow Process	5
Снаі	PTER 3: PROPERTIES AND THERMODYNAMIC PROCESSES OF GAS	59–102
3.1	Introduction	5
3.2	General Gas Equation	5
3.3	Equation of State and Characteristic Equation of Gas	6
3.4	Universal Gas Constant or Molar Gas Constant	6
3.5	Specific Heat of Gas	6
3.6	Linear Relation Between C_p and C_v	6
3.7	Ratio of C_p and C_v	6
3.8	Enthalpy of a Gas	6
3.9	Thermodynamic Processes	6
3.10	Classification of Thermodynamic Process	6
311	Heating and Expansion of Gases in Non-flow Process	6
5.11	Constant Volume or Isochoric Process	6
3.12		
3.12 3.13	Constant Pressure or Isobaric Process	6

Conte	nts	(xiii)
3.16	Adiabatic Process	71
3.17	Polytropic Process	74
3.18	General Laws for Expansion and Compression	75
3.19	Real Gas	76
3.20	Real Gas and Compressibility Factor	76
3.21	Law of Corresponding State and Generalized Compressibility Chart	77
Сна	PTER 4: FIRST LAW OF THERMODYNAMICS AND ITS APPLICATION	103–131
4.1	Introduction	103
4.2	Joule's Experiment	103
4.3	First Law of Thermodynamics for a System Undergoing a Thermodynamics Cycle	104
4.4	The Important Consequences of the First Law of Thermodynamics	105
4.5	Limitations of First Law of Thermodynamics	108
4.6	Application of First Law Thermodynamics to a Non-flow Process	109
4.7	Application of First Law of Thermodynamic to a Steady Flow Process	111
4.8	Mass Balance (Continuity Equation)	113
4.9	Work Done in a Steady Flow Process	113
4.10	Work Done Various Steady Flow Process	114
4.11	Throttling Process	115
4.12	Application of Steady Flow Energy Equation to Engineering System	115
Сна	PTER 5: SECOND LAW OF THERMODYNAMICS	133–155
5.1	Introduction	133
5.2	The Second Law of Thermodynamics	135
5.3	Equivalence of Kelvin-Planck and Clausius of Second Law of Thermodynamics	137
5.4	Reversible Cycle	138
5.5	Irreversible Cycle	138
5.6	Reversibility and Irreversibility of Thermodynamic Process	138
5.7	Carnot Cycle	139
5.8	Reversed Carnot Cycle	142
5.9	Carnot Theorem	142
Сна	PTER 6: ENTROPY	157–182
6.1	Introduction	157
6.2	Importance of Entropy	157
6.3	Units of Entropy	158
6.4	Clausius Theorem	158
6.5	Entropy—A Point Function or a Property of a System	161
6.6	Clausius Inequality	162
6.7	Principle of Increase of Entropy	162
6.8	Entropy and Temperature Relation	163
6.9	General Expression for Change of Entropy of a Perfect Gas	164
6.10	Change of Entropy of Perfect Gas During Various Thermodynamic Process	165
6.11	Irreversibility	167
		107

Сна	PTER 7: PROPERTIES OF PURE SUBSTANCES	183–227
7.1	Introduction	183
7.2	Phase Equilibrium of a Pure Substance on T-V Diagram	183
7.3	Temperature and Total Heat Graph During Steam Formation	185
7.4	Phase Equilibrium at Higher Pressure	185
7.5	Phase Equilibrium at Lower Pressures	187
7.6	Thermodynamic Surface	187
7.7	p−∀ Diagram of a Pure Substances	188
7.8	Important Terms of Steam	189
7.9	Entropy of Steam	213
7.10	Entropy of Water	213
7.11	Entropy Increases During Evaporation	213
7.12	Entropy of Wet and Dry Steam	214
7.13	Entropy for Superheated Steam	214
7.14	External Work of Evaporation	215
7.15	Internal Latent Heat	215
7.16	Internal Energy of Steam	215
7.17	Temperature-Entropy Diagram of Water and Steam	216
7.18	Isothermal Lines on I-S Diagram	216
7.19	Isentropic Lines	216
7.20	Enthalpy–Entropy (n-s) Diagram for water and Steam or Mollier Chart	217
7.21	Dryness Fraction Lines on n-s Diagram	219
7.22	Constant Volume Line	219
7.25	Loothermol Line	219
7.24	Isontronio Line on (h.s.) Diagram	219
7.25	Throttling Lines on h-s Diagram	220
,. <u>2</u> 0		220
CHAI	PTER 8: I HERMODYNAMIC AIR STANDARD CYCLES	229–248
8.1	Introduction	229
8.2	Assumptions in Thermodynamic Cycles	229
8.3	Classification of Thermodynamic Cycles	230
8.4	Important Parameters in Air Standard Cycle Analysis	230
8.5	Important Terms used in Thermodynamic Cycles	230
8.6	Types of Thermodynamic Cycles	232
8.7	Carnot Cycle	232
8.8	Otto Cycle	232
8.9	Joule's Cycle	234
8.10	Diesel Cycle	235
8.11	Comparison between the Efficiency of Otto and Diesel Cycle for same Compression Ratio	0 237
Сна	PTER 9: STEAM POWER CYCLE	249–258
9.1	Introduction	249
9.2	Rankine Cycle	249
9.3	Vapour Compression Refrigeration Cycle	251
	BIBLIOGRAPHY	259

(xiv)

Group–B: Fluid Mechanics

Сна	PTER 1: INTRODUCTION AND FUNDAMENTAL CONCEPTS	263–288
1.1	Definition of Fluid	263
1.2	Fluid Mechanics and its Perview	263
1.3	Fluid as a Continuum	263
1.4	Properties of Fluid	264
Сна	PTER 2: Fluid Statics	289–327
2.1	Forces on Fluid Element	289
2.2	Equilibrium of Static Fluid Element	290
2.3	Solution of Euler's Equations	291
2.4 2.5	Gauge Pressure and Absolute Pressure	294 205
2.5	Manometer	293 296
2.0	Hydrostatic Force on Submerged Plane Surface	298
2.8	Hydrostatic Force on Submerged Curved Surface	300
2.9	Buoyancy and Archimedes Law	301
2.10	Equilibrium and Metacentre	302
Сна	PTER 3: KINEMATICS OF FLUID FLOW	329–355
3.1	Introduction	329
3.2	Scalar and Vector Field	329
3.3	Description of Fluid Flow	329
3.4	Classification of Flow	330
3.5	Description of Flow Patterns	333
3.6	Conservation of Mass	335
Сна	PTER 4: DYNAMICS OF FLUID FLOW	357–392
4.1	Introduction	357
4.2	Equation of Steady Motion Along Streamline	358
4.3	Bernoulli's Equation	359
4.4	Different Heads	360
4.5	Hydraulic Grade Line (HGL) and Energy Grade Line (EGL)	360
4.6	Major and Minor Head Loss	361
4./ 1 9	Absolute and Kelative Kougnness Bounold's Number	301
4.0 4 0	Application of Bernoulli's Equation	362
4.10	Static Pressure and Stagnation Pressure	365
4.11	Pitot Tube	366
	BIBLIOGRAPHY	393
	QUESTIONS PAPER	395–408
		409–416

Group–A *Thermodynamics*



BASIC CONCEPT AND SOME DEFINITIONS

1.1 INTRODUCTION

Thermodynamics is the science of the conversion of energy, in terms of heat and work and their mutual relationship along with properties of the system. It also can be defined as the science of Energy and Entropy. The principles of energy conversion have been formulated into different thermodynamic laws, which are known as *Zeroth Law, First, Second* and *Third Law* of Thermodynamics.

The field of Engineering Science which deals with the applications of thermodynamics and its laws to energy conversion devices, in order to understand their function and improve their performance, is known as *Thermal Engineering*.

A system, which converts heat into mechanical works or vice versa, is known as *Heat Engine*. In the heat engine, heat is generated by the combustion of fluid, which may be solid, liquid or gas. In *external combustion engines* (i.e., steam engines or steam turbines) solid fuel is used as working substance and combustion takes place outside of the engine. In *internal combustion engine* a mixture of air and fuel is used as working substance and combustion takes place inside the engine's cylinder.

1.2 MACROSCOPIC AND MICROSCOPIC APPROACH

There are two point of views to study a thermodynamics problem from which the behaviour of matter can be studied. They are known as *macroscopic approach* and *microscopic approach*.

In the *macroscopic approach*, we study the gross or time averaged effects of the particles which may be observable and measurable by instruments. The macroscopic approach is used in *Classical Thermodynamics* which is the subject matter of the text. In this point of view we deal with volumes that are considerably large compared to molecular dimensions. It is not concerned with the behaviour of individual molecules. Therefore, it treats the matter as continuous, or the whole of this as a *Continuum*.

For example, when a container contains gas, the gas exerts pressure on the walls of its container. The pressure results from the change in momentum of the gas molecules as they collide against the wall of the container. In this approach, it is not concerned with the collisions of the molecules, but with the time averaged value of force exerted on the unit area of the surface of the container, which can be measured by a pressure gauge.

In *microscopic approach*, we make an attempt to analyse system by considering it as comprising of discrete particles which are its atoms and molecules. It is difficult to adopt it in practice. The *modified microscopic approach* is employed to simplify this problem in which we deal with average value for all particles under consideration making use of the *theory of probability*. This modified approach is employed in statistical thermodynamics and kinetic theory of gases. It is particularly helpful, when dealing with the system in which the mean free path of the molecules is large compared with the dimension of the system such as in high vacuum technology.

1.3 THERMODYNAMIC SYSTEM

Thermodynamics dealing with the energy transfer within a definite quantity of matter of fixed mass or definite area or region in space or control volume which is bounded by a closed surface upon which attention is focused is known as system. The systems are shown in Venn diagram in figure 1.1.





- (a) Boundary: Thermodynamic system (or simply known as system) is bounded by a closed surface which separates it from the other systems and universe. This closed surface is known as boundary or system boundary. These boundaries may be fixed or movable. A tank of certain mass of compressed fluid and a certain volume of water in a pipe/line are examples of fixed and movable boundary, respectively. Boundary may be real or imaginary in nature.
- (b) Surrounding: The space or anything outside the system boundaries is known as surrounding.
- (c) Universe: Systems and surroundings together is known as universe.

1.4 CLASSIFICATION OF THERMODYNAMIC SYSTEMS

The thermodynamic systems may be classified into three groups:- (a) Closed system, (b) Open system, and (c) Isolated system.

1.4.1 Closed System

In this system, a definite quantity of matter of fixed mass and identity is considered whose boundaries are determined by the space of the matter occupied in it. A closed system is shown in figure 1.2(a). In this system energy can cross its boundary but the system does not permit any mass transfer along its boundary. In other words, the heat and work energy can cross the boundary of the system during the process, but during the process there is no addition or loss of the original mass of the working substance. Hence, the mass of the working substance which occupies in the control volume of the system, is fixed.



1.4.2 Open System

In this system, both energy (heat and work energy) and mass of the working substance can cross the boundary of the system, during the process. The mass of the working substance within the control volume of system may not be constant. It is shown in figure 1.2(b).





1.4.3 Isolated System

In this system, neither energy (heat energy & work energy) nor mass of the working substance can cross the system boundary during the process. During the process ongoing, there is no addition or loss of the energy and mass of working substance which occupy within the control volume. It is shown in figure 1.2.(c).



1.5 CONTROL MASS AND CONTROL VOLUME

The region in space of definite quantity of matter of fixed mass and identity is bounded by a closed surface (boundary), upon which attention is focused for study. Mass inside the system boundary is known as control mass as shown in figure 1.3(a). In closed system the mass does not cross the system boundary, so mass of system is fixed, though its volume can change against a flexible boundary. Therefore fixed mass of closed system is called control mass.





A control volume is a region in space bounded by a closed envelope on which attention is focused for energy analysis. The control volume bounded by a prescribed boundary is known as control surface as shown in figure 1.3(b). The control volume need not be fixed in size as well as shape. It may not be fixed in position also. However, in most of the applications we deal with control volume which are fixed in size and shape as well as fixed in position relative to the observer. An open system is equivalent in every respect to a control volume.



Fig. 1.3 (b)

6

1.5.1 Comparison of Properties of a Control Mass and a Control Volume

Control mass	Control volume
1. Refers to a definite quantity of matter on which attention is focused.	 Refers to a defined region of space on which attention is focused.
2. Bounded by a closed boundary which may be real or imaginary.	2. Enclosed by a control surface which may be real or imaginary.
3. Matter does not cross the boundaries of a control mass.	3. Matter continuously flows in and out of the control volume.
 Heat and work interaction are present across the system boundary. 	4. Control volume can exchange heat and work through control surface.

1.6 THERMODYNAMIC CO-ORDINATES

Every system has certain observable characteristics by which its physical condition may be described. Such characteristics like location, pressure, volume, temperature etc are called properties. Properties are the co-ordinates to describe the state of the system. These co-ordinates for locating thermodynamic state points are defined as the thermodynamic co-ordinates.

1.7 STATE OF A SYSTEM

The state of a system is the configuration of the system at any particular moment which can be identified by the statement of its properties, such as pressure, volume, temperature etc., so that one state point may be distinguished from the other states.

1.8 PROPERTIES OF A SYSTEM

Properties of a system are thermodynamic co-ordinates of the system which describe its physical condition such as location, volume, pressure, temperature etc., So any observable characteristics of the system is a property. A sufficient number of independent properties exist to represent the state of a system. A property of the system depends only upon the state of the system and not upon how that state may have been reached. The properties are measurable quantities in terms of numbers and units of measurements.

Mathematically it can be explained whether any quantity is a property or not as follows.

If x and y are two properties of a system, then dx and dy are their exact differentials. In general,

if the differential is of the form of
$$(Mdx + Ndy)$$
, the test for exactness is $\left(\frac{\partial M}{\partial_y}\right)_x = \left(\frac{\partial N}{\partial_x}\right)_y$

1.9 CLASSIFICATION OF PROPERTIES OF A SYSTEM

The thermodynamic properties of a system may be classified into the following two groups: (a) Extensive properties and (b) Intensive properties.

1.9.1 Extensive Properties

The properties of a system, whose values for the entire system depending on the mass of the system are equal to the sum of their values for the individual parts of the system are called extensive properties. Total volume, total energy and total mass of the system are extensive properties of the system.

The ratio of an extensive property to the mass is called the specific property or the property per unit mass. For example

Specific volume $(v_s) = \frac{\text{total volume}}{\text{total mass}} = \frac{\forall}{m} \text{ m}^3/\text{kg}$

The ratio of an extensive property to the mole number is called the molar property, like

Molar volume $(v_m) = \frac{\text{total volume}}{\text{mole number}} = \frac{\forall}{n} \text{ m}^3/\text{kg}$

1.9.2 Intensive Properties

The properties of a system where value do not depend upon the mass of the system are called intensive properties, such as temperature of the system. It does not depend on the mass of the system and whatever remains the mass of the system it is same for the entire system. It is also true for specific volume, density and pressure of the system.

1.9.3 Difference between Extensive and Intensive Properties with an Example

The properties which depend on the mass of a system are called extensive properties, whereas the properties which do not depend on the mass of a system are called intensive properties. So intensive properties are independent of the mass of the properties. At normal temperature and pressure i.e., at 0°C and 1 standard atmospheric pressure (at NTP) one mole of a gas occupies a volume of 22.4 litres. Consider two systems *P* and *Q* such that system *P* contains 1 mole of oxygen at NTP and system *Q* contains 5 moles of oxygen at NTP. If the volume occupied by the systems *P* and *Q* can be measured, it will be found that the system *Q* occupies a volume of 22.4 × 5 = 112.0 litres which is equal to five times the volume occupied by system P. So the volume occupied by a system depends on the mass of the system and hence volume is an extensive thermodynamic property. Measurement of temperature and pressure show that systems *P* and *Q* will have the same temperature (0°C) and pressure (one standard atmospheric pressure). Hence, the properties pressure and temperature do not depend on the mass of the system and so are called intensive properties.

1.10 PHASE

When a quantity of matter is homogeneous throughout in terms of chemical composition and physical structure is called a phase. There are three phases, solid, liquid and gas. A system consisting of a single phase is called a *homogeneous system*, while a system consisting of more than one phase is known as a *heterogeneous system*.

1.11 THERMODYNAMIC EQUILIBRIUM

A system is said to be in thermodynamic equilibrium, when no change in any one of the properties do occur. A system will be in a state of thermodynamic equilibrium, if it satisfies the following three

requirements of equilibrium (a) Mechanical equilibrium, (b) Chemical equilibrium (c) Thermal equilibrium.

1.11.1 Mechanical Equilibrium

The system is said to be in mechanical equilibrium, when there is no unbalanced force within the system.

1.11.2 Chemical Equilibrium

The system is said to be in chemical equilibrium, if there is no chemical reaction and no transfer of matter from one part of the system to another.

1.11.3 Thermal Equilibrium

When a system is in contact with its surroundings across a diathermal wall(wall through which heat can flow) and if there is no spontaneous change in any of the properties of the system, the system is said to exist in thermal equilibrium with its surroundings. Also, a system may exist in thermal equilibrium with another system. Adjacent figure 1.4 illustrates thermal equilibrium between two systems or between a system and its surroundings, temperature being one property must be the same.



1.12 PATH

Properties are the thermodynamic co-ordinates of the state of a system. So the properties are statevariable of the system. When any one or more of the properties of a system change, it is called a change of state. When a system passes through a series of states during a change of state from the initial state to the final state, it is called the path of the change of state.

1.13 PROCESS

When a system passes through a successive states during a change of state from the initial state to the final state, with a completely specified path for each successive change in states, the change of state is defined as a process, e.g., a constant volume process, constant pressure process. It is shown in figure 1.5, where 1-2 is a constant volume process and 2-3 is a constant pressure process.

A process is designated by the path followed by the system in reaching the final equilibrium state from the given initial state.



1.14 CYCLIC PROCESS OR THERMODYNAMIC CYCLE

When a process or processes are performed on a system in such a way that the initial and the final states will be same, then the process is called thermodynamic cycle or cyclic process. In figure 1.6 1-*A*-2 and 2-*B*-1 are two simple processes whereas 1-*A*-2-*B*-1 is a cyclic process, whose final and initial states are the same.



1.14.1 Quasi-static Process or Quasi-equilibrium Process

Quasi-static process is a process carried out in such a way that at every instant, the deviation of the state from the thermodynamic equilibrium will be infinitesimally small. Every state passing through the system will be in closely approximating succession of equilibrium state. The locus of all these equilibrium points passed through the system is a quasi-static process. A quasi-static process is shown in figure 1.7. *AB* is a quasi-static process and at the successive states (shown by the points at 1, 2, 3 etc.,) the system is very nearly of thermodynamic equilibrium. Only Quasi-static process can be represented on a thermodynamic plane.



1.14.2 Non-equilibrium Process

Non-equilibrium process is a process carried out in such a way that the initial state-point and the final state-point are in equilibrium but the intermediate state-points, through which the system is passing, are in non-equilibrium state. Figure 1.8 shows the non-equilibrium process whose initial and final equilibrium states are joined by a dotted line which has got no-meaning otherwise.



1.14.3 Reversible Process

Reversible process is a process carried out in such a way that at every instant, the system deviation is only infinitesimal from the thermodynamic state, and also which can be reversed in direction and the system retraces the same equilibrium states. Thus in reversible process, the interactions between the system and the surroundings are equal and opposite in direction. The Quasi-static or Quasi-equilibrium process is also known as reversible process. In reversible process the work done could be written in the form $W = \int p d \forall$ When there is a change in system boundaries.

1.14.4 Irreversible Process

A process is said to be irreversible, while initial and final states both being in equilibrium, when reversed, the system and the surroundings do not come to the original initial state and a trace of history of the forward process is left. In actual practice, most of the processes are irreversible, due

to turbulence in the system, temperature gradients in the system and due to friction. In irreversible processes, the network output is less than $\int pd \forall$ and is given by $W_{\text{net}} = \int pd \forall - W_{\text{dissipated}}$ Reversible and Irreversible processes are shown in figure 1.9.



1.14.5 Flow Process

The process occuring in the control volume of open system which permits the transfer of mass to and from the system is known as flow process. The working substance, in flow processes, enters the system and leaves after doing the work. The flow processes may be classified as (1) steady flow processes and (2) unsteady flow processes. The conditions which must be satisfied for a steady flow process are as following:

- (*i*) The mass flow rate through the system remains constant.
- (ii) The rate of heat transfer is constant.
- (*iii*) The rate of work transfer is constant.
- (*iv*) The characteristics of the working substance, like velocity, pressure, density etc., at any point do not change with time.

If any one of these conditions are not satisfied, then the flow process is said to be an unsteady flow process.

1.14.6 Non-flow Process

The process in which mass of working substance is not permitted to cross the boundary of the control volume of the system, is called non-flow process. Generally non-flow processes occur in the closed system.

1.15 POINT FUNCTION AND PATH FUNCTION

The values of properties of a system in a given state are independent of the path followed to reach that final state from initial state. Because of this characteristics, the state is a point function. Consider the change of state of a system from an initial equilibrium state 1 to a final equilibrium state 2 by following many different paths such as 1-A-2 or 1-B-2 or 1-C-2 or 1-D-2 etc, as shown in figure1.10. The values of properties of states 1 and 2 are p_1 , \forall_1 and p_2 , \forall_2 . respectively. The state 2, whether is reached via 1-A-2 or 1-B-2 or 1-C-2 etc, will have the same values of properties p_2 and \forall_2 . So the change in the value of a property between any two given states is the same irrespective of the path

between the two states. Thus in the present illustration, the differences of values of properties, i.e., $\int_{1}^{2} dp = p_2 - p_1$ and $\int_{1}^{2} d\forall = \forall_2 - \forall_1$ are always the same between the states 1 and 2 whether the path followed is 1-*A*-2 or 1-*B*-2 or any other one. It is also true for any quantity being a point function, the change of which is independent of path.

The path followed to reach the final state from a initial state is called the process. The quantity, the value of which depends on the path followed during a change of state is a path function. For example, in the above figure, the areas under 1-*A*-2, 1-*B*-2 etc, are different i.e.,



Where initial value A_1 and the final value A_2 have no meaning. The integral value of such quantities as

$$\int_{1}^{2} dA = {}_{1}A_{2} \text{ or } A_{12}$$

The area for process between state 1 and 2 can only be determined, when the path followed is known.

1.16 UNIT

The primary quantities are measured in terms of the basic or fundamental units and the secondary quantities are measured in terms of derived units.

1.16.1 Fundamental Units

Fundamentals units are the basic unit normally which are unit of mass (M), unit of length (L) and unit of time(T), but in the International System of units, there are seven fundamental units and two supplementary units, which cover the entire field of science and engineering. These units are shown in the Table 1.1.

Sl.No.	Physical Quantity	Unit	Symbol
1.	Mass (M)	kilogram	kg
2.	Length (L)	metre	m
3.	Time (t)	second	S
4.	Temperature (T)	kelvin	K
5.	Electric current (I)	ampere	А
6.	Luminous intensity (I_v)	candela	Cd
7.	Amount of substance (n)	mole	mole

 Table: 1.1. Fundamental and Supplementary Units

 Fundamental Unit:

Supplementary Unit:

1.	Plane angle $(\alpha, \beta, \theta, \phi)$	radian	rad
2.	Solid angle (Ω)	steradian	Sr

1.16.2 Derived Units

Some units expressed in terms of other basic units, which are derived from fundamental units are known as derived units. The derived units, which will be commonly used in this book, are given in the following table.

Sl.No.	Physical Quantity	Unit	Symbol
1.	Area	m ²	А
2.	Angular velocity	rad/s	ω
3.	Angular acceleration	rad/s ²	α
4.	Linear velocity	m/s	V
5.	Linear acceleration	m/s ²	а
6.	Mass density	kg/m ³	ρ
7.	Force, weight	N	F, W
8.	Work, energy, enthalpy	J	W, E, H
9.	Pressure	N/m ²	р
10.	Power	Watt	Р
11.	Absolute or dynamic viscosity	N-s/m ²	μ
12.	Kinematic viscosity	m²/s	υ
13.	Characteristic gas constant	J/kg·K	R

Table: 1.2. Derived Units

14

Contd..

14.	Universal gas constant	J/kgmol.K	R _m
15.	Frequency	Hz, 1 Hz = 1cps	f
16.	Thermal conductivity	W/mK	k
17.	Specific heat	J/kg.K	С
18.	Molar mass or molecular mass	kg/mol	М
19.	Sp. weight or wt. density	kgf/m ³	W _S
20.	Sp. volume	m³/kg	V _S
21.	Volume	m ³	\forall

1.17 SYSTEMS OF UNITS

There are only four systems of units, which are commonly used and universally recognized. These are known as:

(1) C.G.S systems (2) F.P.S systems (3) M.K.S systems (4) S.I (System Internationale or International system of units). The internationally accepted prefixes in S.I to express large and small quantities are given below in the table.

Factor of Multiplication	Prefix	Symbol
1012	tera	Т
10^{9}	giga	G
10^{6}	mega	М
10 ³	kilo	k
12 ²	hecto	h
10 ¹	deca	da
10-1	deci	d
10-2	centi	с
10-3	milli	m
10-6	micro	m
10-9	nano	n
10 ⁻¹²	pico	р

Table: 1.3. Prefix Factors

1.18 MASS (M)

Mass is the amount of matter contained in a given body and it does not vary with the change in its position on the earth's surface.

1.19 WEIGHT (W)

The weight is the amount of force of attraction, which the earth exerts on a given body. The weight of the body will vary with its position on the earth's surface, because force of attraction vary with variation of distance between the two bodies.

1.20 FORCE (F)

Force may be defined as an agent which produces or tends to produce, destroy or tends to destroy the motion. According to Newton's Second Law of Motion, the applied force or impressed force is directly proportional to the rate of change of momentum

Hence,

$$F \propto \frac{mv - mu}{t}$$

$$\Rightarrow \qquad F \propto m\left(\frac{v - u}{t}\right)$$

$$\Rightarrow \qquad F \propto m a$$

$$\therefore \qquad F = k \cdot m \cdot a$$

Where k is the constant of proportionality. If the unit of force adopted so that it produces unit acceleration to a body of unit mass, then

or,	Force = mass × acceleration
\Rightarrow	F = ma
\Rightarrow	k = 1
.: .	$1 = k \cdot 1 \cdot 1$

In S.I the unit is newton(N) and $1N = 1 \text{ kg m/s}^2$.

There are two types of units of force, absolute and gravitational. When a body of mass 1 kg is moving with an acceleration of 1 m/s^2 , the force acting on the body is 1 newton(N). This is the absolute unit of force.

When a body of mass 1 kg is attracted towards the earth with an acceleration of 9.81m/s², the force acting on the body is 1 kilogram-force, briefly written as 'kgf' or kg-wt. The unit of force in kgf is called gravitational or engineers, unit of force or metric unit of force. From the Newton's Second Law of Motion

$$F = k \cdot m \cdot a = m \cdot \frac{g}{g_c} = \text{weight}$$

 $k = \frac{1}{2}$ while

Where

$$a = g$$
 (acceleration due to gravity), and

$$g_c = 9.80665 \text{ kg} \frac{\text{m}}{\text{s}^2} \approx 9.81 \text{ kg} \cdot \frac{\text{m}}{\text{s}^2} \cdot \frac{1}{\text{kgf}}$$

and

$$1 \text{ kgf} = \frac{1 \text{ kg} \times 9.81 \text{ m/s}^2}{g_c} = \text{weight}$$

If, local value of g is numerically the same as g_c then the weight of 1 kg becomes equal to 1 kgf.

The gravitational unit of force is g' times greater than the absolute unit of force or S.I unit of force, as

$$1 \text{ kgf} = 1 \text{ kg} \times 9.81 \text{ m/s}^2 = 9.81 \text{ m/s}^2 = 9.81 \text{ N}$$

1.21 SPECIFIC WEIGHT (w_s)

It is the weight per unit volume. It is also known as the weight density. It may be expressed in kgf/m³, in MKS system of unit and newton/m³ in S.I.

Specific weight
$$(w_s) = \frac{W}{\forall} = \frac{mg}{\forall} = \rho.g$$

1.22 SPECIFIC VOLUME (v_s)

It is defined as the volume per unit mass. It may be expressed in m³/kg.

Specific volume
$$(v_s) = \frac{\forall}{m} = \frac{1}{\rho}$$

1.23 PRESSURE

Pressure is the normal force exerted by a system against unit area of the boundary surface. The unit of pressure depends on the units of force and area. In S.I, the practical units of pressure are N/mm², N/m², kN/m², MN/m² etc.

A bigger unit of pressure known as bar, such that

1 bar = $1 \times 10^5 \text{ N/m}^2 = 0.1 \times 10^6 \text{ N/m}^2 = 0.1 \text{ MN/m}^2$

Other practical units of pressure are Pascal (Pa), kilopascal (kPa) & mega Pascal (MPa), such that

1 Pa =
$$1N/m^2$$

1 kPa = $1 \text{ kN/m}^2 = 10^3 \text{ N/m}^2$
1 MPa = $1 \times 10^6 \text{ N/m}^2 = 10^3 \text{ kPa} = 1 \text{ N/mm}^2$

1.24 ABSOLUTE, GAUGE AND VACUUM PRESSURE

The pressure is measured in two different systems. In one system, it is measured above the absolute zero or complete vacuum, and is defined as absolute pressure. In other system, pressure is measured above the atmospheric pressure, and is defined as gauge pressure. So

- (a) Absolute pressure: It is defined as the pressure which is measured with reference to absolute zero pressure.
- (b) Gauge pressure: It is defined as the pressure which is measured with reference to atmospheric pressure. It is measured with the help of a pressure measuring instrument. It is a pressure above the atmospheric pressure.

(c) Vacuum pressure: It is the pressure below the atmospheric pressure. Sometimes it is called as negative gauge pressure. The relationship between the absolute pressure, gauge pressure and vacuum pressure are shown in figure 1.11.





Mathematically

(*i*) Absolute pressure = atmospheric pressure + gauge pressure

 P_{abs} (*ii*) Vacuum pressure

 $p_{abs} = p_{atm} + p_{gauge}$ ssure = atmospheric pressure – absolute pressure = $p_{atm} - p_{abs}$

1.25 PRESSURE MEASUREMENT BY MANOMETER

A manometer is normally used to measure pressure. In manometer the pressure is determined according to the hydrostatic formula. The manometric liquid may be mercury, water, alcohol, etc. A *U*-tube manometer is shown in figure 1.12. Since manometric fluid is in equilibrium, the pressure along a horizontal line AB is the same for either limb of manometer, then

$$p + \frac{\rho_1 g \, z_1}{g_c} = p_{\text{atm}} + \frac{\rho_2 g \, z_2}{g_c}$$

Where p is the absolute pressure in the bulb, p_{atm} is the atmospheric pressure exerted on the free surface of liquid and ρ_1 and ρ_2 are the densities of the liquid in the bulb and manometer respectively.

If ρ_1 is small as compared to ρ_2 i.e., $\rho_1 << \rho_2$, $\frac{\rho_1 \cdot g \ z_1}{z_c} \approx 0$

then

$$p - p_{\text{atm}} = \frac{\rho_2 \cdot g \, z_2}{g_c} = p_{\text{gauge}}$$





Table: 1.4. Conversion Factor for Pressure

	Bar	dyne/cm ²	kgf/cm² or ata	N/m ² or Pa	mm Hg at 21°C or torr	mm H ₂ O at 21°C	atm
Bar	1	106	1.01972	105	750.062	10197.2	0.986923
dyne/cm ²	10-6	1	1.01972 × 10 ⁻⁶	0.1	750.062 × 10 ⁻⁶	10197.2 × 10 ⁻⁶	0.986923 × 10 ⁻⁶
kgf/cm ² or ata	0.980665	0.980665 × 10 ⁶	1	0.980665	735.559	10000	0.967838
N/m ² or Pa	10-5	10	10.1972 × 10 ⁻⁶	1	750.062 ×10 ⁻⁵	10197.2 ×10 ⁻⁵	0.986923 ×10 ⁻⁵
mm of Hg at 21°C	1.333233 × 10 ⁻³	1.333233 × 10 ³	1.3595 × 10 ⁻³	1.333233	1	13.5951	1.31578 × 10 ⁻³
mm of H ₂ O at 21°C	98.0665 × 10 ⁻⁶	98.0665	10-4	9.80665	0.073556	1	96.7838 × 10 ⁻⁶
atm	1.01325	1.01325 × 10 ⁶	1.03323	1.01325 × 10 ⁵	760	103523	1
1.26 NORMAL TEMPERATURE AND PRESSURE (N.T.P)

Normal temperature is at 0°C or 273 K temperature and normal pressure is 760 mm of Hg. Normal temperature and pressure are briefly written as N.T.P.

1.27 STANDARD TEMPERATURE AND PRESSURE (S.T.P)

The temperature and pressure of any gas, under standard atmospheric condition, is taken as 15°C (288K) and 760 mm of Hg respectively.

1.28 ENERGY

The simplest definition of energy is the capacity for doing work. In other words, a system is said to posses energy when it is capable of doing work. The energy can be classified as (i) Stored energy and (ii) Transit energy.

The stored energy is a thermodynamic property as it depends on the point, not upon the path. The stored energy is the energy which is contained within the system boundaries. Examples of stored energy are (i) potential energy (ii) kinetic energy (iii) internal energy etc.

The transit energy is in transition and crosses the system boundaries. Examples of transit energy are (i) heat (ii) work (iii) electrical energy etc. The transit energy is not a thermodynamic property as it depends upon the path.

1.29 TYPES OF STORED ENERGY

The potential energy, kinetic energy or an internal energy are the different types of stored energy and are discussed in detail, as follows:

1.29.1 Potential Energy

The energy possessed by a body, or a system for doing work, by virtue of its location or configuration is called potential energy. If a body of mass m is at an elevation of z above the datum plane, the potential energy(P.E) possessed by the body is given by

$$PE = mgz = W \cdot z$$

where *g* is the acceleration due to gravity.

1.29.2 Kinetic Energy

The energy possessed by a body, or a system for doing work, by virtue of its motion is called kinetic energy. If a body of mass of m moves with a velocity v the kinetic energy(KE) possessed by the body is given by

$$KE = \frac{1}{2}mv^2$$

The sum of the potential energy and kinetic energy of a body is called the Mechanical energy of the body.

1.29.3 Internal Energy

This energy is possessed by a body, or a system due to its molecular arrangement and motion of the molecules. It is usually represented by U and the change in internal energy (dU). It depends upon the change in temperature of the system.

:. change of internal energy $dU = C_v (T_2 - T_1)$

Where C_v is specific heat at constant volume &

20

 T_1 and T_2 are the temperature at state points

 \therefore The total energy of the system (*E*) is equal to the sum of the P.E, K.E and internal energy.

 $\therefore \qquad E = P.E + K.E + U$

[any other form of the energy such as chemical, electrical energy etc. are neglected]

Again

$$E = mgz + \frac{1}{2}mv^2 + U = ME + U$$

while Mechanical energy $(M.E) = P.E. + K.E = mgz + \frac{1}{2}mv^2$

For unit mass, total energy

$$e = gz + \frac{1}{2}v^2 + u$$

When the system is stationary and the effect of gravity is neglected, then,

$$E = U$$
 and $e = u$

1.30 LAW OF CONSERVATION OF ENERGY

The law of Conservation of Energy states that *The energy can neither be created nor destroyed, though it can be transformed from one form to any other form, in which the energy can exist.*

1.31 POWER

:.

Power may be defined as the rate of doing work or work done per unit time or rate of energy transfer or storage. Mathematically,

Power =
$$\frac{\text{work done}}{\text{time taken}}$$

= $\frac{\text{energy storage or transfer}}{\text{time taken}}$

The unit of power in S.I is watt (W)

1

$$W = 1 N \cdot m/s = 1 J/s$$

A bigger unit of power called kilowatt (kW) or megawatt (MW)

$$1000 \text{ W} \text{ and } 1 \text{ MW} = 10^6 \text{ W} = 1000 \text{ kW}$$

If T is the torque transmitted expressed in N·m or J and the angular speed is ω in rad/s, then

Power
$$(P) = T \times \omega$$

1 kW =

$$= T \times \frac{2\pi N}{60} \text{ watt, } \omega = 2\pi N/60$$

N is speed in r.p.m

Hence, Efficiency $(\eta) = \frac{power output}{power input}$

Multiple Choice Questions

1.	In a	a closed thermodynamic system there is		
	(a)	only mass transfer	(b)	only energy transfer
	(c)	both mass and energy transfer	(d)	none of the above
2.	In f	flow system there is		
	(a)	no mass transfer across the boundaries	(b)	neither mass nor energy transfer
	(c)	both mass and energy transfer	(d)	none of the above
3.	In a	an isolated system there is		
	(a)	no mass transfer	(b)	no energy transfer
	(c)	neither mass nor energy transfer	(d)	both mass and energy transfer
4.	Wh	nich of the following are the properties of t	he sy	stem?
	(a)	∫pd∀	(b)	$\int \forall dp$
	(c)	$\int (pd\forall + \forall dp)$	(d)	$\int dT/T = C_{V} dp/T$
5.	Wh	ich of the following is an intensive proper	ty?	
	(a)	volume	(b)	temperature
	(c)	density	(d)	entropy
6.	Wh	nich of the following is an extensive proper	ty?	
	(a)	pressure	(b)	temperature
	(c)	density	(d)	volume
7.	1 tc	orr is equivalent to		
	(a)	1 kgf/cm ²	(b)	1 N/m ²
	(c)	1 atm	(d)	1mm of Hg
8.	The	e expression $\int pd \forall$ may be applied for obtain	ning	work of
	(a)	non-flow reversible process	(b)	steady flow reversible process
	(c)	steady flow non-reversible process	(d)	steady flow adiabatic reversible process
9.	Eac	ch of heat and work is		
	(a)	point function	(b)	path function
	(c)	property of a system	(d)	state description of a system
10.	The	e property which depends only on tempera	ture i	IS
	(a)	internal energy	(b)	enthalpy
	(c)	entropy	(d)	none of above
11.	Ad	lefinite area or a space where some thermo	dyna	mic process takes place is known as
	(a)	thermodynamic cycle	(b)	thermodynamic process
	(c)	thermodynamic system	(d)	thermodynamic law
12.	Wh	nen either of mass or energy is not allowed	to cr	ross the boundary of a system; it is then called
	(a)	closed system	(b)	open system
	(c)	isolated system	(d)	none of these

- 13. Which of the following is not a thermodynamic property?
 - (a) pressure (b) temperature
 - (c) heat (d) specific volume
- 14. When a process or processes are performed on a system in such a way that the final state is identical with the initial state, it is then known as
 - (a) thermodynamic cycle
 - (b) thermodynamic property
 - (c) thermodynamic process
 - (d) zeroth law of thermodynamics
- 15. Atmospheric pressure is equal to
 - (a) 1.013 bar
 - (c) 760 mm of Hg
- 16. 1 mm of H₂O is equal to
 - (a) 100×10^{-6} bar
 - (c) 9.80665 Pa
- 17. -40° C is equal to
 - (a) -40° F (1)
 - (c) 400°R
- 18. A centrifugal fan forms
 - (a) closed system (b) open system
 - (c) isolated system (d) none of the above
- 19. Which of the following statement is correct?
 - (a) isolated system uninfluenced by surrounding is called universe.
 - (b) system and surrounding combine to constitute universe whether there are interaction with each other or not.
 - (c) system which only interacts with surrounding is part of the universe.
 - (d) system and surrounding put together form universe only if there is interaction between them.
- 20. Thermodynamic system may be defined as a quantity of matter upon which attention is focussed for study if
 - (a) it is only bounded by real surface
 - (b) the boundary surfaces are constant in shape and volume.
 - (c) it is not bounded by imaginary surface
 - (d) it is bounded by either real surfaces or imaginary surfaces irrespective of shape or volume.

			nswers		
3. (c)	4. (c)	5. (b)	6. (d)	7. (d)	8. (a)

1. (b)	2. (c)	3. (c)	4. (c)	5. (b)	6. (d)	7. (d)	8. (a)	9. (b)	10. (a)
11. (c)	12. (c)	13. (c)	14. (a)	15. (d)	16. (c)	17. (a)	18. (b)	19. (b)	20. (d)

- (b) 101.3 kN/m^2
- (d) all of the above
- (b) 0.001 kgf/ cm^2
- (d) 0.077 mm Hg
- (b) 230K
- (d) -72°F

NUMERICAL EXAMPLES

EXAMPLE 1

If $pv_s = RT$ (where $v_s = specific$ volume, p = pressure, R = a constant and T = temperature) determine whether the following quantities (i) $\int \left(\frac{dT}{T} - \frac{v_s dp}{T}\right)$ and (ii) $\int \left(\frac{dT}{T} + \frac{p dv_s}{v_s}\right)$ can be used as properties.

SOLUTION

Each of the differential is of the form (Mdx + Ndy).

Therefore, apply the test
$$\left(\frac{\partial M}{\partial y}\right)_x = \left(\frac{\partial N}{\partial x}\right)_y$$

[Condition of exact differential equation]

(i) Thus, for
$$\left(\frac{dT}{T} - \frac{v_s dp}{T}\right)$$
 we can write the condition

$$\left[\frac{\partial \left(\frac{1}{T}\right)}{\partial p}\right]_T = \left[\frac{\partial \left(\frac{-v_s}{T}\right)}{\partial T}\right]_p$$
or,
$$\left[\frac{\partial \left(\frac{1}{T}\right)}{\partial p}\right]_T = \left[\frac{\partial \left(\frac{-R}{p}\right)}{\partial T}\right]_p$$
or,
$$0 = 0$$

 $\begin{bmatrix} \text{as } pv_s = RT \\ \therefore \frac{v_s}{T} = \frac{R}{p} \end{bmatrix}$

Thus, $\left(\frac{dT}{T} - \frac{v_s dp}{T}\right)$ is an exact differential equation and may be written as $\frac{dT}{T} - \frac{v_s dp}{T} = dS$ where S is a point function and hence a property.

(ii) For
$$\left(\frac{dT}{t} + \frac{pdv_s}{v_s}\right)$$
, we have

$$\left[\frac{\partial\left(\frac{1}{T}\right)}{\partial v_s}\right]_T = \left[\frac{\partial\left(\frac{p}{v_s}\right)}{\partial T}\right]_{v_s}$$
or,
$$\left[\frac{\partial\left(\frac{1}{T}\right)}{\partial v_s}\right]_T = \left[\frac{\partial\left(\frac{RT}{v_s^2}\right)}{\partial T}\right]_{v_s}$$

or,

$$0 = \frac{R}{v_s^2} \left[\frac{\partial(T)}{\partial T} \right]_{v_s} \qquad \left[A_s \ p = \frac{RT}{v_s} \right]$$
or,

$$0 = \frac{R}{v_s^2}$$

or,

But $\frac{R}{v^2} \neq 0$. So, it is not an exact differential.

Therefore,
$$\left(\frac{dT}{T} + \frac{pdv_s}{v_s}\right)$$
 is not a point function. So, it is not a property

EXAMPLE 2

Explain if the following can be used a properties (i) $\int pd\forall$ (ii) $\int \forall dp$ (iii) $\int (pd\forall + \forall dp)$.

SOLUTION

- (i) $\int p d\forall$ is an expression where p is function of \forall and they are connected by a line path on p and \forall plane. The value $\int pd \forall$ depends on the area under the line path on p and \forall planes. Thus it is not an exact differential, it is a path function, not a point function. So it is not a property.
- (ii) $\int \forall dp$ is such an expression where \forall is a function of p and they are connected by a line path on p and \forall planes. So it is a path function not a point function. Thus, it is not an exact differential and thus not a property.
- (iii) $\int p d \forall + \int \forall dp = \int d(p \cdot \forall) = p \cdot \forall$

So, it is an exact differential and hence it is a property.

EXAMPLE 3

A manometer contains a fluid having a density of 1200 kg/m³. The difference in height of two columns is 400 mm. What pressure difference is indicated thus? What would be the height difference, if the same pressure difference is to be measured by a mercury manometer having mass density of mercury 13600 kg/m3.

SOLUTION

⇒

The pressure difference indicated

 $p_1 = \rho g h = 1200 \times 9.81 \times 400/1000 = 4708.8 \text{ Pa} = 4.7088 \text{ kPa}$ $p_1 = 4708.8$ Again $4708.8 = \rho_g g h_g = 13600 \times 9.81 \times h_g$ \Rightarrow $h_g = \frac{4708.8}{13600 \times 9.81} = 53.3 \text{ mm}$

So, length of mercury column 53.3 mm.

EXAMPLE 4

The pressure of steam inside a boiler, measured by pressure gauge is 1 N/mm². The barometeric pressure of the atmosphere is 765 mm of mercury. Determine the absolute pressure of steam in N/m², kPa, bar and N/mm².

SOLUTION	
Gauge pressure = $1 \text{ N/mm}^2 = 1 \times 10^6 \text{ N/m}^2$	
Atmospheric pressure = 765 mm of Hg	
We know that atmospheric pressure	
$=765 \mathrm{mm}\mathrm{of}\mathrm{Hg}$	
$= 765 \times 133.3 = 0.102 \times 10^6 \text{N/m}^2$	$(:: 1 \text{ mm of Hg} = 133.3 \text{ N/m}^2)$
:. Absolute pressure of steam	
= atmospheric pressure + gauge pressure	
$= 0.102 \times 10^6 + 1 \times 10^6 = 1.102 \times 10^6 \text{N/m}^2$	
= 1102 kPa	$(:: 1 \text{ kPa} = 10^3 \text{ N/m}^2)$
= 11.02 bar	$(:: 1 \text{ bar} = 10^5 \text{ N/m}^2)$
$= 1.102 \mathrm{N/mm^2}$	$(:: 1 \text{ N/mm}^2 = 10^6 \text{ N/m}^2)$

EXAMPLE 5

In a condenser of a steam power plant, the vacuum pressure is recorded as 700 mm of mercury. If the barometer reading is 760 mm of mercury, determine the absolute pressure in the condenser in N/m^2 , kPa, bar and N/mm^2 .

SOLUTION

Vacuum pressure = 700 mm of Hg

Barometer reading = 760 mm of Hg

We know that absolute pressure

= atmospheric pressure – vacuum pressure

= barometer pressure – vacuum pressure

 $= 760 - 700 = 60 \,\mathrm{mm}\,\mathrm{of}\,\mathrm{Hg}$

 $=60 \times 133.3 = 7998 \text{ N/m}^2$

= 7.998 kPa

 $= 0.07998 \, \text{bar}$

 $= 0.007998 \,\text{N/mm}^2$

EXAMPLE 6

Compute the quantity of heat required to raise the temperature of a steel forging of mass 180 kg from 300 K to 1265 K. The specific heat of steel = 0.49 kJ/ kg K.

SOLUTION

Given: m = 180 kg, $T_1 = 300 \text{ K}$, $T_2 = 1265 \text{ K}$, C = 0.49 kJ/kg K

We know that quantity of heat required

= mass × specific heat × rise in temperature = $mC(T_2 - T_1)$ = 180 × 0.49 (1265 - 300) = 85113 kJ

EXAMPLE 7

...

The forced draught fan supplies air to furnace of the boiler at draught of 30 mm of water. Determine the absolute pressure of air supply if the barometer reads 760 mm mercury, in kgf/cm², bar and kPa. **SOLUTION**

Since it is a forced draught, the absolute pressure is above atmospheric pressure.

Absolute pressure = atmospheric pressure + gauge pressure
= barometer reading + forced draught reading
We know, 1 mm of water $\frac{1}{13.595} = 0.073556 \mathrm{mm Hg}$
$30 \mathrm{mm} \mathrm{of} \mathrm{water} = 30 \times 0.073556 \mathrm{Hg}$
$= 2.20668 \mathrm{mm}\mathrm{Hg}$
Thus, absolute pressure = $760 + 2.20668 = 762.20668 \text{ mm Hg}$
$= \frac{762.20668}{735.559} \text{ kgf/cm}^2$
$= 1.03623 \text{ kgf/ cm}^2$
$= 1.03623 \times 0.981 = 1.01654$ bar
$=\frac{1.01654\times10^5}{10^3}=101.654\mathrm{kPa}$

EXAMPLE 8

The pressure in a gas pipe is measured by a mercury manometer as shown in the figure. One leg of manometer is open to atmosphere. If the difference in the height of mercury column in the two legs is 450 mm, compute the gas pressure in the pipe in kPa, bar and in atm. The barometeric reading is 755 mm Hg. The local acceleration due to gravity is 9.81 m/sec² and the mass density of mercury is 13595 kg/m³. **SOLUTION**

Dollo Holi

Referring figure,

The pressure of gas at plane a - a is given by $p = p_a + \rho g h$ [N]

 $p_a = \rho g h_0$

[Neglecting the density of gas as compared to that of mercury]

And

where ρ is density of mercury



Given the vessel shown in the figure fitted with pressure gauges indicating reading of A = 300 kPa and B = 100 kPa gauge. If the barometer reading is 750 mm of Hg, determine the reading of gauge C and convert this value to absolute value.

SOLUTION

Let the pressure in one region be YkPa absolute and the pressure in other region XkPa absolute and barometer

pressure = 750 mm of Hg =
$$\frac{750 \times 13.595 \times 9.81}{10^3}$$
 kPa ≈ 100 kPa
Threaded for connection
Threaded for connection
Closed Box
Y
100 kPa
Thus
 $Y-100$ = reading on $A = 300$ kPa
and
 $Y-X$ = reading on $B = 100$ kPa
 $X-100$ = reading on C
Solving (i) and (ii), $Y = 400$ kPa. (abs) and $X = 300$ kPa. (abs)

Substituting the value of X in (iii) we get reading on C = 200 kPa gauge.

:. Absolute pressure for reading on C = 200 + 100 = 300 kPa

EXAMPLE 10

An inclined manometer with angle of inclination $\theta = 30^{\circ}$ is shown in the figure. It is filled to measure the pressure of a gas. The fluid inside the manometer has a density of 0.75 g/cm³ and the manometer reading is

labelled on the diagram as x = 0.3 m. If the atmospheric pressure is 101.25 kPa and the acceleration due to gravity is g = 9.7 m/sec², determine the absolute pressure of the gas in kPa.

SOLUTION

For inclined manometer, verticle liquid column

$$h = 0.3 \sin 30^\circ = 0.15 \,\mathrm{m}$$

The gauge pressure is given by

$$p_{gauge} = \rho g h$$

$$= \frac{0.75 \times 10^{-3}}{\left(10^{-2}\right)^3} \times 9.81 \times 0.15$$

= 1103.625 Pa= 1.103625 kPa

Thus absolute pressure is given by

$$p_{abs} = p_{gauge} + p_{atm}$$

= 1.103625 + 101.25 = 102.353625 kPa



EXAMPLE 11

Convert (a) 3 kgf/cm^2 absolute to kgf/cm² gauge, (b) 45 cm vacuum to cm of Hg absolute and to kgf/cm² absolute (ata). (c) 0.5 kgf/cm^2 absolute i.e., 0.5 ata to cm of Hg vacuum, (d) 25 cm of Hg gauge to cm of Hg absolute and to atmosphere, (e) 1 ata to kPa. Barometer may be assumed to be 760 mm of Hg. **Solution**

Here	e, $1 \text{ kgf/cm}^2 = 735.559 \text{ mm of Hg} = 735.6 \text{ mm of Hg}$
So,	$760 \mathrm{mm}\mathrm{Hg} = 1.03323 \mathrm{kgf/cm^2} = 1.033 \mathrm{kgf/cm^2}$
(a)	Absolute pressure – Barometer pressure = Gauge pressure
or,	3 - 1.033 = Gauge pressure
<i>.</i> :.	Gauge pressure = 1.967 kgf/cm^2
(b)	Barometer pressure – Absolute pressure = Vacuum pressure
or,	76 cm of Hg – Absolute pressure = 45 cm of Hg
<i>:</i> .	Absolute pressure = $(76 - 45)$ cm Hg
or,	= 31 cm of Hg absolute
(c)	Barometer pressure – Absolute pressure = Vacuum pressure
or,	$1.033 \text{ kgf/cm}^2 - 0.5 \text{ kgf/cm}^2 = \text{Vacuum pressure}$
or,	$0.533 \text{ kgf/cm}^2 = \text{Vacuum pressure}$
÷	Vacuum pressure = 0.533 kgf/cm^2
	$= 0.533 \times 73.56$ cm of Hg vacuum
	= 39.20748 cm of Hg vacuum.

```
(d) Abs pressure – Barometer pressure = Gauge pressure

Abs pressure – 76 cm of Hg = 25 cm of Hg

Abs pressure = 101 cm of Hg absolute

= 101/76 atm = 1.33 atm

(e) 1 ata = 1 kgf/cm<sup>2</sup>

= 0.981 \times 10^2 kPa

= 98.1 kPa.
```

EXERCISE

- 1. Explain the terms (a) thermodynamic state (b) thermodynamic process (c) thermodynamic cycle.
- 2. What do you understand by macroscopic and microscopic view points of thermodynamics?
- 3. What is thermodynamic system? Explain its different types.
- 4. What do you understand by property of a system? Distinguish between extensive and intensive properties of a system with the help of an example.
- 5. What is a thermodynamic process and a cyclic process?
- 6. Explain the non-equilibrium and quasi-static process. Is the quasi-static process a reversible process?
- 7. Distinguish between gauge pressure and absolute pressure. How the gauge pressure is converted into absolute pressure?
- 8. What do you understand by N.T.P and S.T.P? What are their values?
- 9. Compare the control volume and control mass.
- 10. Define energy. What is stored energy and transit energy. Discuss the types of stored energy.
- 11. Distinguish between absolute pressure and gauge pressure. How is one related to the other in case of vacuum.
- 12. Define the following
 - (a) Point function and Path function
 - (b) Specific weight
 - (c) Specific volume
 - (d) Pressure
 - (e) Temperature
 - (f) Density
 - (g) Flow process & Non-flow process.
- 13. Mercury of density 13.59508 g/cm³ is used as monometer fluid. What gauge pressure in bar is exerted by a column of mercury of 760 mm ?

Ans. [1.01325 bar]

14. Which of the following can be used as properties of the system

(a)
$$\int \left(\frac{dT}{T} + \frac{\forall dp}{T}\right)$$

(b)
$$\int \left(\frac{dT}{T} - \frac{\forall dp}{\forall}\right)$$

- 15. A piston has area of 5 cm². What mass must the piston have if it exerts a pressure of 50 kPa above atmospheric pressure on the gas enclosed in the cylinder.
- 16. A pressure gauge reads 2.4 bar and the barometer reads 75 cm of Hg. Calculate the absolute pressure in bar and in the standard atmosphere.
- 17. A manometer has a liquid of density 800 kg/cm², the difference in level of the two legs is 300 mm. Determine the pressure difference read by it in kgf/m²; bar; kPa.
- The pressure of steam inside a boiler is recorded by a pressure gauge as 1.2 N/mm². If the barometer reads the atmospheric pressure as 770 mm of Hg, find the absolute pressure of steam inside the boiler in N/m², kPa and bar.

Ans. [1.3026×10⁶ N/m², 1302.6 kPa, 13.026 bar]

 In a condenser, the vacuum is found to be 145 mm of mercury and the barometer reads 735 mm of mercury. Find the absolute pressure in a condenser in N/m², kPa and N/mm².

Ans. [78647 N/m², 78.647 kPa, 0.078647 N/mm²]

This page intentionally left blank



ZEROTH LAW AND TEMPERATURE

2.1 TEMPERATURE

The temperature is an intensive thermodynamic property of the system, whose value for the entire system is not equal to the sum of the temperature of its individual parts. It determines the degree of hotness or the level of heat intensity of a body or a system. A body is said to be at a high temperature or hot, if it shows high level of heat intensity in it and a body is said to be at a low temperature or cold, if it shows a low level of heat intensity.

2.2 ZEROTH LAW OF THERMODYNAMICS

This law states, "When each of two systems are in thermal equilibrium with a third system, then the two systems are also in thermal equilibrium with one another."

Let a body X is in thermal equilibrium with a body Y, and also separately with a body Z, then following above law, Y and Z will be mutually in thermal equilibrium with each other. A system is said to be in thermal equilibrium, when there is no temperature difference between the parts of the system or between the system and the surroundings. Zeroth law provides the basis of temperature measurement.

2.3 MEASUREMENT OF TEMPERATURE

The temperature of a system is a property that determines whether or not a system is in thermal equilibrium with other system. The temperature of a system or body is measured with the help of an instrument known as *Thermometer*. A thermometer may be in the form of a glass tube containing mercury in its stem. Or any physical body with at least one measurable property that changes as its temperature changes can be used as a thermometer, for example, a length of a column of mercury in an evacuated capillary tube. The height of mercury column in a thermometer, therefore, becomes a thermometric property. There are other methods of temperature measurement which utilize various other properties of a materials, which are functions of temperature as thermometric properties. The particular substance that exhibits changes in thermometric properties is called thermodynamic substance. Commonly used properties of materials employed in temperature-sensing devices or thermometers are given below and the names of the corresponding thermometric properties employed.

Sl.No.	Thermometer	Thermometric Property	Symbol
1.	Mercury alcohol in glass thermometer	Length	L
2.	Constant volume gas thermometer	Pressure	р
3.	Constant pressure gas thermometer	Volume	A
4.	Electric resistance	Resistance	R
5.	Thermocouple	Electromotive force	Е
6.	Radiation (pyrometer)	Intensity of radiation	I or J
7.	Optical pyrometer	Monochromatic radiation	

Table: 2.1. Types of Thermometer

2.4 CONSTANT VOLUME GAS THERMOMETER AND TEMPERATURE SCALE

A schematic diagram of a constant volume gas thermometer is shown in figure 2.1. It is so exceptional in terms of precision and accuracy that it has been adopted internationally as the standard instrument for calibrating other thermometers.

It consists of a constant volume gas bulb which is connected to a mercury column through a flexible tube. The thermometric substance is the gas (normally *hydrogen* or *helium* or some other gas), filled in gas bulb treated as system with system boundaries shown in figure 2.1 and the pressure exerted by it is measured by open tube mercury manometer. When temperature in the system increases, the gas expands and forces the mercury up in the open tube. The constant volume of the gas is maintained by raising or lowering the mercury reservoir. Then the pressure (p) of the gas is measured from the height (L) of the mercury column in the limb. A similar measure is made by immersing the gas bulb in a bath which is maintained at the triple point of the water.



Let p_{tp} denotes the pressure of the gas when bulb is in thermal equilibrium with the bath at the triple point temperature T_{tp} .

For an ideal gas at constant volume

We have,

$$\frac{T}{T_{tp}} = \frac{p}{p_{tp}}$$

:.

 $T = \frac{p}{p_{tp}} \times T$ The triple point temperature of water has been assigned a value of 273.16 K. Therefore, equation (2.1) can be rewritten as

$$T = 273.16 \frac{p}{p_m}$$
(2.2)

The measured pressure (p) at the system temperature (T) as well as the pressure (p_{p}) at the triple point temperature change depending on the quantity of gas in the gas bulb. A plot of T calculated from equation (2.2) as a function of p_{tp} appears as shown in Fig. 2.2. It shows that the value of T and $p_{y} \rightarrow 0$ is identical for different gases. This behaviour can be expected because all gases behave like ideal gases as $p \rightarrow 0$. Hence to obtain the correct temperature one should ensure that p_{tp} is as low as possible. Therefore, to obtain the actual temperature of the system, equation (2.2) is modified as



Fig. 2.2

Therefore, in the limit when the pressure tends to zero, the same value is obtained for each gas. Thus, the temperature scale is defined as

$$T = 273.16 \operatorname{Lt}_{p, p_{p \to 0}} \left(\frac{p}{p_{tp}} \right)$$
(2.3)

(2.1)

The temperature value which has been measured does not depend on the properties of the substance is called Thermodynamics Temperature Scale.

2.4.1 International Practical Temperature Scale (IPTS-68)

In view of the practical difficulties associated with the use of ideal gas thermometer, the International Committee on Weight and Measures has adopted the *International Practical Temperature Scale* in 1968 based on a number of easily reproducible fixed points. The standard fixed point which is easily reproducible by international agreement is triple point of water, the state of equilibrium between steam, ice and liquid water. The temperature at this point is defined as 273.16 Kelvin, abbreviated as 273.16 K, as a matter of convenience. This makes the temperature interval from ice point = 273.16 K to steam point = 373.16 K equal to 100 K, at pressure of 1atm.

2.4.2 Celsius or Centigrade Scale

The ice point of water on this scale is marked as zero, and the boiling point of water, at a pressure of 1atm is 100. The space between these two points has 100 equal divisions and each division represents one degree Celsius (written as °C); then

$$t = T - 273.16^{\circ} \text{C} \tag{2.4}$$

Thus, the Celsius temperature (t_s) at which steam condenses at a pressure of 1 atm

$$t_s = T_s - 273.16^{\circ}\text{C}$$

= 373.16°C - 273.16°C = 100.0°C

2.4.3 Electrical Resistance Thermometer

In the resistance thermometer the change in resistance of a metal wire due to its change in temperature is the thermodynamic property. The wire, frequently, may be incorporated in a Wheatstone Bridge circuit. In a restricted range, the following quadratic equation is often used

$$R_{t} = R_{0} \left(1 + at + bt^{2} + ct^{3} + \cdots \right)$$
(2.5)

Where R_0 is the resistance of the platinum wire when it is surrounded by melting ice, R_t is the resistance at temperature t and a and b are constants.

The triple point represents an equilibrium state between solid, liquid and vapour phases of substance. Normal boiling point is the temperature at which the substance boils at standard atmospheric pressure of 760 mm Hg. Normal freezing point is the solidification or melting point temperature of the substance at standard atmospheric pressure.

Sl.No.	Equilibrium State	Assigned value of temperature				
		T(K)	t(°C)			
1.	Triple point of hydrogen	13.81	-259.34			
2.	Boiling point of hydrogen at 33.306 kPa	17.042	-266.108			
3.	Normal boiling point of hydrogen	20.28	-252.87			
4.	Normal boiling point of neon	27.102	-246.048			
5.	Triple point of oxygen	54.361	-218.789			

Table: 2.2. Fixed Points of the IPTS - 68

6.	Normal boiling point of Oxygen	90.188	-182.962
7.	Triple point of Water	273.16	0.01
8.	Normal boiling point of Water	373.15	100.1
9.	Normal freezing point of Antimony (Antimony point)	630.74	357.59
10.	Normal freezing point of Zinc (Zinc point)	692.73	419.58
11.	Normal freezing point of Silver (Silver point)	1235.08	961.93
12.	Normal freezing point of Gold (Gold point)	1337.58	1064.43

The whole temperature scale may be divided into four ranges, based on the available method of measurement. The equation for interpolation for each range are as follows:

(i) *From* – 259.34°*C* (*Triple point of Hydrogen*) to 0°*C*: A platinum resistance thermometer of a standard design is used. A polynomial equation of the following form is deployed.

$$R_{t} = R_{0} \left(1 + at + bt^{2} + ct^{3} \right)$$
(2.6)

Here R_0 is the resistance at the ice point is fitted between resistance of the wire (R_1) and the temperature (t).

(ii) From $0^{\circ}C$ to 357.59°C (Antimony point): It is also based on platinum resistance thermometer. The diameter of the platinum wire must lie between 0.05 and 0.2 mm and governing equation is

$$R_{t} = R_{0} \left(1 + at + bt^{2} \right) \tag{2.7}$$

(iii) *From 357.59°C to 1064.43°C(Gold point)*: It is based on standard platinum versus platinum thermocouple. A three term equation, as following, is used.

$$E = a + bt + ct^2 \tag{2.8}$$

Here *E* is the EMF of the thermocouple.

(iv) *Above 1064.43* °C: The temperature is calculated from Planck's equation for black body radiation.

2.5 HEAT AND HEAT TRANSFER

The heat is defined as the form of energy which is transferred, without transfer of mass, across a boundary by virtue of a temperature difference between the system and the surroundings. It is a form of transit energy which can be identified only when it crosses the boundary of a system.

The temperature difference is the 'potential' or 'force' and heat transfer is the 'flux'. The heat can be transferred in three distinct ways, i.e., *conduction, convection* and *radiation*. The process of heat transfer between two bodies in direct contact is called conduction. The heat may be transferred between two bodies separated by empty space or gases by electromagnetic waves and the process is known as radiation. A third method of heat transfer is convection through fluid in motion.

The heat always flows from higher temperature to lower temperature without external energy. So, heat is transferred across a boundary from a system at a higher temperature to a system at lower temperature by virtue of the temperature difference. The heat is a form of transit energy, so it is not a thermodynamic property. It is not a point function, it is a path function. Conventionally when heat flowing into a system, it is considered positive and heat flowing out of a system, it is considered negative, shown in Fig.2.3.



2.6 SPECIFIC HEAT

The specific heat of a substance is defined as the amount of heat required to raise the temperature of unit mass of any substance by one degree. It is generally denoted by C and its unit in S.I is taken as kJ/kg·K. Heat required to raise the temperature of m kg mass of a substance from an initial temperature of T_1 to a final temperature of T_2 , then

Heat required
$$Q = mC (T_2 - T_1) \text{ kJ}$$
 (2.4)

where T_1 and T_2 may be either in 0°C or in K and C is specific heat in kJ/kg.K.

The solids and liquids do not change or change very negligibly in volume on heating, therefore they have only one specific heat. But gases have the following two specific heats depending upon the process adopted for heating the gas.

- (i) Specific heat at constant pressure C_p and
- (ii) Specific heat at constant volume C_{v}

It is noted that C_p is always greater than C_v . Relation between two specific heats is: $C_p - C_v = R$, where R is known as characteristic gas constant and its unit is J/kg K or kJ/kg.K. Value of gas constant (R) is different for different gases. For air it is taken as 287J/kg K or 0.278 kJ/kg.K. in S.I.

Ratio of C_p and C_v is, $\frac{C_p}{C_v} = \gamma$ where γ is adiabatic index. For air $\gamma = 1.4$ [For details see chapter

on properties of perfect gas]

2.7 THERMAL OR HEAT CAPACITY OF A SUBSTANCE

It is defined as the amount of heat required to raise the temperature of whole mass of a substance through one degree. Mathematically,

Thermal or heat capacity of a substance = mC kJ

Where m = mass of the substance in kg, and C = specific heat of the substance in kJ/kg.K.

2.8 WATER EQUIVALENT OF A SUBSTANCE

It may be defined as the quantity of water requires to raise unit temperature by the same quantity of heat which requires the substance to raise its temperature through one degree. Mathematically,

Water equivalent of a substance = mC kg

Where

m = mass of the substance in kg, and

C = specific heat of the substance in kJ/kg.K.

The thermal capacity and water equivalent of a substance are equal in numerical value, but in different units.

2.9 MECHANICAL EQUIVALENT OF HEAT

Heat and work are mutually convertible which was established by *Joule*. He established, experimentally, that certain amount of work is required to produce unit amount of heat. The ratio between work (W) and heat (H) i.e., is denoted by J (named after Joule) and is known as *Joule's equivalent or mechanical equivalent of heat*.

2.10 WORK

In mechanics, work is the effect of force. The work is said to be done by a force when the force acts upon a body and body moves in the direction of applied force. The magnitude of mechanical work is the product of the force and the distance moved parallel to the direction of applied force. Mathematically

$$W = \int_{1}^{2} dW = F dS$$

where F is applied force, and dS is the differential displacement parallel to direction of applied force.

The unit of work depends on the unit of force and unit of distance moved. In S.I, unit of force is newton and unit of distance is meter, so the practical unit work is N-m. It is known also as joule (briefly written as J). So, 1 N-m = 1 J.

In thermodynamics, work may be defined as the energy in transition between the system and surrounding. So, the work may be defined as follows:

(a) According to Obert: Work is defined as the energy transferred (without the transfer of mass) across the boundary of a system because of an intensive property difference other than temperature that exists between the system and surrounding.

The pressure difference, i.e., the intensive property difference (between the system and surrounding) at the surface of the system gives rise to a force and the action of this force over a distance is called *mechanical work*.

Electrical work is also the same case. In this case the intensive property difference is the electrical potential difference between the system and surrounding. And the resulting energy transfer across the system and boundary is known as *electrical work*.

(b) According to Keenan: Work is said to be done by a system during a given operation if the sole effect of the system on things external to the system (surroundings) can be reduced to the raising of a weight.

The weight may not be raised actually but the net effect external to the system should be the raising of a weight.



To give example of this, consider a system consisting of a battery as shown in Fig.2.4(*a*). The terminals connected to a resistance through a switch is on for a certain period of time, then the current will flow through the battery and resistance. As a result the resistance will be warmer. This clearly shows that the system (battery) has interaction with the surroundings. In other words, energy transfer(electrical energy) has taken place between the system and the surroundings because of potential difference(not the temperature). According to the definition of work, by laws of mechanics, if there is no force which moves through a distance, then no work is done by the system. However, according to the thermodynamic definition, the work is done by system because the resistance can be replaced by a ideal motor(100% efficient) driving a winding drum, thereby raising an weight, as shown in Fig. 2.4(*b*). Thus, the sole effect external to the system (surrounding) has been reduced to the raising of an weight. Hence, thermodynamic work is done by the system.



Fig. 2.4 (b)

Also, it is relevant to mention that work is manifested at the system boundary only during any interaction between the system and surroundings. Before interaction no work is present. Only work transfer takes place.

2.11 SIGN CONVENTION OF WORK

In thermodynamics, by convention, work done by the system is taken to be positive i.e., when work leaves from the system. The work done on the system is taken to be negative i.e., if work enters into the system from surrounding, as shown in Fig.2.5. In any process when the system does positive work, its surroundings do an equal amount of negative work. Thus, in any process algebraic sum of work done by the system and its surrounding is zero. Mathematically,

$$W_{\text{system}} + (-W_{\text{surrounding}}) = 0$$
(2.9)

System Boundary

 $W (\text{out})$

 $(+ \text{ ve})$

 $W (\text{in})$

 $(- \text{ve})$

Fig. 2.5

2.12 WORK DONE DURING A QUASI-STATIC OR QUASI-EQUILIBRIUM PROCESS

From the concept of mechanics, the differential work done (dW) is given by

$$dW = F \cdot dS \tag{2.10}$$

where F is the applied force and dS is the differential displacement.

Consider a certain amount of gas contained in a cylinder-piston assembly as shown in the Fig.2.6(*a*) in a schematic representation of work by a gas. The system, enclosed by the dotted line, constitute the gas contained in the cylinder and Fig. 2.6(*b*) is representation of work done on a $p-\forall$

diagram. The shaded area represents the work done $\int_{1}^{2} pd \forall$ by the gas. In a differential time dt, let the piston of cross-sectional area A moves a differential distance dL while the motion of piston is opposed by an external pressure (p - dp) where $dp \rightarrow 0$. Then, the differential work done by the gas is given by

$$dW = F \cdot dS$$

= (pA) dL [:: F = p × A]
= p \cdot AdL



It may be noted that the external pressure (p - dp) is always infinitesimally smaller than the gas pressure inside the system. Hence, the expansion process can be reversed at any time by increasing the external pressure infinitesimally more than the gas pressure. That is, if the external pressure is (p + dp), the gas undergoes a compression process.



5 . ()

During the process, the expansion or compression of the gas creating a force on the moving boundary are balanced and hence the process is reversible. Being a reversible one, it must be a

Quasi-Static or Quasi-Equilibrium Process, because 'a reversible process is a quasi-static process, but a quasi-static process need not be a reversible process'.

2.13 WORK AND HEAT TRANSFER — A PATH FUNCTION

Consider that a system, from an initial equilibrium state 1 reaches to a final equilibrium state 2 by two reversible process 1-A-2 and 1-B-2 as shown in Fig.2.7(*a*). When the system changes from its initial state 1 to final state 2, the quantity of work or heat transfer will depend upon the intermediate stages through which the system passes, i.e., its path.



Figure 2.7(*a*) shows two different reversible paths 1-*A*-2 and 1-*B*-2 connecting the given initial and final states of system. Since $\int_{1}^{2} p \cdot d \forall$ represents the area under the curve, the work done during path 1-*A*-2 is different from the work done during the path 1-*B*-2. In other words, the work done depends on the path, followed in going from one state to another i.e., by a system.

Hence, it is not possible to evaluate the work done purely from a knowledge of its initial and final states of a system. The work done by a system is a path function and hence it is not a property of a system. Work is an interaction, that is energy transfer, between a system and its surroundings. Hence work is energy in transit. Mathematically work is an inexact differential and written as δW . On integrating for path 1-*A*-2,

$$\int_{1,A}^{2} \delta W = \left[W\right]_{1,A}^{2} = \left(W_{1-2}\right)_{A}$$
(2.12)

It should be noted that in the case of inexact differentials such as work, $\int_{1}^{2} \delta W \neq W_2 - W_1$

because work is not a point function. We never speak of work W_1 in the system in state 1 or work W_2 in state 2. In contrast to this, there is only work, either W_{in} or W_{out} in transition, which manifests itself

at the system boundary during a process, and can only be found by integration of $(pd\forall)$ along the path followed in the quasi-equilibrium process.



Like work, heat transfer during from its initial state point 1 to final state point 2 will depend on the intermediate state through which the system passes, i.e., its path. In other words, heat is a path function and is an inexact differential and written as δQ . On intergrating for the path 1-A-2 in Fig. 2.7(b)

$$\int_{1,A}^{2} \delta Q = [Q]_{1,A}^{2} = (Q_{1-2})_{A}^{2} = ({}_{1}Q_{2})_{A}^{2}$$
(2.13)

It may be noted that in the case of inexact differentials such as heat, $\int_{1,A}^{2} \delta Q \neq Q_2 - Q_1$, because heat is not a point function and it would be more appropriate to write $\int_{1,A}^{2} \delta Q = (Q_{1-2})_A$. It can be

written as the integral of the product of intensive property (T) and differential change of an extensive

property, say entropy (S), i.e. $\int_{-\infty}^{2} T \delta S$. The *T*-S diagram is shown in the figure.

2.14 COMPARISON OF HEAT AND WORK

Similarities between heat and work are follows :

- (i) Heat and Work are neither thermodynamic properties, nor point functions. The system itself does not posses heat and work. When a system undergoes a change, heat transfer or work done may occur. They are both transient phenomena.
- (ii) The heat and work are boundary phenomena. They are observed at the boundary of the system.

- (iii) Heat and work are both transit energy crossing the boundary of the system.
- (iv) The heat and work both are path functions and hence are inexact differentials. They are written as δQ and δW .
- (v) The area under the pressure-volume (i.e., p-∀) diagram represents the work done during the process and is given by ∫ pd∀. Similarly the area under the temperature-entropy i.e., (T-S) diagram represents the heat transfer during the process and is given by ∫ TdS.
- (vi) Work is written by the integral of the product of the intensive property(p) and differential change of the extensive property (\forall) .

Heat transfer can also be written as the integral of product of the intensive property (T) and the differential change of an extensive property, entropy(S).

2.15 EXAMPLE OF WORK

2.15.1 Paddle Wheel Work

This is also known as Stirring Work where $\int p d \forall = 0$, but work is done.

The paddle wheel work is an illustration of shaft-work. Paddle wheel work process is a process involving friction in which the volume of the system does not change at all, and still work is done on the system.

Representation of the process is provided by a system in which a paddle wheel turns a fixed mass of fluid as shown in Fig.2.8. Consider that in the system weight is lowered, paddle wheel runs. The work is transferred across the system boundary in the fluid system. The volume of the system remains constant and the work, $\int p d \nabla = 0$.



Fig. 2.8

If *m* is the mass of the weight lowered through a distance dz and *T* is the torque transmitted by the shaft in rotating through an angle $d\theta$, the differential work transfer to the fluid is given by

$$\delta W = mgdz = Td\theta$$

Work transfer
$$(W) = \int_{1}^{2} \delta W$$

= $\int_{1}^{2} mg \, dz = \int_{1}^{2} T \cdot d\Theta$ (2.14)

Thus, $\int pd \forall$ does not represent work for this case, although work has been done on the system. So work may be done on a closed system even though there is no volume change.

2.15.2 Extension of Solid Rod

Consider an elastic rod of length L and cross-sectional area A, which is fixed at one end as shown in Fig.2.9 as the system. The solid rod is fixed at x = 0, and a tensile force equal to F is applied at the other end of the rod.



Thus, $F = \sigma A$ where s is normal stress acting at the end of the rod. The rod elongates a distance dx due to elastic property. The work done due to elongation dx is given by

$$\delta W = -\sigma A dx$$

The negative sign shows that work is done on the rod when dx is positive. The work done for change of length from x_1 to x_2 is given by,

$$W = -\int_{x_1}^{x_2} (\sigma A dx)$$
(2.15)

2.15.3 Stretching of a Liquid Film (Surface Tension)

Suppose a liquid film is held between a wire from one side of which can be moved with the help of a sliding wire as shown in Fig.2.10 as a system. The interfacial tension or surface tension which is the force per unit length normal to a line on the surface is given by τ . Thus, force is given by, $F = 2L\tau$; when the factor 2 indicates two film surfaces act at the wire length *L*.



The work for displacement dx is given by,

$$\delta W = -2L\tau \cdot dx$$

Negative sign shows that work is done on the system when dx is positive. The surface in contact with wire is dA = 2Ldx

Thus,

...

$$\delta W = -\tau dA$$

$$W = -\int_{A_1}^{A_2} \delta W$$

$$= -\int_{A_1}^{A_2} \tau dA = -\tau (A_2 - A_1) \qquad (2.16)$$

2.15.4 Free Expansion Process

The free expansion, or unresisted expansion, process is an irreversible non-flow adiabatic process in which the volume of a closed system increases, and still no work at all is done. So, here $\int pd \forall$ is finite but work done is zero. Representation of this unresisted expansion process is shown in Fig. 2.11(*a*) & (*b*). A free expansion occurs when a fluid is allowed to expand suddenly into a vacuum chamber through an orifice of large dimensions.



Fig. 2.11(a)

Consider two chambers A and B separated by a membrane. The chamber A contains fluid having volume \forall_1 pressure p_1 and temperature T_1 and the chamber B is completely evacuated, pressure i.e., $p_{\text{ext}} = 0$. The fluid is in state 1 in chamber A. Let the membrane gets ruptured. The fluid then fills both the chambers and reaches state 2. Initial pressure p_1 of the fluid dropped to p_2 in the final state and volume \forall_1 reaches to volume \forall_2 in the final state. Both the chambers are completely insulated so that heat transfer is zero.

Let us first consider the fluid and vacuum together as the system shown in Fig.2.11(*a*), so as no work crosses the system boundary. Next we consider only the fluid as the system as in the figure 2.11(*b*). We observe that the system boundary moves and volume of the system change from \forall_1 to \forall_2 . But it is not quasi-equilibrium process and $p_{ext} = 0$. Hence, $\int p_{ext} d\forall$ is also zero and therefore, no work is done in the process. Free expansion process is thus an example of an expansion process in which $\int p d\forall$ is finite, but still $W_{1,2} = 0$. Hence, it is adiabatic process where $Q_{1,2} = 0$ and $U_{1,2} = 0$.



2.15.5 Shaft Work

When a shaft within a system is rotated, by a motor, there is work transfer into the system as shown in the Fig.2.8. This is so since the shaft can rotate a drum which in turn can raise on weight. If *T* is the torque applied to shaft running at *N* r.p.m, (angular velocity ω) then the rate of work done would be

$$W = T \cdot \omega = \frac{2\pi NT}{60} \tag{2.17}$$

Paddle wheel work is an example of negative shaft-work.

2.15.6 Electrical Work

Electrical work done by the flow of a current through electrolytic cell or a conductor or a resistor, etc., is given by

$$_1W_2 = -\int_1^2 \mathbf{E} \cdot \mathbf{I} \cdot \mathbf{dt}$$

where E is the potential difference, I is the current and dt is the time change.

So, power can be written by

$$P = -\frac{\delta W}{\delta t} = -E \cdot I \tag{2.18}$$

2.15.7 Flow Work

The flow work is the work required to move a fluid into or out of a system. To calculate the work required, consider an element of fluid of mass dm and volume $d\forall$. It is required to push this fluid into the control volume of an open system through a passage against an existing pressure p as shown in Fig.2.12 (*a*).Cross-sectional area of the passage is A. The fluid behind this element of fluid acts as a piston, and pushes the fluid into the control volume. Consider the imaginary piston is placed behind this elemental fluid, as shown in figure 2.12 (*b*). The distance moved by piston is dx, when the fluid with volume $d\forall$ has been pushed into the control volume.



Fig. 2.12 (a & b)

Force $F = p \cdot A$ and $dx = \frac{d \forall}{A}$

Now the required flow work is

 $\delta W = \text{ force } \times \text{ distance moved}$ $= F \times dx$

$$= p \cdot A\left(\frac{d\forall}{A}\right) = p \cdot d\forall$$

Work done per unit mass of the fluid flow is

$$W_{\rm flow} = \frac{\delta W}{dm}$$

$$= \frac{p \cdot d \forall}{dm}$$
$$= p \cdot \left(\frac{d \forall}{dm}\right) = p \cdot v_s$$
(2.19)

So the magnitude of the flow work per unit mass of fluid flow is simply the product of the two properties of the fluid, i.e., pressure and specific volume.

2.15.8 Non-flow Work

Consider a closed system process, say a cylinder and piston arrangement as shown in the Fig.2.6(b) from 1 to 2. The non-flow work done by the system in the process per unit mass of the substance is

$$W_{\rm nonflow} = \int_{1}^{2} p dv_{s} \tag{2.20}$$

and represented by the area 1-2-*b*-a-1, this is pdv_s work.

2.16 WORK IN NON-FLOW PROCESS VERSUS FLOW PROCESS

A process undergone by a closed system of fixed mass is referred to as a non-flow process. A change of state of a substance taking place while flowing through a control volume in an open system is referred to as a flow process.

The work done in a non-flow process in closed system is called non-flow work. The work done in a flow process in an open system is equal to the non-flow work plus the flow work.

i.e., $W_{\text{flow}} = \int p dv_s + p \cdot v_s$

Multiple Choice Questions

1. Temperature is a (b) non property (a) property 2. Thermometers used in thermodynamic applications are constant volume gas thermometer (a) (b) constant pressure gas thermometer (c) pyrometer all of these (d) 3. When heat flows into a system, conventionally it is considered negative (b) positive (a) The unit of specific heat of any substance is 4. J/kg (a) (b) Jkg (c) kg/J (d) J/kgK 5. The ratio of C_p and C_y of any substance is called (a) adiabatic constant (b) isothermal index (d) isobaric constant (c) pressure index

6.	Heat	and v	vork	are														
	(a)	exac	et inte	egrals	5				(b)	inexa	act ii	ntegral	s					
7.	Tem	perati	ire is	an														
	(a)	exac	et inte	egral					(b)	inexa	act ii	ntegral						
8.	The	amou	nt of	work	c don	e in fr	ee ex	pansio	n pro	cess is								
	(a)	non	zero						(b)	zero								
	(c)	unit	у						(d)	none	of	these						
9.	In co	onstan	t vol	ume g	gas th	ermo	meter	tempe	eratur	e is me	asu	red in t	erms o	of				
	(a)	heig	ht of	merc	ury c	olum	1		(b)	diam	eter	of mer	cury c	olumn				
	(c)	volu	ime o	of mer	cury	colum	n		(b)	none	of	these						
10.	The	S.I. u	nit fo	or wat	ter eq	uivale	ent of	a subs	tance	e is	7							
	(a)	kg V							(b) (d)	J/kgł	S.	thasa						
11	(C) The	к ст.,,	nit fo	r haa	taan	ooitu	ofoo	ubstan	(u)	none	01	ulese						
11.	(a)	5.1. u ka	IIIt IC	n nea	u cap	acity		uostan	(b)	I/kal	7							
	(a) (c)	K							(d)	J								
12.	The	gener	allv a	accep	ted v	alues	of ad	iabatic	cons	stant fo	r air	is						
	(a)	1.04		····P					(b)	0.104								
	(c)	10.4							(d)	1.4								
13.	The	S.I. u	nit fo	or adia	abatic	c cons	tant i	S										
	(a)	J/kg							(b)	kg/J								
	(c)	J/kg	K						(d)	unitl	ess,	pure n	umbei	ſ				
14.	The	speci	fic he	eat fo	r any	gased	ous su	ıbstanc	e in o	constar	nt vo	olume p	proces	s is ger	nerally	^y symł	oolise	ed by
	(a)	C_{v}							(b)	C_{p}		41						
17	(C)	<i>C</i> _{<i>r</i>} .	<u>ر</u> ۱					1 /	(a)	none	01	these			11	1	1.	11
15.	i ne	speci	nc ne		r any	gaseo	ous su	ostanc	$e \ln c$	constan	t pre	essure p	proces	s is gei	herally	/ symt	oonse	ed by
	(a)	C_{v}							(d)	none	of	these						
	(-)	r							()									
								Ar	ISWE	ers								
1. (a)	2.	(d)	3.	(b)	4.	(d)	5.	(a)	6.	(b)	7.	(a)	8.	(b)	9.	(a)	10.	(a)
11. (d)	12.	(d)	13.	(d)	14.	(a)	15.	(b)										
						NU	JME	RICA	AL E	EXAN	IPL	.ES						

EXAMPLE 1

Over a specified temperature range, the electrical resistance thermometer holds the relationship as: $R_t = R_0 [1 + \alpha (t - t_0)]$, where R_0 is resistance in ohm, measured at reference temperature t_0 in 0°C and α is material constant with unit of (°C)⁻¹. The test result are: $R_t = 51.39$ ohm at t = 0°C, $R_t = 51.72$ ohm at t = 91°C. Determine the reading of the resistance for 70°C, 50°C and 100°C on this thermometer.

SOLUTION		
Given	$t_0 = 0^{\circ}\text{C}, t = 0^{\circ}\text{C}, R_t = 51.39$	
We have	$R_t = R_0 [1 + \alpha (t - t_0)]$	
So,	$51.39 = R_0 (1+0)$	
··	$R_0 = 51.39$	
Again,	$t_0 = 0^{\circ}$ C, $t = 91.0^{\circ}$ C, $R_t = 51.72$	
<i>.</i> .	$R_t = R_0 [1 + \alpha (t - t_0)]$	
	$51.72 = 51.39 [1 + \alpha (91 - 0)]$	
	$= 51.39 + \alpha \times 51.39 \times 91$	
Therefore,	$\alpha = \frac{51.72 - 51.39}{51.39 \times 91} = 7.05657 \times 10^{-5/\circ} \text{C}$	
Thus,	$R_{70^{\circ}} = R_0 [1 + \alpha (t - t_0)]$	
	$= 51.39 [1 + 7.05657 \times 10^{-5} \times 70] = 51.644 \text{ ohm}$	
	$R_{50^{\circ}} = 51.39 [1 + 7.05657 \times 10^{-5} \times 50] = 51.57 \text{ ohm}$	
	$R_{100^{\circ}} = 51.39 [1 + 7.05657 \times 10^{-5} \times 100] = 51.75 \text{ ohm}$	
EXAMPLE 2		

The resistance of a platinum wire is found to be 11 ohms at the ice point, 15.247 ohms at steam point, and 28.887 ohms at the Zinc point. Find the constant *a* and *b* in the equation. $R_t = R_0 (1 + at + bt^2)$.

SOLUTION

Here,	$R_0 = 11$ ohm, $R_{\text{steam}} = 15.247$ ohm, $R_z = 28.887$ ohm and Zinc point = 419.58°	°C
Following the relation,	$R_t = R_0 \left(1 + at + bt^2 \right)$	
	$15.247 = 11 (1 + a \times 0.01 + b \times 0.01^2)$	
or,	1.3861 = 1 + 0.01a + 0.0001b	
or,	0.3861 = 0.01a + 0.0001b	
or,	38.61 = a + 0.01b	(i)
And	$28.887 = 11 \left[1 + a \times 419.58 + b \left(419.58 \right)^2 \right]$	
or,	2.6261 = 1 + 419.58 a + 176047.3764 b	
or,	1.6261 = 419.58 a + 176047.3764 b	(ii)
Solving equations (i) and (i	i), we obtain the constants as a = 38.6109, b = -0.092	

EXAMPLE 3

The temperature *t* on a thermometric scale is defined in terms of a property *K* by the relation $t = a \ln K + b$ where *a* and *b* are constants. The values of *K* are found to be 1.83 and 6.78 at the ice point and the steam point, the temperatures of which are assigned the numbers 0 and 100 respectively. Determine the temperature corresponding to a reading of *K* equal to 2.42 on the thermometer.

Solution		
From the relation	$t = a \ln K + b$, we have	
	$0 = a \ln 1.83 + b$	(i)
and	$100 = a \ln 6.78 + b$	(ii)
Solving above equations,	a = 78.68951, and $b = -47.5533282$	
Hence, relation is	$t = 78.68951 \ln K - 47.5533282$	
For $K = 42$, we find	$t = 78.68951 \ln 2.42 - 47.5533282$	
	= 21.989 °C	

EXAMPLE 4

A constant volume gas thermometer containing nitrogen is brought into contact with a system of unknown temperature and then into contact with a system maintained at the triple point of water. The mercury column attached to the device has reading of 59.2 and 2.28 cm respectively for the two systems. If the barometric pressure is 960 m bar (96.0 kPa), what is the unknown temperature in Kelvin, if g = 9.806 m/ sec². Specific gravity of mercury may be taken as 13.6.

SOLUTION

We know

$$T = 273.16 \frac{p}{p_{tp}}$$

where T is the unknown temperature, p is the gauge pressure at unknown temperature, p_{p} is the gauge pressure at the triple point of water.

So, p_{gauge} at unknown temperature

$$= \left(13.6 \times 10^{3} \times 9.807 \times \frac{59.2}{100}\right) \text{ N/m}^{2}$$

= 78.95 kPa = 789.5 m bar
∴ $p_{abs} = (789.5 + 960) \text{ m bar}$
= 1749.5 m bar = 174.95 kPa = p

Now, p_{gauge} at triple point temperature

$$= \left(13.6 \times 10^{3} \times 9.807 \times \frac{2.28}{10^{3}}\right) \text{ N/m}^{2}$$

= 3.0406 kPa = 30.406 m bar
= (30.406 + 960) m bar = 990.406 m bar
= 99.0406 kPa = p_{ip}

Hence, the unknown temperature

$$T = 273.16 \left(\frac{174.95}{99.0406}\right) = 482.52 \text{ K}$$

(iii)

EXAMPLE 5

A constant volume gas thermometer immersed in boiling water of $t_1 = 100^{\circ}$ C reads $z_1 = 30.0$ cm of Hg. Later the thermometer is immersed in a radiator fluid at temperature of 150°C. That time, barometric pressure is 76.0 cm Hg. Determine the height z_2 of the mercury column, in cm. Assume the gas in the thermometer as perfect gas.

SOLUTION

For perfect gas we can write

$$\frac{p_1 \forall_1}{T_1} = \frac{p_2 \forall_2}{T_2} = \text{constant}$$
(i)

At constant volume above relation reduce to

$$\frac{p_1}{T_1} = \frac{p_2}{T_2}$$
 (ii)

So, Here, $\frac{z_0 + z_2}{t_2} = \frac{z_0 + z_1}{t_1}$ $T_1 = 100 + 273 = 373 \text{ K}, \qquad z_1 = 30 \text{ cm}$ $T_2 = 150 + 273 = 423 \text{ K}, \qquad z_0 = 76 \text{ cm}$

Applying equation (iii), we get $z_2 = 44.21 \text{ cm}$

EXAMPLE 6

A soap film is suspended on a 5 cm \times 5 cm wire frame ABCD as shown in figure. The movable wire BC is displaced by 1 cm to B'C' by an applied force while the surface tension of soap film remains constant at 25×10^{-5} N/cm². Estimate the work done in stretching the film in N·m.

SOLUTION

The film is a closed system with moving boundary. Work done, on both side of the film.



EXAMPLE 7

A gas in the cylinder and piston arrangement comprises the system is shown in Fig.2.6(a). It expands from 1.5 m³ to 2 m³ while receiving 200 kJ of work from a paddle wheel. The pressure on the gas remains constant at 600 kPa. Determine the net work done by the system.

SOLUTION

The work done by the piston is given by $\int p d \forall$ and this is the work done by the system

$$W_{\text{gas}} = \int_{1.5}^{2} p d \forall$$
$$= p \int_{1.5}^{2} d \forall$$
$$= p [\forall_2 - \forall_1]_{1.5}^{2}$$
$$= 6 \times 10^5 (2 - 1.5) \text{ N·m}$$
$$= 3 \times 10^5 \text{ J}$$

[pressure is constant]

The gas also receives 200 kJ from the paddle wheel and this is work done on the system.

Thus, the net work done by the system is,

$$W_{\text{net}} = 3 \times 10^5 - 200 \times 10^3$$

= 1 × 10⁵ N·m = 100 kJ

EXAMPLE 8

A gas is at a pressure of 3 bar in cylinder attached with frictionless movable piston shown in Fig.2.6(*a*). The spring force exerted through the piston is proportional to the volume of gas. Also additional atmospheric pressure of 1 bar acts on the spring side of piston. Determine the work done by gas in expansion from 0.1 m^3 to 0.5 m^3 .

or J

SOLUTION

The pressure exerted by the spring

$$p_s = \text{gas pressure} - \text{atmospheric pressure}$$

= 3 - 1 = 2 bar = 2 × 10⁵ N/m²

Thus, the spring force is given by

$$F_s = p_s \times A = (2 \times 10^5 \times A) \,\mathrm{N}$$

where A is cross-sectional area of piston in m². This force gives spring displacement of $x_1 = \frac{\text{volume}}{\text{area}} = \frac{0.1}{A}$

Thus, spring constant is given by $k_s = \frac{2 \times 10^5 \times A}{\left(\frac{0.1}{A}\right)} = 20 \times 10^5 A^2 \text{ N/m}$

Similarly we can write

$$x_2 = \frac{0.5}{A}$$

The work done against the spring

$$W_{S} = \int_{x_{1}}^{x_{2}} 20 \times 10^{5} \times A^{2} x \, dx$$
$$= 20 \times 10^{5} A^{2} \left[\frac{x^{2}}{2} \right]_{x_{1}}^{x_{2}}$$

$$= 10 \times 10^{5} A^{2} \qquad [x_{2}^{2} - x_{1}^{2}]$$

$$= 10 \times 10^{5} A^{2} \qquad \left[\frac{0.5^{2}}{A^{2}} - \frac{0.1}{A^{2}} \right]$$

$$= 2.4 \times 10^{5} \text{ N·m.}$$
The atmosphere $W_{\text{atm}} = 1 \times 10^{5} (0.5 - 0.1) = 0.4 \times 10^{5} \text{ N·m}$
The area by the area $W = W + W$

Work done against Thus, total work done by the gas $W = W_{S} + W_{atm}$

= $2.4 \times 10^5 + 0.4 \times 10^5 = 2.8 \times 10^5$ N·m or J.

EXAMPLE 9

For a system whose mass is 4.5 kg undergoes a process and the temperature changes from 50°C to 100°C. Assume that the specific heat of the system is a function of temperature only. Calculate the heat transfer

during the process for the following relationship $C_n = \left[0.3 + \frac{20}{T + 150}\right] \text{ kJ/ kg K.}$

373

SOLUTION

Here

and C_n

$$T_1 = 273 + 50 = 323 \text{ K}, T_2 = 273 + 100 = 373 \text{ K}, m = 4.5 \text{ kg}.$$

= $0.3 + \left(\frac{20}{T + 150}\right)$

.:.

$$\begin{aligned} Q_{1-2} &= \int_{323} mC_n dT \\ &= 4.5 \int_{323}^{373} \left[0.3 + \frac{20}{T + 150} \right] dT \\ &= 4.5 \left[0.3 \int_{323}^{373} dT + \int_{323}^{373} \frac{dT}{T/20 + 150/20} \right] \\ &= \left\{ \left[0.3T \right]_{323}^{373} + \left[\frac{1}{0.05} \ln \left(0.05T + 75 \right) \right]_{323}^{373} \right\} \\ &= 4.5 \left\{ 0.3 \times 50 + \frac{1}{0.05} \left[\ln \left(18.65 + 75 \right) - \ln \left(16.15 + 75 \right) \right] \right\} \\ &= 4.5 \left\{ 15 + \frac{1}{0.05} \ln \frac{93.65}{91.15} \right\} \\ &= 69.935 \, \text{kJ} \end{aligned}$$

EXERCISE

- 1. Define temperature. Name the different temperature scale in common use. Establish relation between Celsius and Fahrenheit scale?
- 2. What is absolute temperature ?
- 3. What do you understand by N.T.P and S.T.P? What are their values?
- 4. Distinguish between gauge pressure, absolute pressure and atmospheric pressure. What are their relation.
- 5. Define energy. What is stored energy and transit energy? Discuss the types of stored energy.
- 6. Define thermodynamic work. Show that both heat and work are a path function not a point function.
- 7. Work done is given by $W = \int p d \forall$, Is there be any situation where $d \forall = 0$ and yet $dW \neq 0$?
- 8. Which of the following characteristics are true for work? (a) It is a path functions (b) It is not a property of a system. (c) Its differential is not exact. (d) It is not stored in a system (e) It is energy in transit. (f) It can be indentified only during the course of interaction.
- 9. Which of the characteristics given in question no. 8 are true for heat.
- 10. State the Zeorth Law of Thermodynamics. How is mercury in thermometer able to find the temperature of a body using the Zeroth Law of Thermodynamics.
- 11. Power transmission through a rotating shaft is a common engineering practice. Derive a relation to determine the power transmitted through a rotating shaft.

Ans. $[P = T\omega]$

12. An aqueous soap solution is maintained in the form of thin film, on a wire frame as shown in Fig 2.10. If the surface tension of the liquid is 72.75 mN/m. Calculate the work done when the movable wire is displaced by 1 cm in the direction indicated.

Ans. $[W = -29.1 \times 10^{-6} \text{ J}]$

13. One end of steel rod of 1 cm in diameter and 50 cm in length is fixed and a force is continuously applied at the other end in a direction which is parallel to the length of the rod. Calculate the work done by the force on the steel rod to change its length by 1 cm. Also calculate the stress and force applied on the rod. The Young's modulus of steel is 2×10^{11} N/m².

Ans.
$$[W = -1.5708 \text{ kJ}, F = 3.1416 \times 10^5 \text{ N}]$$

14. A piston has area of 5 cm². What mass must the piston have if it exerts a pressure of 50 kPa above atmospheric pressure on the gas enclosed in the cylinder.

Ans. [7.71 kg]

- 15. A pressure guage reads 2.4 bar and the barometer reads 75 cm of Hg. Calculate the absolute pressure in bar and in the standard atmosphere.
- 16. A manometer has a liquid of density 800 kg/m³, the difference in level of the two legs is 300 mm. Determine the presure difference. Read it in kgf/m², bar and kPa.

This page intentionally left blank



PROPERTIES AND THERMODYNAMIC PROCESSES OF GAS

3.1 INTRODUCTION

Vapour is partially evaporated liquid which are in contact with liquid surface and contains some *particles* of liquid in suspension. It does not obey the gas laws. Wet and dry steams are the examples of vapour.

Gas is completely evaporated liquid which are not in contact with liquid surface and does not contain any particles of liquid in suspension, and it obeys the gas laws. Oxygen, hydrogen and superheated steam etc., are the examples of gas.

An *ideal* or perfect gas may be defined as a state of a substance, which *strictly obeys* all the gas laws under all conditions of temperature and pressure. In actual practice, there is no real or actual gas which strictly obeys Boyle's Law, Charle's Law and Gay-Lussac's Law over the entire range of temperature and pressure. But, the real gases which are ordinarily difficult to liquefy (oxygen, hydrogen and air), within certain temperature and pressure limits may be regarded as perfect gases.

3.2 GENERAL GAS EQUATION

The gas laws give us the relation between the two variables when the third variable is constant. But in actual practice, all the three variables i.e., pressure, volume and temperature change simultaneously. In order to deal with all practical cases, the Boyle's Law and Charle's Law are combined together to form a general gas equation.

According to Boyle's Law:

$$p \propto \frac{1}{\forall}$$
 or $\forall \propto \frac{1}{p}$ [T = constant]

According to Charle's Law:

 $p \forall = KT$, or $\frac{p \forall}{T} = K$

 $\forall \propto T$ Combining these two relations, we have

$$\forall \propto T$$

Ì

[p = constant]

(3.1)

or

where, K = constant of proportionality and its value depends on mass and properties of the gas. So, the general equation of gas can be

written as,
$$\frac{p_1 \forall_1}{T_1} = \frac{p_2 \forall_2}{T_2} = \text{constant}$$
 (3.2)

3.3 EQUATION OF STATE AND CHARACTERISTIC EQUATION OF GAS

It is modified form of equation of gas. Let us consider \forall volume of gas of mass *m* kg and *n* number of moles having molecular weight *M*. The volume of 1 kg mole gas is v_m and the volume of 1 kg gas is known as *specific* volume (v_s) .

So,
$$v_s = \frac{\forall}{m}$$
, $v_m = \frac{\forall}{n}$ and $M = \frac{m}{n}$. Now the equation (3.1) is known as equation of state.

When we consider volume of 1 kg mass, the equation (3.1) will be

$$pv_s = RT$$
(3.3)

where *R* is known as characteristic gas constant or simply the gas constant. It is different for different gases. In S.I., R = 287 J/kgK or 0.287 kJ/kgK.

Again,

 \Rightarrow

$$p\frac{\forall}{m} = RT$$

$$p\forall = mRT$$
(3.4)

The equation (3.4) is known as characteristic equation of gas. Unit of R may be obtained as

$$R = \frac{p\forall}{mT} = \frac{\left(\text{N/m}^2\right) \times \text{m}^3}{\text{kg} \cdot \text{K}} = \text{N} \cdot \text{m/kg} \cdot \text{K} = \text{J/kg} \cdot \text{K}$$

3.3.1 Physical Significance of R

We know characteristic gas constant is $R = \frac{p \forall}{mT}$. So, *R* is defined as the amount of work per kg of mass of gas per unit change of absolute temperature.

3.4 UNIVERSAL GAS CONSTANT OR MOLAR GAS CONSTANT

We know the equation of state is $p \forall = KT$. When we consider volume of 1 kg mole gas (v_m) , the equation of state will be

$$pv_m = R_m T \tag{3.5}$$

Here, K is replaced by R_m , which is called *universal gas constant* or *molar gas constant*. It is constant for all gases.

But,

$$p\frac{\forall}{n} = R_m T \tag{3.6}$$

$$\Rightarrow \qquad p \forall = nR_m T = \frac{m}{M}R_m T = m\frac{R_m}{M}T$$

Comparing eqns. (3.4) and (3.6),

This is the relation between universal or *molar gas constant* (R_m) and *characteristic gas constant*. So, the universal gas constant of a gas is the product of characteristic gas constant and the molecular mass of the gas. In S.I., $R_m = 8314$ J/kg mole·K = 8.314 kJ/kgmole·K

3.5 SPECIFIC HEAT OF GAS

The specific heat of a substance may broadly be defined as the *amount of heat required to raise* the temperature of its unit mass through one degree. All the liquids and solids have only one specific heat. But a gas can have any number of specific heat, lying between zero to infinity, depending upon the condition, under which it is heated. Only two types of specific heat of gas, specific heat at constant volume (C_v) and specific heat at constant pressure (C_p) are important from the subject point of view.

3.5.1 Specific Heat at Constant Volume (C)

It is the amount of heat required to raise the temperature of unit mass of gas through one degree when it is heated at a constant volume. It is generally denoted by C_v . Whenever, a gas is heated at constant volume, total heat energy supplied is utilized to increase the temperature, because no work is done by the gas,

$$W_{1-2} = \int_{1}^{2} p d \forall = p (\forall_2 - \forall_1) = 0$$
(3.8)

Hence, total heat (Q_{1-2}) supplied to the gas at constant volume equals with change of internal energy of the gas (dU).

$$(Q_{1-2})_{\nu} = mC_{\nu} (T_2 - T_1) = dU = U_2 - U_1$$
(3.9)

So,

$$C_{v} = \frac{(Q_{1-2})_{v}}{m(T_{2} - T_{1})}$$
(3.10)

3.5.2 Specific Heat at Constant Pressure (C_p)

It is the amount of heat required to raise the temperature of the unit mass of a gas through one degree, when it is heated at constant pressure. It is generally denoted by (C_p) . Total heat supplied to the gas at constant pressure,

$$(Q_{1-2})_p = mC_p (T_2 - T_1)$$
(3.11)

Here, total heat supplied, when gas is heated at constant pressure is utilized

(a) to raise the temperature of gas which increases internal energy of the gas,

$$dU = mC_{v} (T_2 - T_1)$$
(3.12)

(b) to do some external work during expansion $\delta W = pd \forall$ which on integration,

$$\int_{1}^{2} \delta W = \int_{1}^{2} p d \forall$$

 \Rightarrow

 \Rightarrow

$$W_{1-2} = p(\forall_2 - \forall_1) = mR(T_2 - T_1)$$
(3.13)

Hence, from above we can write,

$$\delta Q = dU + \delta W \tag{3.14}$$

Integrating within limits, we obtain

$$(Q_{1-2})_p = mC_v (T_2 - T_1) + mR(T_2 - T_1)$$
(3.15)

3.6 LINEAR RELATION BETWEEN C_{p} AND C_{v}

Specific heat at constant pressure (C_p) is greater than that at constant volume (C_p) . This can be established as following.

From eqns. (3.15) and (3.11), we can write

$$mC_{p} (T_{2} - T_{1}) = mC_{v} (T_{2} - T_{1}) + mR (T_{2} - T_{1})$$

$$\Rightarrow \qquad C_{p} = C_{v} + R \qquad (3.16)$$
Hence, characteristic gas constant $R = C_{v} - C_{v}$

$$(3.17)$$

Hence, characteristic gas constant $R = C_p - C_v$

3.7 RATIO OF C_p AND C_v

The ratio of C_p and C_v of a gas is an important parameter in the field of thermodynamics. It is known as *adiabatic index* and symbolized by γ . Dividing eqn. (3.16), by C_{ν} we obtain,

$$\frac{C_p}{C_v} = 1 + \frac{R}{C_v}$$

$$\gamma = 1 + \frac{R}{C_v}$$
(3.18)

The values of C_p and C_v of some common gases along with their γ are given in table 3.1.

Sl.No.	Name of gases	C _p (kJ/kgK)	C _v (kJ/kgK)	γ
1	Air	1.000	0.720	1.400
2	Oxygen (O ₂)	0.913	0.653	1.390
3	Hydrogen (H ₂)	14.257	10.133	1.400
4	Nitrogen (N ₂)	1.043	0.745	1.400
5	Carbon-di-oxide (CO_2)	0.846	0.657	1.290
6	Carbon-mono-oxide (CO)	1.047	0.749	1.400
				Contd

Table: 3.1. The values of C_{p} , C_{y} and γ

7	Ammonia (NH ₃)	2.177	1.692	1.290
8	Methane (CH ₄)	2.169	1.650	1.310
9	Helium (He)	5.234	3.153	1.660
10	Argon (A)	5.234	3.153	1.660

3.8 ENTHALPY OF A GAS

One of the basic qualities most frequently recurring in thermodynamics is *enthalpy*. It is the sum of internal energy (U) and the product of pressure and volume ($p\forall$). It is symbolized by H.

So,

$$H = U + p \forall \tag{3.19}$$

Since, $(U + p \forall)$ is made up entirely of property, therefore, enthalpy (*H*) is also a property. Specific enthalpy (*h*) is enthalpy per unit mass.

So, by unit mass consideration we have,

$$h = u + pv_s \tag{3.20}$$

Again we know,

$$(Q_{1-2})_{p} = (U_{2} - U_{1}) + p (\forall_{2} - \forall_{1})$$

= $(U_{2} + p \forall_{2}) - (U_{1} + p \forall_{1}) = H_{2} - H_{1}$ (3.21)
 $q_{1-2} = h_{2} - h_{1}$ (3.22)

For unit mass,

Thus, for a constant temperature process, total heat supplied to the gas is equal to the change of enthalpy.

3.9 THERMODYNAMIC PROCESSES

(a) When a system changes from one equilibrium state to another equilibrium state, the path of successive states through which the system has passed, is known as a *thermodynamic process*. But in actual practice, no system is in true equilibrium during the process because the properties such as pressure, volume, temperature etc., are changing. However, if the process is assumed to take place sufficiently slowly so that the deviation of the properties at the intermediate states are infinitesimally small, then every state passed through by the system will be in equilibrium and the process is named as *quasi-static* or *reversible* process. It is represented by a *continuous curve* on the pressure-volume diagram in Fig. 3.1 (a) and (b).



Fig. 3.1 (a) and (b)

(b) If the process takes place in such a manner that the properties at the intermediate states are not in equilibrium state, except the initial and final states, then the process is said to be *non-equilibrium* or *irreversible* process, represented by broken line on pressure-volume diagram.

3.10 CLASSIFICATION OF THERMODYNAMIC PROCESS

All the thermodynamic processes are classified into two groups:

- (i) Non-flow process: It is that process which occurs in a closed system and does not permit the transfer of mass across their boundaries. In these processes, mass transfer does not occur but energy crosses the system boundary in the form of heat and work. Work done in this process is, $W_{1-2} = \int p d \forall$.
- (ii) Flow process: It is that process which does occur in open system and permits the transfer of mass into and from the system. The flow processes may be steady or non-steady flow processes. The common examples of steady flow processes are flow through nozzles, turbine,

compressors etc. Work done in flow process is $W_{\text{flow}} = \int p d\forall + p_{V_s}$.

3.11 HEATING AND EXPANSION OF GASES IN NON-FLOW PROCESS

The heating and expansion of a gas may be performed in many ways. Following are the different non-flow processes as applied to perfect gas. It may be reversible or irreversible.

- (a) Reversible non-flow process:
 - (i) Constant volume (Isochoric) process
 - (ii) Constant pressure (Isobaric) process
 - (iii) Hyperbolic process
 - (iv) Constant temperature (Isothermal) process
 - (v) Adiabatic (Isentropic) process
 - (vi) Polytropic process
- (b) Irreversible non-flow process: The free expansion process is an irreversible non-flow process. In case of cooling and compression the above mentioned processes are also applicable. Cooling is negative heating and compression is negative expansion.

3.12 CONSTANT VOLUME PROCESS OR ISOCHORIC PROCESS

When gas is heated at constant volume, its pressure and temperature will increase. The process is shown on the pressure-volume ($p-\forall$) and pressure-temperature (p-T) diagrams in Fig. 3.2(a) and (b) respectively. Here follows some relationship within this process.

(a) Pressure, volume-temperature ($p - \forall -T$) relationship: The general equation of gas

$$\frac{p_1 \forall_1}{T_1} = \frac{p_2 \forall_2}{T_2}$$

is



Since, the gas is heated at constant volume, $\forall_1 = \forall_2$

So,
$$\frac{p_1}{T_1} = \frac{p_2}{T_2}$$
 or $\frac{p}{T} = \text{constant}$ (3.23)

Thus, the constant volume process is governed by Gay-Lussac's law.

(b) Work done by the gas:

We know that, $\delta W = pd \forall$



On integrating from state 1 to state 2, we obtain

$$\int_{1}^{2} \delta W = \int_{1}^{2} p d \forall = p \int_{1}^{2} d \forall$$

$$\Rightarrow \qquad W_{1-2} = p(\forall_{2} - \forall_{1}) = 0$$

$$(3.24)$$

(c) Change in internal energy: We know that, change in internal energy $dU = mC_v dT$ On integrating from state 1 to state 2, we obtain

$$\int_{1}^{2} dU = \int_{1}^{2} mC_{\nu} dT$$

$$\Rightarrow \qquad U_{2} - U_{1} = mC_{\nu} (T_{2} - T_{1}) \qquad (3.25)$$

(d) Supply of heat: We know heat supplied or heat transfer

$$\delta Q = dU + \delta W$$

On integrating from state 1 to state 2, we have

$$\int_{1}^{2} \delta Q = \int_{1}^{2} dU + \int_{1}^{2} \delta W$$
$$Q_{1-2} = (U_{2} - U_{1}) + W_{1-2} = mC_{v} (T_{2} - T_{1}) \qquad (3.26)$$

 \Rightarrow

Since, the gas is heating at constant volume, $W_{1-2} = 0$

(e) Change in enthalpy

$$dH = dU + d(p\forall)$$

On integrating from state 1 to state 2, we have

$$\int_{1}^{2} dH = \int_{1}^{2} dU + \int_{1}^{2} d(p \forall)$$

$$\Rightarrow \qquad H_{2} - H_{1} = (U_{2} - U_{1}) + p_{1} \forall_{1} - p_{1} \forall_{2}$$

$$= mC_{v} (T_{2} - T_{1}) + mR (T_{2} - T_{1})$$

$$= m (T_{2} - T_{1}) (C_{v} + R)$$

$$= m(T_{2} - T_{1}) C_{p}$$

$$[\because C_{p} - C_{v} = R]$$

$$(3.27)$$

3.13 CONSTANT PRESSURE PROCESS OR ISOBARIC PROCESS

When a gas is heated at constant pressure, its temperature and volume both will increase. Since, there is a change in its volume, therefore the heat supplied to the gas is utilized to increase the internal energy of the gas and for some external work.

Consider *m* kg of a certain gas being heated at a constant pressure from an initial state 1 to a final state 2 shown in the pressure-volume diagram $(p - \forall)$ and pressure-temperature (p - T) diagram in Fig. 3.3 (a) and (b). Now let us derive the following relations.

(a) Pressure, volume-temperature ($p - \forall -T$) relationship

The general equation of gas is
$$\frac{p_1 \forall_1}{T_1} = \frac{p_2 \forall_2}{T_2}$$

Since, the gas is heated at constant pressure, $p_1 = p_2$



So,
$$\frac{\forall_1}{T_1} = \frac{\forall_2}{T_2}$$
 or $\frac{\forall}{T} = \text{constant}$ (3.28)

Thus, the constant pressure process is governed by Charle's law.

(b) Work done by the gas

We know that, $\delta W = pd \forall$



Fig. 3.3. (b)

On integrating from state 1 to state 2, we have

2

$$\delta W = \int_{1}^{2} p d \forall = p \int_{1}^{2} d \forall$$

 \Rightarrow

$$W_{1-2} = p(\forall_2 - \forall_1) = mR(T_2 - T_1)$$
(3.29)

(c) Change in internal energy

We know that, change in internal energy $dU = mC_v dT$ On integrating from state 1 to state 2, we obtain

$$\int_{1}^{2} dU = \int_{1}^{2} mC_{\nu} dT$$

$$U_{2} - U_{1} = mC_{\nu}(T_{2} - T_{1})$$
(3.30)

 \Rightarrow

(d) Supply of heat

We know heat supplied or heat transfer $\delta Q = dU + \delta W$

On integrating from state 1 to state 2, we have

$$\int_{1}^{2} \delta Q = \int_{1}^{2} dU + \int_{1}^{2} \delta W$$

$$\Rightarrow \qquad (Q_{1-2})_{p} = (U_{2} - U_{1}) + (W_{1-2})_{p}$$

$$= mC_{v}(T_{2} - T_{1}) + mR(T_{2} - T_{1})$$

$$= m(T_{2} - T_{1}) (C_{v} + R)$$

$$= m(T_{2} - T_{1})C_{p} \qquad [\because C_{p} - C_{v} = R] \qquad (3.31)$$

(e) Change in enthalpy

$$dH = dU + d(p\forall)$$

On integrating from state 1 to state 2, we have

$$\int_{1}^{2} dH = \int_{1}^{2} dU + \int_{1}^{2} d(p \forall)$$

$$\Rightarrow \qquad H_{2} - H_{1} = (U_{2} - U_{1}) + (p_{1} \forall_{1} - p_{1} \forall_{1})$$

$$= mC_{v}(T_{2} - T_{1}) + mR(T_{2} - T_{1})$$

$$= m(T_{2} - T_{1}) (C_{v} + R)$$

$$= m(T_{2} - T_{1})C_{p} \qquad \left[\because C_{p} - C_{v} = R\right]$$

$$(3.32)$$

3.14 HYPERBOLIC PROCESS

It is a process, in which the gas is heated as expanded in such a way that the product of its pressure and volume remains constant. The hyperbolic process follows law $p \forall = \text{constant}$, so it is governed by Boyle's law. The hyperbolic process is shown in diagram in Fig. 3.4. Its practical application is *isothermal process*, which is discussed below.



3.15 CONSTANT TEMPERATURE PROCESS OR ISOTHERMAL PROCESS

During this process, the temperature of the working substance remains constant. This process is a very slow process. Let the Mass of the working substance is m kg and being headed at constant temperature from initial state 1 to final state 2 which is shown on $((p - \forall))$ and (p - T) diagrams in Fig. 3.5 (a) and (b) respectively. The relations for the reversible constant temperature process or isothermal process are discussed below.

(a) Pressure, volume-temperature $((p - \forall - T))$ relationship



Since, the gas is heated at constant temperature, $T_1 = T_2$

So,
$$p_1 \forall_1 = p_2 \forall_2$$
 or $p \forall = \text{constant}$ (3.33)

Thus, the constant temperature process is governed by Boyle's law.

(b) Work done by the gas

 \Rightarrow

We know that, $\delta W = pd \forall$

On integrating from state 1 to state 2

$$\int_{1}^{2} \delta W = \int_{1}^{2} p d \forall$$
$$(W_{1-2})_{T} = \int_{1}^{2} p d \forall$$



In isothermal process $p \forall$ = constant



So,

 $p = \frac{p_1 \forall_1}{\forall}$

 $p \forall = p_1 \forall_1$

Hence,

 \Rightarrow

$$(W_{1-2})_T = \int_{\forall_1}^{\forall_2} \frac{p_1 \forall_1}{\forall} d\forall = p_1 \forall_1 \int_{\forall_1}^{\forall_2} \frac{d\forall}{\forall} = p_1 \forall_1 \ln\left(\frac{\forall_2}{\forall_1}\right)$$
(3.34)

Again we know, expansion ratio (r) is the ratio of volume at the end of expansion to volume at the beginning of expansion and compression ratio (r) is the ratio of volume at the beginning of

compression to volume at the end of compression and $r = \frac{\forall_2}{\forall_1}$

But

Hence, finally
$$(W_{1-2})_T = mRT \ln r$$
 (3.35)

(c) Change in internal energy

We know that, change in internal energy $dU = mC_v dT$ In constant temperature condition, $T_1 = T_2$ and dU = 0So, initial internal energy $(U_1) =$ final internal energy (U_2) .

 $p_1 \forall_1 = p_2 \forall_2 = mRT$

(d) Supply of heat

We know heat supplied or heat transfer $\delta Q = dU + \delta W$ On integrating from state 1 to state 2

$$\int_{1}^{2} \delta Q = \int_{1}^{2} dU + \int_{1}^{2} \delta W$$

$$\Rightarrow \quad Q_{1-2} = (U_{2} - U_{1}) + W_{1-2} = W_{1-2} \int_{1}^{2} \delta Q = p_{1} \forall_{1} \ln\left(\frac{\forall_{2}}{\forall_{1}}\right) = mRT \ln\left(\frac{\forall_{2}}{\forall_{1}}\right) \qquad (3.36)$$

(e) Change in enthalpy

$$\int_{1}^{2} dH = H_2 - H_1 = mC_p (T_2 - T_1)$$

In constant temperature condition, $T_1 = T_2$

$$\int_{1}^{2} dH = H_{2} - H_{1} = 0$$

$$H_{2} = H_{1}$$
(3.37)

 \Rightarrow

So, initial enthalpy (H_1) = final enthalpy (H_2) (3.38)

3.16 ADIABATIC PROCESS

During the process, if the working substance neither receives nor gives out heat to its surroundings, is called adiabatic process. It may be *reversible* or *irreversible*. The reversible or *frictionless adiabatic process* is known as *isentropic process*. In isentropic process entropy is constant, i.e.,

$$\int dS = 0$$

Therefore, an isentropic process need not be adiabatic or reversible. If the isentropic process is reversible, it must be adiabatic. An adiabatic process need not be an isentropic, since entropy can also increase due to friction etc. But if the process is reversible adiabatic it must be isentropic. Now consider *m* kg of certain gas being heated adiabatically from an initial state 1 to final state 2 shown in $(p - \forall)$ diagram in Fig. 3.6. Relation for a reversible adiabatic process is discussed below.



(a) Pressure, volume-temperature ($p - \forall -T$) relationship

We know, heat supplied during process, $\delta Q = dU + \delta W$ In an adiabatic process, $\delta Q = 0$, So finally,

$$dU + \delta W = 0$$

 $\Rightarrow \qquad mC_v dT + pd \forall = 0$

$$\Rightarrow \qquad dT = -\frac{pd\forall}{mC_v dT}$$

Again we know, $p \forall = mRT$

On differentiation we get,

$$pd\forall + \forall dp = mRdT$$

⇒	$dT = \frac{pd\forall + \forall dp}{mR} = \frac{pd\forall + \forall dp}{m(C_p - C_v)}$
Hence,	$-\frac{pd\forall}{mC_{v}dT} = \frac{pd\forall + \forall dp}{m(C_{p} - C_{v})}$
⇒	$\frac{C_p - C_v}{C_v} = -\frac{pd\forall + \forall dp}{pd\forall}$
⇒	$\frac{C_p}{C_v} - 1 = -1 - \frac{\forall dp}{pd\forall}$
\Rightarrow	$\gamma = -\frac{\forall dp}{pd\forall}$
\Rightarrow	$\gamma \frac{d \forall}{\forall} = -\frac{dp}{p}$
Integrating b	both sides,
	$\gamma \ln \forall + \ln p = \text{constant}$

$$\Rightarrow \qquad \ln(p\forall^{\gamma}) = \text{constant}$$

$$\Rightarrow \qquad p\forall^{\gamma} = \text{constant}$$

$$\Rightarrow \qquad p\forall^{\gamma} = \text{constant}$$

$$\Rightarrow \qquad p_{1}\forall_{1}^{\gamma} = p_{2}\forall_{2}^{\gamma} = \text{constant} \qquad (3.38)$$
The shows constant one can be a written in this form as

The above equation can also be written in this form as,

$$\frac{p_1}{p_2} = \left(\frac{\forall_2}{\forall_1}\right)^{\gamma} \tag{3.39}$$

We know from general gas equation,

$$\frac{p_1}{p_2} = \frac{T_1}{T_2} \times \frac{\forall_2}{\forall_1}$$

Hence we get, $(\forall, \gamma)^{\gamma}$

$$\left(\frac{\forall_2}{\forall_1}\right)^r = \frac{T_1}{T_2} \times \frac{\forall_2}{\forall_1}$$
$$\frac{T_1}{T_2} = \left(\frac{\forall_2}{\forall_1}\right)^{\gamma-1}$$
(3.40)

 \Rightarrow

Again we can write,

$$\frac{\forall_2}{\forall_1} = \left(\frac{p_1}{p_2}\right)^{1/\gamma} \tag{3.41}$$

Substituting the values, in equations (3.40), we obtained

$$\frac{T_1}{T_2} = \left(\frac{p_1}{p_2}\right)^{\frac{\gamma-1}{\gamma}}$$
(3.42)

Comparing the values, of equations (3.40) and (3.42) we can write

$$\frac{T_1}{T_2} = \left(\frac{p_1}{p_2}\right)^{\frac{\gamma-1}{\gamma}} = \left(\frac{\forall_2}{\forall_1}\right)^{\gamma-1}$$
(3.43)

The above equation is known as Poisson's Equation.

(b) Work done

We know that, work done $\delta W = pd \forall$ On integrating from state 1 to state 2, we obtain

$$\int_{1}^{2} \delta W = \int_{1}^{2} p d \forall$$
$$W_{1-2} = \int_{1}^{2} p d \forall$$

 \Rightarrow

In adiabatic process of gas

$$p \forall^{\gamma} = p_1 \forall_1^{\gamma} = \text{constant}$$

 \Rightarrow

 $p = p_1 \left(\frac{\forall_1}{\forall}\right)^{\gamma}$

Substituting the value of *p*, we have

$$W_{1-2} = \int_{1}^{2} \frac{p_{1} \forall_{1}^{\gamma}}{\forall^{\gamma}} d\forall$$
$$= p_{1} \forall_{1}^{\gamma} \left[\frac{\forall^{1-\gamma}}{1-\gamma} \right]_{1}^{2} = \frac{p_{1} \forall_{1}^{\gamma} \forall_{2}^{1-\gamma} - p_{1} \forall_{1}^{\gamma} \forall_{1}^{1-\gamma}}{1-\gamma}$$
$$= \frac{p_{2} \forall_{2}^{\gamma} \forall_{2}^{1-\gamma} - p_{1} \forall_{1}^{\gamma} \forall_{1}^{1-\gamma}}{1-\gamma} = \frac{p_{2} \forall_{2} - p_{1} \forall_{1}}{1-\gamma}$$
(3.44)

The equation for work done may also be expressed as follows. We know that $p_1 \forall_1 = mRT_1$ and $p_2 \forall_2 = mRT_2$. Substituting the above values, we have

$$W_{1-2} = \frac{p_2 \forall_2 - p_1 \forall_1}{1 - \gamma} = \frac{mRT_2 - mRT_1}{1 - \gamma} = mR \frac{T_2 - T_1}{1 - \gamma}$$
(3.45)

It is also expressed as

$$W_{1-2} = \frac{p_1 \forall_1 - p_2 \forall_2}{\gamma - 1} = \frac{p_1 \forall_1}{\gamma - 1} \left(1 - \frac{p_2 \forall_2}{p_1 \forall_1} \right) = \frac{mRT_1}{\gamma - 1} \left(1 - \frac{p_2 \forall_2}{p_1 \forall_1} \right)$$
(3.46)

(c) Change in internal energy: Change in internal energy

$$\int_{1}^{2} dU = U_{2} - U_{1} = mC_{\nu} (T_{2} - T_{1})$$
(3.47)

(d) Supply of heat: Heat supplied during the process

$$Q_{1-2} = 0 \tag{3.48}$$

(e) Change in enthalpy:

$$\int_{1}^{2} dH = H_{2} - H_{1} = mC_{p} (T_{2} - T_{1})$$
(3.49)

3.17 POLYTROPIC PROCESS

The polytropic process is also known as the general law for expansion and compression of gases and it is given by the relation $p \forall^n = \text{constant}$. Here *n* is a polytropic index, which may assume any value from zero to infinity, depending upon the manner in which the expansion and compression has taken place. The various equations for polytropic process may be expressed by changing the index (*n* for γ in the adiabatic process). Let usconsider *m* kg of a certain gas being heated polytropically from an initial state 1 to a final state 2, shown in Fig. 3.7 in the ($p - \forall$) diagram. The following relations for the polytropic process can be derived below.



(i)
$$p_1 \forall_1^n = p_2 \forall_2^n = \text{constant}$$
 (3.50)

(ii)
$$\frac{T_1}{T_2} = \left(\frac{p_1}{p_2}\right)^{\frac{n-1}{n}} = \left(\frac{\forall_2}{\forall_1}\right)^{n-1}$$
(3.51)



(b) Work done: The equation for the work done during a polytropic process may also be expressed by changing the index *n* for g in the adiabatic process. So we have,

$$W_{1-2} = \frac{p_2 \forall_2 - p_1 \forall_1}{1 - n} = \frac{mRT_2 - mRT_1}{1 - n} = mR\frac{T_2 - T_1}{1 - n}$$
(3.52)

$$W_{1-2} = \frac{p_1 \forall_1 - p_2 \forall_2}{n-1} = \frac{p_1 \forall_1}{n-1} \left(1 - \frac{p_2 \forall_2}{p_1 \forall_1} \right) = \frac{mRT_1}{n-1} \left(1 - \frac{p_2 \forall_2}{p_1 \forall_1} \right) (3.53)$$

(c) Change in internal energy: Change in internal energy

$$\int_{1}^{2} dU = U_{2} - U_{1} = mC_{\nu} (T_{2} - T_{1})$$
(3.54)

 \smile

(d) Supply of heat: Heat supplied during the process

$$\delta Q = dU + \delta W$$

On integration, we have,

$$Q_{1-2} = U_{1-2} + W_{1-2} = mC_{\nu}(T_2 - T_1) + \frac{p_1 \nabla_1 - p_2 \nabla_2}{n-1}$$

$$= mR \frac{T_2 - T_1}{\gamma - 1} + mR \frac{T_1 - T_2}{n-1}$$

$$= mR(T_1 - T_2) \left[\frac{1}{n-1} - \frac{1}{\gamma - 1} \right]$$

$$= mR(T_1 - T_2) \left[\frac{\gamma - n}{(n-1)(\gamma - 1)} \right]$$
(3.55)

 \smile

(e) Change in enthalpy

$$\int_{1}^{2} dU = H_2 - H_1 = mC_p (T_2 - T_1)$$
(3.56)

3.18 GENERAL LAWS FOR EXPANSION AND COMPRESSION

The general laws for expansion and compression of a perfect gas is $p \forall^n = \text{constant}$, when *n* is the index of process and it gives the relationship between pressure and volume of a given quantity of gas. The value of *n* depends upon the manner of expansion and compression. The following values of *n* are important from the subject point of view.

(a) When n = 0 the equation will be $p \forall^0 = p = \text{constant}$. So, expansion and compression will occur under constant pressure and it is *constant pressure process*.

- (b) When n = 1 then equation will be $p \forall^1 = p \forall = \text{constant}$. The process is *isothermal* or *hyperbolic process*.
- (c) When $\gamma > n > 1$, the expansion or compression is *polytropic*, i.e. $p \forall^n = \text{constant}$.
- (d) When $n = \gamma$, the expansion or compression will be *adiabatic*, i.e., $p \forall^{\gamma} = \text{constant}$.
- (e) When $n = \alpha$, the equation will be

$$p \forall^{\alpha} = C$$

$$\Rightarrow \qquad \forall = \left(\frac{C}{p}\right)^{\frac{1}{\alpha}} \alpha \sqrt{\frac{C}{p}} = \text{constant, i.e., constant volume process}$$

For expansion of a perfect gas for different values of *n* is shown in pressure-volume diagram in Fig. 3.8.

3.19 REAL GAS

A hypothetical gas which obeys the law $pv_s = RT$ at all pressure and temperature is called an *ideal* gas. But in actual practice, there is no actual gas which strictly obeys this gas law over the entire range of temperature and pressure and is called the *real gas*. The real gases are ordinarily difficult to liquefy. The real gas such as oxygen, hydrogen and air within certain temperature and pressure limits may be regarded as ideal or perfect gases. There is small variation of behaviour in real gas from ideal gas.

3.20 REAL GAS AND COMPRESSIBILITY FACTOR

We are assuming specific volume of an ideal gas where the equation of gas is $p(v_s)_{ideal} = RT$.

So,
$$(v_s)_{ideal} = \frac{RT}{p}$$

If specific volume of a real gas is considered, then, $p(v_s)_{real} = RT$ and in that case,

 $(v_s)_{ideal} = Z \frac{RT}{p}$. Here the term Z is being introducing to signify its variation from ideal

gas behaviour and is named as compressibility factor, or just compressibility.

By definition the value of Z,

$$Z = \frac{p(v_s)_{\text{real}}}{RT} = \frac{(v_s)_{\text{real}}}{RT/p} = \frac{(v_s)_{\text{real}}}{(v_s)_{\text{ideal}}}$$
(3.57)



So, compressibility factor or compressibility is the ratio between *specific volumes of real gas* and *specific volume of ideal gas*. It is a dimensionless factor. For ideal gas Z = 1 and for real gas Z is equal to or greater than or less than 1. The *non ideality of gases* is the differences of the real gas volume and volume predicted by ideal gas equation, denoted as α or \forall^R where,

$$= \alpha = (v_s)_{ideal} - (v_s)_{real} = (1 - Z)\frac{RT}{p}$$
(3.58)

 \forall^R or α is also termed as the *residual volume of the gas*. In terms of molar quantities,

$$Z = \frac{pv_m}{R_m T}$$
 or $v_m = \frac{R_m T}{p}$.

An equation of state can be written in terms of Z as a function of p and T instead of v_s as a function of p and T. So, Z becomes as alternative property. Thus the experimental $p - v_s - T$ data of pure substance can also be fitted into an equation of state of the form Z = f(p, T) and can be plotted on a graph of Z versus p with T as parameter. Such a plot is referred to as the compressibility chart. Fig. 3.9 presents a typical compressibility chart for nitrogen.



3.21 LAW OF CORRESPONDING STATE AND GENERALIZED COMPRESSIBILITY CHART

For a gas, the compressibility factor Z is a function of p and T. Thus, graphs in Fig. 3.10 can be plotted for p versus Z at the line of constant temperature and from these graphs, the value of Z can be found for any value of p and T and then volume can be calculated from the equation.

$$v_s = Z \frac{RT}{p}.$$
(3.59)

For different substance a different compressibility chart is needed. However, one single generalized chart can be developed. The basis for this is the law of corresponding states. For the purpose we consider the dimensionless property called reduced properties p_r , T_r , $v_{s(r)}$, known as the reduced pressure, reduced temperature and reduced specific volume respectively. The reduced temperature of a substance is defined as the ratio of the temperature of the substance to its critical

temperature, i.e.,
$$T_r = \frac{T}{T_c}$$
. Similarity reduced pressure



 $p_r = \frac{P}{P_c}$ and reduced specific volume $v_{s(r)} = \frac{v_s}{v_{s(c)}}$, where p_c , T_c , $v_{s(c)}$ are the critical properties of the substance.

Now the law of corresponding state may be stated as follows: If two or more substance have the same reduced pressure and reduced temperature then they will have the some reduced volume. Mathematically, it can be expressed in the form $Z = f(p_r, T_r)$. This relation is referred to as the generalized equation of state.

This functional relationship can be expressed in terms of van der Walls equation,

i.e.,
$$\left(p + \frac{a}{v_s}\right)(v_s - b) = RT$$
(3.60)

Here

 $a = \frac{27R^2T_c^2}{64p_c}$ and $b = \frac{RT_c}{8p_c}$

Hence







Now, if the actual value of Z is plotted against the values of p_r for constant T_r , one would get the some curve for all substances. Such a plot, as shown in Fig. 3.11 is known as the generalized compressibility chart. The following observations can be made from this chart:

- (i) At all temperatures $Z \to 1$ as $p_r \to 1$.
- (ii) At temperature equals to twice the critical temperature $T_r = 2$ and above, Z = 1 over a wide range of pressures up to five times the critical pressure $p_r = 5$.
- (iii) The compressibility factor at the critical point $Z_c = \frac{p_c v_{s(c)}}{RT} = \frac{p_c v_m}{R_m T_c}$ is called the critical

compressibility and is found to be equal to 0.275 for all substances.

	1	2	3	4
Reversible non-flow process $\frac{p_1 \forall_1}{T_1} = \frac{p_2 \forall_2}{T_2} = C$	Isobaric (Constant pressure)	Isobaric Isochoric (Constant pressure) (Constant volume)		Adiabatic or Isentropic
Index n	n = 0	$n = \alpha$	<i>n</i> = 1	$n = \gamma$
p−∀ diagram	$p_1 = p_2$ $p_1 = p_2$ $p_2 = p_2$ $p_3 = p_2$ $p_3 = p_3$ $p_3 = p_3$ p_3	p_2 p_2 p_1 p_1 p_1 p_1 p_2	p_1 p_1 p_2	p_1 Adiabatic Expansion $(p\forall^{\gamma} = C)$ p_2 \forall_1 Volume \forall_2
$p - \forall -T$ relation	$\frac{\forall_1}{T_1} = \frac{\forall_2}{T_2}$	$\frac{p_1}{T_1} = \frac{p_2}{T_2}$	$p_1 \forall_1 = p_2 \forall_2$	$p_1 \forall_1^{\gamma} = p_2 \forall_2^{\gamma}$
	[Charle's Law]	[Gay-Lussac's Law]	[Boyle's Law]	$\frac{T_1}{T_2} = \left(\frac{p_1}{p_2}\right)^{\frac{\gamma-1}{\gamma}} = \left(\frac{\forall_2}{\forall_1}\right)^{\frac{\gamma-1}{\gamma}}$



	1	2	3	4
Work done $W_{1-2} = \int_{1}^{2} p d \forall$	$p(\forall_2 - \forall_1)$ or $mR(T_2 - T_1)$	0	$p_{1} \forall_{1} \ln \left(\frac{\forall_{2}}{\forall_{1}} \right)$ or $mRT \ln \left(\frac{\forall_{2}}{\forall_{1}} \right)$	$\frac{p_2 \forall_2 - p_1 \forall_1}{1 - \gamma}$ or $\frac{mR(T_2 - T_1)}{1 - \gamma}$
Specific heat (C)	C_p	C_{v}	α	0
Heat added $Q_{1-2} = \int_{1}^{2} dU + W_{1-2}$	$mC_p(T_2 - T_1)$	$mC_{v}(T_2 - T_1)$	$p_1 \forall_1 \ln \left(\frac{\forall_2}{\forall_1} \right)$	0
Change in internal energy $\int_{1}^{2} dU = U_{2} - U_{1}$	$mC_{v}(T_2 - T_1)$	$mC_{v}(T_2 - T_1)$	0	$mC_{v}(T_2 - T_1)$
Change in enthalpy $\int_{1}^{2} dH = H_2 - H_1$	$mC_p(T_2 - T_1)$	$mC_p(T_2 - T_1)$	0	$mC_p(T_2 - T_1)$

 $\frac{Prc}{}$ rties a чdy nic Pr

	5	
	5	0
Work done $W_{1-2} = \int_{1}^{2} p d \forall$	$\frac{p_2 \forall_2 - p_1 \forall_1}{1 - n}$ or $\frac{mR(T_2 - T_1)}{1 - n}$	
Specific heat (C)		
Heat added	$\frac{mR\left(T_2 - T_1\right)}{1 - n} \left(\frac{\gamma - n}{\gamma - 1}\right)$	
$Q_{1-2} = \int_{1}^{2} dU + W_{1-2}$	or $\frac{p_2 \forall_2 - p_1 \forall_1}{1 - n} \left(\frac{\gamma - n}{\gamma - 1} \right)$	
Change in internal energy	$mC_{\nu}(T_2 - T_1)$	
$\int_{1}^{2} dU = U_{2} - U_{1}$		
Change in enthalpy	$mC_p(T_2 - T_1)$	
$\int_{1}^{2} dH = H_2 - H_1$		

Mechanical Science-II





	5	6
Change in entropy in $T-S$ diagram		General Gas Equation:
		$S_2 - S_1 = m \left[C_v \ln \left(\frac{T_2}{T_1} \right) + (C_p - C_v) \ln \left(\frac{\forall_2}{\forall_1} \right) \right]$ or
	$m(C_p - nC_v) \ln\left(\frac{\forall_2}{\forall_1}\right)$	$m\left[C_p \ln\left(\frac{T_2}{T_1}\right) + (C_p - C_v) \ln\left(\frac{p_1}{p_2}\right)\right] \text{ or }$
Change in entropy = $S_2 - S_1$	or $mC_{\nu} \frac{\gamma - n}{n} \ln\left(\frac{p_1}{p_2}\right)$	$m\left[C_{\nu}\ln\left(\frac{p_2}{p_1}\right) + C_p\ln\left(\frac{\forall_2}{\forall_1}\right)\right]$
Steady flow process work done W_{1-2}	$\frac{n}{n-1}(p_1\forall_1-p_2\forall_2)$	

	Boiler	Condenser	Evaporator	Nozzle			
Steady flow energy equation to engineering systems $(u_1 + p_1 v_{s1}) + \frac{V_1^2}{2} + gZ_1 + q_{1-2}$ $= (u_2 + p_2 v_{s2}) + \frac{V_2^2}{2} + gZ_2 + w_{1-2}$	$q_{1-2} = h_2 - h_1$	$-q_{1-2} = h_2 - h_1$	$q_{1-2} = h_2 - h_1$	$V_2 = \sqrt{2}$	$=\sqrt{V_{1}^{2}+2(h_{1}-h_{2})}$		
or	Turbine	Rotary Compressor	Reciprocating Compressor	1	Heat Exchanger		
$m\left(h_{1} + \frac{V_{1}^{2}}{2} + gZ_{1} + q_{1-2}\right)$ $= m\left(h_{2} + \frac{V_{2}^{2}}{2} + gZ_{2} + w_{1-2}\right)$	$w_{1-2} = h_1 - h_2$	$w_{1-2} = h_2 - h_1$	$w_{1-2} = q_{1-2} + (h_2 - h_2)$	h ₁)	$q_{1-2} = h_1 - h_2$		

Prop

Multiple Choice Questions

1.	If the p	temperature remains constant, the vor ressure this is known as	olume	of a given mass of a gas is inversely proportional to
	(a)	Charle's law	(h)	Boyle's law
	(u) (c)	Joule's law	(d)	Gav-Lussac's law.
2.	The s	state of a substance whose evaporatio	n form	n it's liquid state is complete is known as
	(a)	steam	(b)	vapour
	(c)	air	(d)	perfect gas
3.	The o	characteristic equation of a gas is		
	(a)	$p \forall = \text{constant}$	(b)	$p \forall = mR$
	(c)	$p \forall = mRT$	(d)	$p\forall = RT^m$
4.	The	value of gas constant (R) is		
	(a)	287 J/kg	(b)	28.7 J/kg
	(c)	2.87 J/kg	(d)	0.287 J/kg
5.	The	value of universal gas constant (R_n) is	5	
	(a)	8.314 J/kg K	(b)	83.14 J/kg K
	(c)	831.4 J/kg K	(d)	8314 J/kg K
6.	The g	gas constant (R) is equal to the		of two specific heats
	(a)	sum	(b)	difference
	(c)	product	(d)	ratio
7.	The s	specific heat at constant pressure is _		that of specific heat at constant volume
	(a) (c)	equal to more than	(b)	less than
8	The	ratio of specific heat at constant press	aure (C) and specific heat at constant volume (C) is
0.	(a)	equal to one	(b)	less than one
	(c)	more than one	(d)	none of these
9.	The v	value of C/C for air is		
	(a)	p = v	(b)	1.4
	(c)	1.8	(d)	2.3
10.	When	n the gas is heated at constant pressu	ire the	n the heat supplied
	(a)	raises the temperature of the gas		
	(b)	increases the internal energy of the	gas	
	(c)	does some external work during exp	ansior	1
	(d)	both (a) and (b)		
	(e)	both (b) and (c)		
11.	When	n a gas is heated at constant volume		
	(a)	its temperature will increase		
	(b)	its pressure will increase		

- (c) both temperature and pressure will increase
- (d) neither temperature nor pressure will increase

- 12. The heating of a gas at constant pressure is governed by
 - Boyle's law (b) Charle's law (a)
 - (c) Gay-Lussac's law (d) Joule's law
- 13. A process in which the gas is heated or expanded in such a way that the product of it's pressure and volume remains constant, is called
 - (a) isothermal process isobaric process (b)
 - (c) adiabatic process (d) polytropic process
- 14. The hyperbolic process is governed by
 - Boyle's law Charle's law (a) (b)
 - (c) Gay-Lussac's law (d) Joule's law
- 15. The heating of gas at constant is governed by Boyle's law
 - (a) volume (b) pressure none of these
 - temperature (d) (c)
- 16. In an isothermal process

(a)

(c)

(c)

- internal energy increases (b) internal energy decreases
- there is no change in internal energy (c)
- (d) internal energy first increases and then decreases
- 17. The expansion ratio (r) is the ratio of

(a)
$$\frac{\forall_1}{\forall_2}$$
 (b) $\frac{\forall_2}{\forall_1}$

(c)
$$\frac{\forall_1 + \forall_2}{\forall_2}$$
 (d) $\frac{\forall_1 + \forall_2}{\forall_1}$

- 18. When the expansion or compression of the gas takes places according to the law $p \forall^n = \text{constant}$, then the process is known as
 - isothermal process isobaric process (a) (b)
 - polytropic process adiabatic process (d)
- 19. An adiabatic process is one in which
 - no heat enters or leaves the gas (a)
 - (b) the temperature of the gas changes
 - the change in internal energy is equal to the work done (c)
 - all of the above (d)

20. The general law of expansion or compression is the process is said to be hyperbolic if n is equal to

(b)

isobaric process

- 0 (a) (b) 1
- (c) 2 (d) none of these
- 21. If the value of n = 0 in the general law, then the process is called
 - isochoric process (a)
 - isothermal process (d) isentropic process

22.	Ifall	the variable of a steam are independe	nt of ti	me it is said to be in				
	(a)	steady flow	(b)	unsteady flow				
	(c)	uniform flow	(d)	closed flow				
	(e)	constant flow						
23.	A cor	ntrol volume refers to						
	(a)	a fixed region in space	(b)	a specified mass				
	(c)	an isolated system	(d)	reversible process only				
	(e)	a closed system						
24.	Intern	hal energy of a perfect gas depends o	n					
	(a)	temperature, specific heats and pres	ssure					
	(b)	temperature, specific heats and enth	nalpy					
	(c)	temperature, specific heats and enth	nalpy					
	(d)	temperature only						
25.	In rev	versible polytropic process						
	(a)	true heat transfer occurs	(b)	the entropy remains constant				
	(c)	the enthalpy remains constant	(d)	the internal energy remains constant				
	(e)	the temperature remains constant						
26.	An is	entropic process is always						
	(a)	irreversible and adiabatic	(b)	reversible and isothermal				
	(c)	frictionless and irreversible	(d)	reversible and adiabatic				
	(e)	none of these						
27.	The r	net work done per kg of gas in a poly	tropic p	process is equal to				
	(a)	$p_1 \forall_1 \ln \frac{\forall_2}{\forall_1}$	(b)	$p_1(\forall_1 - \forall_2)$				
	(c)	$p_2\left(\forall_2 - \frac{\forall_1}{\forall_2}\right)$	(d)	$\frac{p_1 \forall_1 - p_2 \forall_2}{n - 1}$				
	(e)	$\frac{p_1n_1 - p_2n_2}{n-1}$						

- 28. Steady flow occurs when
 - (a) conditions do not change with time at any point
 - (b) conditions are the same at adjacent points at any instant
 - (c) conditions changes steadily with the time
 - (d) $d\forall/dt$ is constant
- 29. A reversible process requires that
 - (a) there be no heat transfer
 - (b) Newton's law of viscosity be satisfied
 - (c) temperature of system and surrounding be equal
 - (d) there be no viscous or coulomb friction in the system
 - (e) heat transfer occurs from surroundings to system only

30.	The f	irst law of thermodynamics for stead	y of lav	N
	(a) (b) (c) (d) (e)	accounts for all energy entering and is an energy balance for a specific r is an expansion of the conservation is primarily concerned with heat tra is restricted in it's application to pe	d leavin mass of of line insfer	ng a control volume f fluid ar momentum ases
31	The c	characteristic equation of gases $n \forall =$	= mRT	holds good for
51.	(a)	monoatomic gases	(h)	diatomic gases
	(c)	real gases	(d)	ideal gases
	(e)	mixture of gases		e
32.	A gas	s which obeys kinetic theory perfectl	y is kn	own as
	(a)	monoatomic gases	(b)	diatomic gases
	(c)	real gases	(d)	pure gases
	(e)	perfect gases		
33.	Work	done in a free expansion process is		
	(a)	zero	(b)	minimum
	(C)	negative	(a)	positive
24	Whice	bh of the following is not a property (of the c	vistem
54.	(a)	temperature	(h)	pressure
	(a)	specific volume	(d)	heat
	(e)	none of these	(4)	
25		$1 \dots 1^n - \dots$		
35.	In a p	polytropic process equation $p \lor = cc$	onstant	. If $n = 0$, the process is termed as
	(a)	constant volume	(b)	constant pressure
	(c) (e)	isothermal	(d)	adiadatic
36.	In a p	polytropic process equation, if n is in	finitely	large, the process is termed as
	(a)	constant volume	(b)	constant pressure
	(c)	constant temperature	(d)	adiabatic
	(e)	isothermal		
37.	The p	process or system that does not invol	lve hea	t are called
	(a)	isothermal process	(b)	equilibrium process
	(C) (e)	adiabatic process	(u)	steady process
38.	Inar	eversible adiabatic process the adiab	atic rat	tio is equal to
		C		
	(a)	$\frac{C_p}{C}$	(b)	$C_p \times C_v$
		C_{v}		
	(c)	$\underline{C_{v}}$	(d)	C_p
	(0)	C_p	(u)	$2C_{v}$
39.	In isc	othermal process		
	(a)	temperature increases gradually	(b)	volume remains constant
	(c)	pressure remains constant	(d)	enthalpy change is maximum
	(e)	change in internal energy is zero		

4	0.	Durir	ng throttling	process	S												
		(a)	internal end	ergy do	bes no	t char	ige	(b)	р	essur	e does	not c	hange				
		(c)	entropy do	es not	chang	e		(d)	eı	nthalp	y does	not c	hange	1			
	1	(e)	volume cha	nge is i	negiig	ible	c				1.1.1	• 1	1.				
4	H.	when	1 a gas is to t	be store	ed, the	type	of con	npres	sion	that v	vould l	be idea	al 1S				
		(a)	isothermal					(b) (d)	a	llabat	IC	me					
		(c) (e)	none of ab	ove				(u)	C	mstan	it volu						
4	12	Ifan	rocess can be	stonne	ed at a	nv sta	ge and	l reve	ersed	so the	at the s	vstem	and si	irround	dings :	are ex	actly
		restor	red to their in	itial st	ates it	is kno	own as	3	1500	So the	te the s	<i>y</i> stern	und be		unigst		uetry
		(a)	adiabatic p	rocess				(b)	is	othern	nal pro	ocess					
		(c)	ideal proce	SS				(d)	fr	iction	less pr	ocess					
		(e)	energy less	proce	SS												
4	3.	The s	tate of a sub	stance	whose	e evap	oration	n froi	n it's	s liqui	d state	is con	nplete	, is kn	own as	s	
		(a)	vapour					(b)	p	erfect	gas						
		(c)	air					(d)	st	eam							
4	4.	In S.I	. the value o	f the ur	nivers	al gas	consta	ant is				_					
		(a)	0.8314 J/mc	le/K				(b)	8.	314 J/	mole/K	-					
		(c) (e)	83.14 J/mole	3/K. ⊳/K				(a)	8.	51.4 J/	mole/K	•					
	15	(c) What	the gas is k	vin Vantad v	at con	stant	process	ra th	a ha	ot cur	nliad						
7	IJ.	(a)	increases th	ne inter	mal er	erov	of the	σas		ai sup	prica						
		(b)	increases th	ie temr	beratu	re of t	he gas	5us S									
		(c)	does some	externa	ıl worl	c durii	ng exp	ansio	n								
		(d)	both (b) an	d (c)													
		(e)	none of the	above	e												
4	6.	The g	gas constant	(<i>R</i>) is e	equal t	o the											
		(a)	sum of two	specifi	ic hea	ts		(b)	di	fferer	nce of t	wo sp	ecific	heats			
	_	(c)	product of	two spo	ecific	heats		(d)	. ra	tio of	two sp	pecific	heats				
4	7.	The l	eat absorbed	l or rejo	ected	during	g a pol	ytrop	nc p	rocess	15						
			$(\gamma - n)$						($n - \gamma$)						
		(a)	$\left(\frac{1}{\gamma-1}\right)$ ×	work c	lone			(b)		$\gamma - 1$)×wo	ork doi	ne				
		(c)	$\gamma \times$ work do	one				(d)	n	γ×wα	ork dor	ne					
							۸.		- r o								
							AI	15W0	12								
1.	(b)	2.	(d) 3. (c) 4.	(a)	5.	(d)	6.	(b)	7.	(c)	8.	(c)	9.	(b)	10.	(d)
11.	(c)	12.	(b) 13. (a) 14.	(a)	15.	(c)	16.	(c)	17.	(b)	18.	(d)	19.	(d)	20.	(b)
21.	(b)	22.	(a) 23. (a) 24.	(d)	25.	(a)	26.	(d)	27.	(d)	28.	(a)	29.	(d)	30.	(a)
31.	(c)	32.	(e) 33. (a) 34.	(d)	35.	(b)	36.	(a)	37.	(e)	38.	(a)	39.	(e)	40.	(d)
41.	(a)	42.	(c) 43. (b) 44.	(e)	45.	(d)	46.	(b)	47.	(a)						

NUMERICAL EXAMPLES

EXAMPLE 1

Some quantity of air at a temperature of 20°C and at a pressure of 2.15 bar is found to occupy 0.20m³. The volume is kept constant when the heat energy equivalent to 20kJ is given to air. Determine the new pressure and temperature. Consider specific heat of air at constant volume 0.706 kJ/kgK. Any missing data may be assumed.

SOLUTION

Here

Let us assume

$$p_1 = \frac{2.15 \times 10^5}{10^3} \text{ kPa},$$

 $\forall_1 = 0.20 \text{ m}^3, C_v = 0.706 \text{ kJ/kgK}, Q_v = 20 \text{ kJ}$
 $R = 0.287 \text{ kJ/kgK}$
We know,
 $p_1 \forall_1 = mRT_1$

T = 20 + 273 = 293 K

We know,

 $2.15 \times 100 \times 0.20 = m \times 0.287 \times 293$

Therefore

 \Rightarrow

 $m = 0.511 \, \text{kg}$

Again, heat supplied at constant volume

 $Q_v = mC_v(T_2 - T_1)$ $20 = 0.511 \times 0.706 (T_2 - 293)$ \Rightarrow $T_{2} = 348.4 \,\mathrm{K}$ \Rightarrow

As, $\frac{p_1}{T_1} = \frac{p_2}{T_2}$ being a constant volume process,

$$p_2 = \frac{p_1 T_2}{T_1} = \frac{2.15 \times 348.4}{293} = 2.55 \text{ bar} = 255 \text{ kPa}$$

EXAMPLE 2

When 0.1421 kg of a gas is heated from 27°C to 127°C, it is observed that the gas requires 202 kJ of heat if heated at constant pressure and 143 kJ of heat if heated at constant volume. Find the adiabatic index, gas constant and molecular weight of the gas.

SOLUTION

Here mass of the gas (m) = 0.1421 kg, initial temperature $(T_1) = 27 + 273 = 300$ K, $(T_2) = 127 + 273 = 400 \text{ K},$ final temperature heat required at constant pressure $Q_n = 202 \text{ kJ}$, heat required at constant volume $Q_v = 143 \text{ kJ}$ $Q_p = mC_p \left(T_2 - T_1\right)$ Now, $202 = 0.1421 (400 - 300)C_p$ \Rightarrow
Therefore	$C_p = 14.125 \text{ kJ/kgK}$
Again,	$Q_{v} = mC_{v}(T_{2} - T_{1})$
\Rightarrow	$143 = 0.1421(400 - 300)C_{v}$
Therefore	$C_v = 10.06 \text{kJ/kgK}$
So, adiabatic index	$\gamma = \frac{C_p}{C_v} = \frac{14.125}{10.06} = 1.41$
Again,	$R = C_p - C_v = 14.125 - 10.06 = 4.065 \text{ kJ/kgK}$
Assuming,	$R_m = 8.314 \text{ kJ/kgK}, M = \frac{R_m}{R} = \frac{8.314}{4.065} \approx 2.$

If 0.112m^3 of gas at 27°C and 1.913 bar is compressed adiabatically to $1/5^{\text{th}}$ of this volume, determine (i) the mass of the gas (ii) the pressure and temperature at end of compression and (iii) the work done on the gas. Assume $C_p = 1.005 \text{kJ/kgK}$, $C_v = 0.718 \text{kJ/kgK}$.

SOLUTION

EXAMPLE 3

Initial volume $\forall_1 = 0.112 \text{ m}^3$, initial temperature $T_1 = (27 + 273) = 300 \text{ K}$, initial pressure $p_1 = 1.013$ bar, final volume $\forall_2 = \forall_1/5$.

Adiabatic index $\gamma = \frac{C_p}{C_v} = \frac{1.005}{0.718} = 1.4 \text{ and}$ $R = C_p - C_v = 0.287 \text{ kJ/kgK}$ Again we know $m = \frac{p_1 \forall_1}{RT_1} = \frac{1.013 \times 100 \times 0.112}{0.287 \times 300} = 0.1317 \text{ kg}$ But $p_2 = \left(\frac{\forall_1}{\forall_2}\right)^{\gamma} p_1 = 1.013 \times 5^{1.4} = 9.642 \text{ bar}$ Again $T_2 = \left(\frac{\forall_1}{\forall_2}\right)^{\gamma-1} T_1 = 300 \times 5^{0.4} = 571.1 \text{ K}$ Work done on the gas $W_{1-2} = \frac{p_2 \forall_2 - p_1 \forall_1}{\gamma - 1} = \frac{mR(T_2 - T_1)}{\gamma - 1} = 25.617 \text{ kJ}$

EXAMPLE 4.

A system contains 0.15 m³ of a gas at a pressure of 3.8 bar and 150°C. It is expanded adiabatically till the pressure falls to 1 bar. The gas is then heated at a constant pressure till it's enthalpy increases by 70 kJ.Determine the total work done, Take $C_p = 1.00 \text{ kJ/kgK}$, $C_v = 0.714 \text{ kJ/kgK}$.



SOLUTION

Here, $\forall_1 = 0.15 \text{ m}^3$, $p_1 = 3.8 \text{ bar}$, $p_2 = 1.0 \text{ bar}$, $T_1 = 150 + 273 = 423 \text{ K}$, dH = 70 kJ.

In Fig., process 1–2 represents adiabatic expansion of the gas and the process 2–3 represents heating at constant pressure. First of all, let us find the temperature T_2 and volume \forall_2 after the adiabatic expansion.

We know that adiabatic index, $\gamma = \frac{C_p}{C_v} = \frac{1.000}{0.714} = 1.4$

Now,
$$\frac{423}{T_2} = \left(\frac{p_1}{p_2}\right)^{\frac{\gamma-1}{\gamma}} = \left(\frac{3.8}{1}\right)^{\frac{1.4-1}{1.4}} = 1.465$$

 \Rightarrow

$$T_2 = (P_2)$$

 $T_2 = 288.7 \,\mathrm{K}$

 $\forall_2 = 0.39 \,\mathrm{m}^3$

Again,

$$\frac{0.15}{\forall_2} = \left(\frac{p_2}{p_1}\right)^{\frac{1}{\gamma}} = \left(\frac{1}{3.8}\right)^{\frac{1}{1.4}} = 0.385$$

 \Rightarrow

Now let us find the temperature T_3 and volume \forall_3 after constant pressure heating. If *m* is the mass of gas contained in the system we know that gas constant

$$R = C_p - C_v = 0.286 \text{ kJ/kgK}$$

 $m = \frac{p_1 \forall_1}{RT_1} = \frac{3.8 \times 10^5 \times 0.15}{286 \times 423} = 0.47 \text{ kg}$

And

enthalpy
$$dH = mC_p(T_3 - T_2)$$

$$70 = 0.47 \times 1 (T_3 - 288.7)$$

$$\Rightarrow$$

$$T_3 = 437.6 \,\mathrm{K}$$

Since the heating is at constant pressure,

$$\forall_3 = \frac{T_3}{T_2} \forall_2 = \frac{437.6}{288.7} \times 0.39 = 0.59 \text{ m}^3$$

Work done during adiabatic expansion,

$$W_{1-2} = \frac{p_1 \forall_1 - p_2 \forall_2}{\gamma - 1} = \frac{(3.8 \times 0.15 - 1.0 \times 0.39) 10^5}{(1.4 - 1) 10^3} = 45 \text{ kJ}$$

And work done during constant pressure heating,

$$W_{2-3} = p_2(\forall_3 - \forall_2) = 1.0 \times 10^2 (0.59 - 0.39) = 20 \text{ kJ}$$

Hence, total work done during constant pressure heating, ı

$$W = W_{1-2} + W_{2-3} = 45 + 20 = 65 \text{ kJ}$$

EXAMPLE 5

0.336 m3 of gas at 10 bar pressure and 150°C temperature expands adiabatically, until its pressure is 4 bar. It is then compressed isothermally, to its original volume. Determine the final temperature and pressure of the gas. Also determine the change in internal energy. Take $C_p = 0.996$ kJ/kgK, $C_v = 0.703$ kJ/kgK.

SOLUTION

Here, $\forall_1 = \forall_3 = 0.336 \text{ m}^3$, $p_1 = 10 \text{ bar}$, $p_2 = 4.0 \text{ bar}$, $T_1 = 150 + 273 = 423 \text{ K}$

In figure process 1-2 represents the adiabatic expansion of the gas and the process 2-3 represents the isothermal compression to its original volume.

We know that adiabatic index,

$$\gamma = \frac{C_p}{C_v} = \frac{0.996}{0.703} = 1.417 \text{ and}$$
Gas constant
$$R = C_p - C_v = 0.293 \text{ kJ/kgK}$$
Now,
$$\frac{423}{T_2} = \left(\frac{p_1}{p_2}\right)^{\frac{\gamma-1}{\gamma}} = \left(\frac{10}{4}\right)^{\frac{1.417-1}{1.417}} = 1.31$$

$$\Rightarrow T_2 = 323 \text{ K}$$
Since the compression is isothermal from 2 to 3,
$$P_2 = \frac{1}{V_1 - V_2} = \frac{V_1}{V_2}$$

 $T_3 = T_2 = 323 \text{ K}$ therefore,

If final pressure of the gas is p_3 , for a constant volume process 3–1,

$$p_3 = \frac{T_3}{T_1} p_1 = \frac{323}{423} \times 10 = 7.6$$
 bar

Let us find the mass of the gas *m* and

$$m = \frac{p_1 \forall_1}{RT_1} = \frac{10 \times 10^5 \times 0.336}{293 \times 423} = 2.7 \text{ kg}$$

Change in internal energy

$$dU = U_3 - U_1 = mC_v(T_3 - T_1) = 2.7 \times 0.703(323 - 423) = -189.8 \text{ kJ}$$

The negative sign indicates that there is a decrease in internal energy.

EXAMPLE 6

A certain quality of air has a volume of 0.028 m³ at a pressure of 1025 bar and 25°C. It is compressed to a volume of 0.0042 m3 according to the law $p \forall^{1.3}$ = constant. Find the constant volume required to bring the air back to its initial temperature.

SOLUTION

Here, $\forall_1 = 0.208 \text{ m}^3$, $\forall_2 = 0.0042 \text{ m}^3$, $p_1 = 1.25 \text{ bar}$, $T_1 = 25 + 273 = 298 \text{ K}$.

Let at the end of compression the final temperature = T_{2} , the final pressure = p_{2} and pressure at a constant volume required to bring the air back to its initial temperature is p_3 .

Now,

$$\frac{298}{T_2} = \left(\frac{\forall_2}{\forall_1}\right)^{n-1} = \left(\frac{0.0042}{0.028}\right)^{1.3-1} = 0.566$$

$$\Rightarrow \qquad T_1 = 526.5 \,\mathrm{K}$$

So,

 \Rightarrow

$$\frac{1.25}{p_2} = \left(\frac{\forall_2}{\forall_1}\right)^n = \left(\frac{0.0042}{0.028}\right)^{1.3} = 0.085$$

$$p_2 = 14.7 \text{ bar}$$

The work done during compression

$$W_{1-2} = \frac{p_2 \forall_2 - p_1 \forall_1}{n-1} = \frac{(0.0042 \times 14.7 - 1.25 \times 0.028) \times 10^5}{(1.3 - 1)10^3} = 8.913 \text{ kJ}$$

For a constant volume process 2–3, maintaining given condition, $T_3 = T_1 = 298$ K

. . .

$$p_3 = \frac{T_3}{T_2} p_2 = \frac{298}{526.5} \times 14.7 = 8.32$$
 bar

• • •

EXAMPLE 7

The properties of a closed system change following the relation between pressure and volume as $p \forall = 3.0$, where p is expressed in bar. Calculate the work done when the pressure increases from 1.5 bar to 7.5 bar.

SOLUTION Here.

So,

$$p_1 = 1.5 \text{ bar}, \quad p_2 = 7.5 \text{ bar}$$

 $\forall_1 = \frac{3}{1.5} = 2 \text{ m}^3 \quad \text{and} \quad \forall_2 = \frac{3}{7.5} = 0.4 \text{ m}^3$
 $W_1 = \int_{-\infty}^{\sqrt{2}} p d\forall = 10^5 \int_{-\infty}^{\sqrt{2}} \frac{3}{-2} d\forall = 10^5 [\ln \forall]^{0.4} \times 3$

Hence, work done

$$W_{1-2} = \int_{\forall_1}^{\forall_2} p d\,\forall = 10^5 \int_{\forall_1}^{\forall_2} \frac{3}{\forall} d\,\forall = 10^5 \left[\ln\forall\right]_2^{0.4} \times 10^5 \,\mathrm{J} = -4.83 \times 10^5 \,\mathrm{J} = -4.83 \,\mathrm{kJ}$$



 $p = C/\forall^2$

¥,

EXAMPLE 8

To a closed system 150 kJ of work is supplied. Let the initial volume is 0.6 m³ and pressure of the system changes as $p = (8 - 4\forall)$, where p is in bar. Determine the final volume and pressure of the system. Solution

SOLUTION

Here, Hence, work done

$$W_{1-2} = \int_{\forall_1}^{\forall_2} p d\forall = 10^5 \int_{\forall_1}^{\forall_2} (8 - 4\forall) d\forall$$
$$-150 \times 10^3 = 10^5 \left[8\forall - 2\forall^2 \right]_{0.6}^{\forall_2}$$
$$= 10^5 \left[8\forall_2 - 2\forall_2^2 - 4.08 \right]$$
$$8\forall_2 + 2.58 = 0$$

 $\forall_1 = 0.6 \,\mathrm{m}^3, W_{1-2} = -150 \,\mathrm{kJ}$

 \Rightarrow

 \Rightarrow

 \Rightarrow

 $\forall_2 = 0.354$ $p_2 = 8 - 4 \times 0.354 = 6.584$ bar

 $2\forall_2^2 - 8\forall_2$

So, final pressure

EXAMPLE 9

A fluid at a pressure of 3bar and with specific volume of 0.18 m³/kg contained in a cylinder behind a piston expands reversibly to a pressure of 0.6 bar according to a law $p = \frac{C}{\forall^2}$, where C is a constant. Calculate the work done by the fluid on the piston.

SOLUTION

Here,
$$\forall_1 = 0.18 \text{ m}^3/\text{kg}, p_1 = 3 \text{ bar}, p_2 = 0.6 \text{ bar}$$

So,

$$C = p_1 \forall_1^2 = 3(0.18)^2 \text{ bar.} (\text{m}^3/\text{kg})^2 = 0.0972 \text{ bar.} (\text{m}^3/\text{kg})^2$$

And

$$\forall_2 = \sqrt{\frac{C}{p_2}} = \sqrt{\frac{0.0972}{0.6}} = 0.402 \text{ m}^3\text{kg}$$

Hence, work done

$$W_{1-2} = \int_{\forall_1}^{\forall_2} p d \forall = \int_{\forall_1}^{\forall_2} \frac{C}{\forall^2} d \forall = -C \left[\frac{1}{\forall}\right]_{\forall_1}^{\forall_2}$$
$$= -0.0972 \times 10^5 \left[\frac{1}{0.402} - \frac{1}{0.18}\right] = 29820.9 \text{ N·m/kg} \cdot 10^{10} \text{ N/m/kg}$$

EXAMPLE 10

A cylinder contains 1 kg of a certain fluid at an initial pressure of 20 bar. The fluid is allowed to expand reversibly behind a piston according to a law $p\nabla^2$ = constant, until the volume is doubled. The fluid is then cooled reversibly at constant pressure until the piston regains its original position. Heat is then supplied reversibly with the piston firmly locked in position until the pressure rises to the original value of 20 bar. Calculate the net work done by the fluid for an initial volume of 0.05 m³.



SOLUTION

Here,

$$m = 1 \text{ kg}, p_1 = 20 \text{ bar}, \forall_1 = 0.05 \text{ m}^3, \forall_2 = 2 \times 0.05 = 0.1 \text{ m}^3$$

Now,

$$p_2 = \left(\frac{\forall_1}{\forall_2}\right)^2 p_1 = \left(\frac{\forall_1}{2\forall_1}\right)^2 20 = 5 \text{ bar}$$

Work done by the fluid from 1 to 2.

$$W_{1-2} = \int_{\forall_1}^{\forall_2} p d\forall = \int_{\forall_1}^{\forall_2} \frac{C}{\forall^2} d\forall = -C \left[\frac{1}{\forall}\right]_{\forall_1}^{\forall_2} = -p_1 \forall_1^2 \left[\frac{1}{\forall}\right]_{\forall_1}^{\forall_2}$$
$$= -20 \times 10^5 \times 0.0025 \left[\frac{1}{0.1} - \frac{1}{0.05}\right] = 50000 \text{ N} \cdot \text{m} = 50 \text{ kNm}$$

And work done during constant pressure process, from 2 to 3

$$W_{2-3} = p_2(\forall_2 - \forall_3) = 5 \times 10^5 (0.1 - 0.05) = 25000 \text{ N} \cdot \text{m} = 25 \text{ kNm}$$

In constant volume process from 3 to 1, $W_{2-3} = 0$. Net work done by the fluid,

$$W = W_{1-2} - W_{2-3} = 50 - 25 = 25 \text{ kN} \cdot \text{m}$$

EXAMPLE 11(a)

A gas is enclosed in a piston cylinder assembly as a system. The gas is initially at a pressure of 500 kPa and occupies a volume of 0.2 m³. The gas is taken to the final state

where $p_2 = 100$ kPa by following two different processes. Calculate the work done by the gas in each case when: (a) the volume of the gas is inversely proportional to the pressure and (b) the process follows the

path $p \forall^{\gamma} = \text{constant}$, where $\gamma = 1.4$.

SOLUTION

(a) The path followed by the system $\forall \propto \frac{1}{p}$ i.e., $p\forall$ = constant



So,

Hence work done,

$$W_{1-2} = \int_{\forall_1}^{\forall_2} p d \forall = p_1 \forall_1 \int_{\forall_1}^{\forall_2} \frac{d \forall}{\forall}$$
$$= p_1 \forall_1 \ln\left(\frac{\forall_2}{\forall_1}\right)$$
$$= p_1 \forall_1 \ln\left(\frac{p_1}{p_2}\right)$$
$$= (500 \times 10^3 \times 0.2 \ln 5) J$$
$$= 160.94 \text{ kJ}$$

(b) The path followed by the system, $p_1 \forall_1^{\gamma} = p_2 \forall_2^{\gamma}$.

So,
$$\left(\frac{\forall_2}{\forall_1}\right)^{\gamma} = \frac{p_1}{p_2}$$

and

$$\forall_2 = \left(\frac{p_1}{p_2}\right)^{1/\gamma} \forall_1 = \left(\frac{500}{100}\right)^{1/1.4} \times 0.2 = 0.6313 \text{ m}^3$$

Hence work done,

$$W_{1-2} = \frac{p_1 \forall_1 - p_2 \forall_2}{\gamma - 1} = \frac{(500 \times 0.2 - 100 \times 0.6313)10^3}{(1.4 - 1) \times 10^3} = 92.175 \text{ kJ}$$

EXAMPLE 11(b)

Determine the total work done by a gas system following an expansion process as shown in the figure. Initial pressure of the system is (50 E + 2) kPa



SOLUTION

Being a constant pressure process,

$$W_{A-B} = 50 \times 100 (0.4 - 0.2) = 1000 \text{ kJ}$$

 $p_1 \forall_1^{\gamma} = p_2 \forall_2^{\gamma}$

Again,

Therefore,	$p_2 = \left(\frac{\forall_1}{\forall_2}\right)^{\gamma} p_1 = \left(\frac{0.4}{0.8}\right)^{1.3} \times 50 \times 100 = 2030.63 \text{ kPa}$
Again,	$W_{B-C} = \frac{p_1 \forall_1 - p_2 \forall_2}{\gamma - 1} = \frac{(50 \times 100 \times 0.4 - 2030.63 \times 0.8)}{1.3 - 1} = 1251.65 \text{ kJ}$
Hence, total workdone	$W_{A-C} = W_{A-B} + W_{B-C} = 1000 + 1251.65 = 2251.65 \text{ kJ}$
EXAMPLE 12	

A piston and cylinder arrangement containing a fluid system has a stirring device as shown in the figure. The piston is frictionless and it is held down against the fluid due to atmospheric pressure of 101.3 kPa.The stirring device is turned 9500 revolutions with an average torque against the fluid of $1.25 \text{ N} \cdot \text{m}$. Meanwhile the piston of 0.65 m diameter moves out 0.6 m. Determine the net work transfer for the system.



SOLUTION

Here, work done by the stirring device upon the system

$$W_1 = 2\pi NT = 2\pi \times 9500 \times 1.25 = 74612.83 \text{ N} \cdot \text{m} = 74.613 \text{ kJ}$$

This is a negative work

Work done by the system upon the surroundings

$$W_2 = pAL = 101.3 \times \left(\frac{\pi}{4} \times 0.65^2\right) \times 0.6 = 20.168 \text{ kJ}$$

Hence, net work transfer for the system.

$$= -74.613 + 20.168 = -54.445 \text{ kJ}$$

EXAMPLE 13

A piston and cylinder arrangement containing a fluid system has a stirring device in the cylinder. The piston is frictionless and it is held down against the fluid due to the atmospheric pressure of 101.325 kPa. The stirring device is turned 10000 revolutions with an average torque against the fluid of 1.275 Nm. Meanwhile the piston of 0.6 m diameter moves out 0.8 m. Compute the net work transfer for the system.

SOLUTION

Here, work done by the stirring device upon the system

$$W_1 = 2\pi NT = 2\pi \times 10000 \times 1.275 = 80110.61 \text{ Nm} = 80.11 \text{ kJ}$$

It is a negative work

Work done by the system upon the surroundings

$$W_2 = pAL = 101.325 \times \left(\frac{\pi}{4} \times 0.6^2\right) \times 0.8 = 22.92 \text{ kJ}$$

Hence, net work transfer for the system.

$$= -80.11 + 22.92 = -57.19 \text{ kJ}$$

EXAMPLE 14

A balloon which is initially collapsed is slowly inflated with helium from a cylinder forming the balloon into a sphere of 5 m in diameter. The atmospheric pressure is 100 kPa. During the filling process the helium inside the cylinder remains constant at 300 K all the time. Determine the work by the cylinder-balloon system.



SOLUTION

Here displacement work

$$W_D = \left(\int pd\forall\right)_{\text{balloon}} + \left(\int pd\forall\right)_{\text{bottle}}$$
$$= p\left(\forall_2 - \forall_1\right) + 0 = 100 \times \frac{4}{3}\pi \left(\frac{5}{2}\right)^3$$
$$= 6544.98 \text{ kJ} = 6.545 \text{ MJ}$$

EXAMPLE 15

A spherical balloon of 1 m diameter contains a gas at 150 kPa. The gas inside the balloon is heated until the pressure increases to 450 kPa. During the process of heating, pressure of the gas inside the balloon is proportional to the cube of the diameter of the balloon. Determine the work done by the gas inside the balloon.

SOLUTION

Here, $p \propto D^3 \implies p = C \cdot D^3$ Substituting p = 150 kPa and D = 1 m, we have

$$C = \frac{p}{D^3} = \frac{150}{1^3} = 150 \text{ kPa/m}^3$$

Properties and Thermodynamic Processes of Gas

When	$p = 450 \mathrm{kPa}$
	$D = \sqrt[3]{\frac{450}{150}} = 1.4422 \text{ m}$
Now we can consider,	$D_1 = 1 \mathrm{m}, D_2 = 1.4422 \mathrm{m}$
Again,	$\forall = \frac{4}{3} \pi \left(\frac{D}{2}\right)^3 = \frac{\pi D^3}{6}$
So,	$d\forall = \frac{\pi}{6} \times 3D^2 \times dD = \frac{\pi D^2}{2} \times dD$
Hence work done	
	$W = \int_{1}^{2} p d \forall = \int_{D_{1}}^{D_{2}} CD^{3} \cdot \frac{\pi D^{2}}{2} dD = \frac{\pi C}{2} \int_{D_{1}}^{D_{2}} D^{5} dD$
	$= \frac{\pi C}{2} \left[\frac{D^6}{6} \right]_{1}^{1.4417}$
	$= \frac{\pi}{2} \cdot \frac{150}{6} [1.4422^6 - 1^6]$
	= 314.09 kJ
	EXERCISE

- 1. Write the characteristic equation of a perfect gas, name symbols and state their unit in S.I. Distinguish between specific gas constant and universal gas constant.
- 2. Derive an expression for heat transfer during polytropic process.
- 3. Establish a relation between the two specific heat of gas and specific gas constant.
- 4. Distinguish between gas and vapour.
- 5. What is perfect gas and what is real gas? Under what condition does a real gas behave as a perfect gas?
- 6. Establish a relation amongst p, \forall, T in adiabatic process.
- 7. What do you understand by enthalpy? Show that for a constant pressure, the heat supplied to the gas is equal to change of enthalpy.
- 8. What do you understand by a thermodynamic process?
- 9. Explain the difference between non flow process and a flow process. Derive the equation for work done during a non-flow process.
- 10. State the law of corresponding states.

- 11. Define the compressibility factor and explain its significance.
- 12. A quantity of gas is compressed from 130 kPa and 38°C to 1.25 MPa and 280°C in an engine cylinder. Find the law of the process and determine per kg of the gas the work transfer. Take R = 0.287 kJ/kg for the gas.

Ans.
$$[p \forall^{1.299} = C, W = -238.28 \text{ kJ/kg}]$$

- 13. A certain gas has $C_p = 1.9681.507 \text{ kJ/kgK}$ and $C_v = 1.507 \text{ kJ/kgK}$. Find its molecular weight and the gas constant.
- 14. A constant volume chamber of 0.3 m³ capacity contain 2 kg of this gas at 5°C. Heat is transferred to the gas until the temperature becomes 100°C. Find the work done, the heat transferred, and the changes in internal energy, enthalpy and entropy.
- 15. 0.2 m^3 of mixture of fuel and air at 1.2 bar and 60°C is compressed until it's pressure becomes 12 bar and temperature becomes 270°C, then it is ignited suddenly at constant volume and its pressure becomes twice the pressure at the end of compression. Calculate the maximum temperature reached and change in internal energy. Also compute the heat transfer during the compression process. Consider mixture as a perfect gas and take = 1.072 kJ/kgK and R = 294 J/kgK.

Ans. [813°C, 143.5 kJ, -16 kJ]

16. An ideal gas at 30° C and 1 bar is compressed adiabatically from 5 m³ to 1 m³. Find the temperature, pressure and work done. Take = 1.4

Ans. [304°C, 9.5 bar, 1.125 kJ]

17. A cylinder contains 0.084 m³ of hydrogen at 7.05 bar and 18°C. It is then compressed adiabatically to 14 bar and then expanded isothermally to the original volume of 0.084 m³. Characteristic constant for hydrogen is 4200 J/kgK and its specific heat at constant pressure is 14.28 kJ/kgK. Determine the final pressure of the gas and the amount of heat which must be added to the gas during isothermal expansion. Also calculate the heat which must be substracted from the gas after expansion in order to reduce it to its initial state of stressor.

Ans. [2.25 bar, 34.5 kJ, 24.14 kJ]



FIRST LAW OF THERMODYNAMICS AND ITS APPLICATION

4.1 INTRODUCTION

When a thermodynamic system undergoes a process or a combination of processes, as a result of which the system can be restored to its initial state, then the system is said to have undergone a *cyclic process* or a cyclic change. It is seen from figure 4.1 that a *thermodynamic cycle* may be the combination of any number of processes. But the minimum number of processes required to complete a cycle is two as shown in figure 4.1(c), where process 1 to 2 and 2 to 1 comprising the cycle.

4.2 JOULE'S EXPERIMENT

Joule carried out a series of experiment during 1843 and 1848 that subsequently led to the formulation of the *first law of thermodynamics*. Of the several experiments conducted by Joule, the *paddle wheel experiment* described below is considered as a classical one.

This experiment consisted of two-process cycle arrived out on a system comprising of a fluid as shown in figure 4.2., where work W was done on the system by means of a paddle wheel. The paddle wheel was made to rotate by lowering the mass which is tied to a string and the string was wound on the paddle wheel as shown in Fig. 4.2 (a). While lowering the mass m through a distance Z, work was done on liquid (the system) and the amount of work done on the system was measured. This caused a rise in temperature of fluid and the system has reached a different







Fig. 4.2 (a)

state. The system was then placed in contact with a water bath as shown in figure 4.2. Heat was thus transferred from the fluid to the water in process 2-1, until the original state of the fluid was restored, as indicated by its temperature. Heat Q was measured from the increase in energy of water bath.



Joule carried out several experiments involving different types of work interaction. Based on these experimental results Joule found that the *net work done on the system is exactly equal to the net heat interaction*, irrespective of the type of the work interaction, the rate at which work was done and the method employed for transferring the energy in the form of heat from the system. That is

 $(\oint \delta Q)_{\text{from the system}} = (\oint \delta W)_{\text{on the system}}$ or $\oint \delta Q - \oint \delta W = 0$. Here \oint denotes integral over a cycle or

cyclic integral.

4.3 FIRST LAW OF THERMODYNAMICS FOR A SYSTEM UNDERGOING A THERMODYNAMICS CYCLE

The conclusion from the Joule's experiments leads to the statement of the first law of thermodynamics as follows:

When a closed system undergoes a thermodynamic cycle, during a cycle the cyclic integral of heat transfer is equal to the cyclic integral of work transfers.

Mathematically
$$\oint \delta Q = \oint \delta W$$
 (4.1)

Here denotes integral over a cycle or cyclic integral \oint has been used to indicate the both Q and W are inexact differentials.

When heat Q and work W are in different units, then a *proportionality factor*, called the *mechanical equivalent of heat*, denoted by a symbol J is required. So, the first law may be stated as follows, Whenever a system undergoes a cyclic change, the algebraic sum of work transfers is proportional to the algebraic sum of heat transfers or work and heat are mutually convertible one into other.

Mathematically,
$$J \oint \delta Q = \oint \delta W$$
 (4.2)

104

In M.K.S unit J = 427 kgfm/kcal, If SI is used for heat and work, then J = 1. So it is not necessary to use any proportionality factor in the first law equation.

4.4 THE IMPORTANT CONSEQUENCES OF THE FIRST LAW OF THERMODYNAMICS

The first law of thermodynamics leads to the following important consequences.

- 1. Heat interaction is a path function.
- 2. Energy is a property of a thermodynamic system.
- 3. The energy of an isolated system is conserved.
- 4. *A perpetual motion machine of the first kind is impossible.*



Fig. 4.3

4.4.1 Heat Interaction is a Path Function

Equation (4.1) states the first law for a thermodynamic cycle. But many times we are concerned with heat and work interaction of a system during a single thermodynamic process. We, therefore need the formulation of first law for a process. Referring figure 4.3, suppose a system follows the path 1-*A*-2 in reaching the final state 2 starting from the initial state 1. Then the system can be restored to the initial state 1 either by the path 2-*B*-1 or 2-*C*-1. Then the combination of the processes leads to two different cycles, namely 1-*A*-2-*B*-1 and 1-*A*-2*C*-1. Applying the first law thermodynamics to the cycles 1-*A*-2-*B*-1 and 1-*A*-2-*C*-1 we get:

$$\int_{1-A-2} \delta Q + \int_{2-B-1} \delta Q - \int_{1-A-2} \delta W - \int_{2-B-1} \delta W = 0$$
(4.3)

and

$$\int_{-A-2} \delta Q + \int_{2-C-1} \delta Q - \int_{1-A-2} \delta W - \int_{2-C-1} \delta W = 0$$
(4.4)

Subtracting equation (4.4) from equation (4.3), we get

$$\int_{2-B-1} \delta Q - \int_{2-C-1} \delta Q - \int_{2-B-1} \delta W + \int_{2-C-1} \delta W = 0$$

or,
$$\int_{2-B-1} \delta Q - \int_{2-C-1} \delta Q = \int_{2-B-1} \delta W - \int_{2-C-1} \delta W = 0$$
(4.5)

We know that work interaction is a path function and $\int \delta W$ represents the area under the $p - \forall$ curve. Hence,

$$\int_{-B-1} \delta W - \int_{2-C-1} \delta W \neq 0 \tag{4.6}$$

So, from equation (4.5) and (4.6) we get,

2.

$$\int_{2-B-1} \delta Q - \int_{2-C-1} \delta Q \neq 0 \text{ or } \int_{2-B-1} \delta Q \neq \int_{2-C-1} \delta Q$$
(4.7)

That is, the heat interaction along the path 2-*B*-1 is different from the heat interaction along the path 2-*C*-1. Thus, heat interaction depends on the path followed by a system and *heat interaction is a path function*. It is not a point function or state function and hence it's differential is not exact. Heat interaction is not a property of a thermodynamic system.

4.4.2 Energy is a Property of a Thermodynamic System

For the purpose, consider the change of state of a system from 1 to 2, along path A and then from 2 to 1 along path B and C, thus completing two different cycles, namely 1-A-2-B-1 and 1-A-2-C-1, as shown in figure 4.3. According to the first law of thermodynamics for both the cycle 1-A-2-B-1 and 1-A-2-C-1 we get from equation (4.5),

$$\int_{2-B-1} (\delta Q - \delta W) = -\int_{2-C-1} (\delta Q - \delta W)$$
(4.8)

We know that $\int \delta Q$ and $\int \delta W$ each depends on the path followed by a system. Equation (4.8) shows that the value of $\int (\delta Q - \delta W)$ connecting the states 1 and 2 is the same for the processes 2-*B*-1 and 2-*C*-1. In other words the quantity $(\delta Q - \delta W)$ does not depend on the path followed by a system, but depends only on the initial and the final states of a system. Hence, the quantity $(\delta Q - \delta W)$ is an exact differential, thus is a property of a system. Thus, property is called the energy *E*, and then the differential change in the energy of a system is given by

$$dE = \delta Q - \delta W \tag{4.9}$$

Here dE is an exact differential. That is, there exits a property called the energy of the system and its differential change is related to the work and heat interactions. We also know that the energy of a system is sum of macroscopic and microscopic modes of energy and is given by

$$E = KE + PE + U \tag{4.10}$$

Where KE = Kinetic Energy, PE = Potential Energy, and U = Internal Energy. Comparing equations (4.9) and (4.10), we get

$$dE = d(KE) + d(PE) + dU = \delta Q - \delta W$$

So, change in the value of energy is the algebraic sum of the heat supply and the work done during any change in state. This is called the *corollary* 1 of the first law of thermodynamics.

4.4.3 The Energy of an Isolated System is Conserved (Law of the Conservation of Energy)

The energy can neither be created nor destroyed though it can be transformed from one form to another. *In an isolated system, the energy of the system remains constant.* According to this law, when a system undergoes a change of state, then both heat transfer and work transfer take place. The net energy transfer is stored within the system and is known as *stored energy* or *total energy* of the system. Mathematically, $dE = \delta Q - \delta W$. On integrating this expression for a change of state from 1 to 2, we get,

$$Q_{1-2} - W_{1-2} = E_2 - E_1 \tag{4.11}$$

When Q_{1-2} = heat transferred to the system during the process from state 1 to 2, W_{1-2} = work done by the system on the surrounding during the process,

$$E_1$$
 = total energy of the system at state 1 = $PE_1 + KE_1 + U_1 = mgZ_1 + \frac{mV_1^2}{2} + U_1$,

 E_2 = total energy of the system at state 2 = $PE_2 + KE_2 + U_2 = mgZ_2 + \frac{mV_2^2}{2} + U_2$.

Thus, the above equation (4.11) may be re-written as

$$Q_{1-2} - W_{1-2} = E_2 - E_1$$

= $(PE_2 + KE_2 + U_2) - (PE_1 + KE_1 + U_1)$
= $m(gZ_2 - gZ_1) + m\left(\frac{V_2^2}{2} - \frac{V_1^2}{2}\right) + (U_2 - U_1)$ (4.12)

For unit mass this expression gets converted to

$$q_{1-2} - w_{1-2} = g\left(Z_2 - Z_1\right) + \left(\frac{V_2^2}{2} - \frac{V_1^2}{2}\right) + \left(u_2 - u_1\right)$$
(4.13)

Case (1): When there is no change in potential energy of the system *i.e.*, when the height of the system from the datum level is same $(Z_2 = Z_1)$, then $PE_2 = PE_1$. Thus, the equation (4.12) will be

$$Q_{1-2} - W_{1-2} = \left(\frac{V_2^2}{2} - \frac{V_1^2}{2}\right) + \left(U_2 - U_1\right)$$
(4.14)

Case (2): When there is no change in potential energy and also there is no flow of the mass into or out of the system, then $PE_2 = PE_1$ and $KE_2 = KE_1$. Thus, the equation (4.12) takes the form

$$Q_{1-2} - W_{1-2} = (U_2 - U_1) = dU$$
(4.15)

In other words, in a closed or no-flow thermodynamic system, PE = 0 and KE = 0. Thus the equation (4.15) is known as *no-flow energy equation*.

Case (3): For an isolated system for which $Q_{1-2} = 0$ and $W_{1-2} = 0$, the above equation (4.11) becomes, $E_2 - E_1 = 0 \Rightarrow E_2 = E_1$. So, the total energy of the system at state 2 equals with total energy at state 1, hence the first law of thermodynamics is sometimes designated as *the law of conservation of energy*.

4.4.4 Perpetual Motion Machine of the First Kind is Impossible

A perpetual motion machine is defined as a machine which produces work energy without consuming an equivalent amount of heat energy from other source or a machine which always receives heat from source without producing any work. It is shown in figure 4.4. So, a machine which violates the first law of thermodynamics *i.e.*, energy can neither be created nor destroyed, but can be transformed from one form to another, is known as *perpetual motion machine of the first kind* (PMM-1). But in actual practice such machine is impossible to obtain, because no machine can produce energy of its own without consuming any other form of energy.



4.5 LIMITATIONS OF FIRST LAW OF THERMODYNAMICS

The following are the limitations of the first law of thermodynamics.

- (a) When a closed system undergoes a thermodynamic cycle, the net heat transfer is equal to the net work transfer.
 - (i) This statement does not specify the direction of flow of heat and work, *i.e.*, whether the heat flows from a hot body to a cold body or from a cold body to a hot body.
 - (ii) It also does not give any condition under which these transfer takes place.
- (b) The heat energy and mechanical work are mutually convertible. Though the mechanical work can be fully converted into heat energy, but only a part of heat energy can be converted into mechanical work. This means that the heat energy and mechanical work are not fully mutually convertible. In other words, there is a limitation on the conversion of one form of energy into another form.

4.6 APPLICATION OF FIRST LAW THERMODYNAMICS TO A NON-FLOW PROCESS

When a system undergoes a change of state or a thermodynamic process then both the heat transfer and work transfer takes place. The net energy transfer is stored within the system and is known as stored energy or total energy of the system and which was discussed in equation (4.11).

$$Q_{1-2} - W_{1-2} = \int_{1}^{2} dE = E_2 - E_1$$
(4.16)

For non-flow process the stored energy is the internal energy only. Thus, equation (4.16) when applied to a non-flow process or a static system may be written as

$$Q_{1-2} - W_{1-2} = \int_{1}^{2} dU = U_2 - U_1$$
(4.17)

It may be noted that heat and work are not properties of the system, but their difference $(Q_{1-2} - W_{1-2})$ during a process is the numerical equivalent of stored energy. Since the stored energy is a property, therefore $(Q_{1-2} - W_{1-2})$ is also a property.

4.6.1 Types of Non-Flow Process

Non-flow processes are of two types

- (a) Reversible non-flow process: All the different reversible non-flow processes as applicable for perfect gas are discussed in detail in the previous chapter.
- (b) Irreversible non-flow process (free or unresisted expansion process): The free expansion process is an irreversible process. A free expansion occurs when a fluid is allowed to expand suddenly into a vacuum chamber through an orifice of large dimension.

Let us consider the insulated container is divided into two chambers A and B by a partition as shown in figure 4.5(a). The chamber A contains a perfect gas having volume \forall_1 , pressure p_1 and temperature T_1 while the chamber B is completely vacuum. As these chambers are insulated, no heat transfer takes

place from its surroundings. Now, if the partition is removed, the gas will expand freely and occupy the whole space as shown in figure 4.5(b). By this, the volume of gas increases to \forall_2 ,

pressure decreases to p_2 and the temperature may also decrease to T_2 . Since there is no expansion of the boundary of the system because it is rigid, therefore no work is done. Thus for a free expansion process of gas following points may be noted:







Fig. 4.5 (b)

- 1. Since the system is perfectly insulated, no heat transfer takes place. So, $Q_{1-2} = 0$, therefore, the expansion of the gas may be called as an *adiabatic expansion*.
- Since the free expansion of the gas from the equilibrium state 1 to the equilibrium of state 2 takes place suddenly, therefore the intermediate states out not be in equilibrium states, as shown on the *p*−∀ diagram in figure 4.5(c), thus the process is reversible and expansion is therefore known as is *reversible adiabatic expansion*.
- 3. Since there is no resistance to overcome during free expansion process, therefore no work is done by the system, *i.e.*, $W_{1-2} = 0$, thus the free expansion process is also known as *unresisted expansion process*.



4. According to the first law of thermodynamics, $Q_{1-2} = W_{1-2} + \int_{1}^{2} dU$ (4.18)

Since for free expansion, $Q_{1-2} = 0$ and $W_{1-2} = 0$, therefore the change in internal energy,

$$\int_{1}^{2} dU = U_{2} - U_{1} = 0 \implies U_{2} = U_{1}$$
(4.19)

In other words, the internal energy of the system in a free expansion process remains constant. Thus, the free expansion process is also known as *constant internal energy process*.

1. We know the change in internal energy $\int_{1}^{2} dU = mC_{\nu}(T_{2} - T_{1})$ (4.20)

$$\int_{1}^{2} dU = 0, \ mC_{\nu}(T_2 - T_1) = 0, \ \text{which leads to } T_2 = T_1$$
(4.21)

So, there is no change in temperature of the system. In other words no temperature of the system, in a free expansion process remains constant. But *it cannot be called an isothermal*

process, because in an actual isothermal process, work done by the gas is $W_{1-2} = p_1 \forall_1 \ln\left(\frac{\forall_2}{\forall_1}\right)$ but

in free expansion process work done. $W_{1-2} = 0$.

Since

2. We know that the change in enthalpy

$$\int_{1}^{2} dT = H_2 - H_1 = mC_p \int_{1}^{2} dT = mC_p (T_2 - T_1)$$
(4.22)

Since,
$$\int_{1}^{2} dT = 0$$
, $mC_{p}(T_{2} - T_{1}) = 0$, this leads to $H_{2} = H_{1}$ (4.23)

In other words, the enthalpy of the system in a free expansion process remains constant. Thus, the free expansion process may also be called *constant enthalpy process*.

APPLICATION OF FIRST LAW OF THERMODYNAMIC TO A STEADY FLOW 4.7 PROCESS

The open systems which permit the transfer of mass to and from the system are known as *flow* process. In this process, the mass (working substance) enters into the system and leaves after doing the work. The flow process may be steady flow process or unsteady flow process.

In steady flow process the rate of mass flow at inlet and outlet is the same, the rate of heat transfer and the rate of work transfer is the same, the state of working substance at any point within the system is the same at all time and there is no change in the chemical composition of the system. Any thermodynamic property will have a fixed value at a particular location and will not alter with time. The property may vary along space co-ordinates (x, y, z), but invariant with time, *i.e.*,

 $\delta(Property)$ = 0. Therefore if the rate of flow of mass and energy across the control surface

of control volume are constant with time then the flow is called steady flow. If the rate of flow of mass and energy across the control surface of control volume are not constant with time then the flow is called unsteady flow.

4.7.1 Steady Flow Energy Equation (S.F.E.E)

Consider an open system through which the working substance flows at a steady rate, as shown in figure 4.6. The working substance enters into

the system at section 1 and leaves the system at section 2 and, there is no accumulation of mass or energy within the control volume of the system. Let

- A_1, A_2 = cross sectional area at section 1 and 2 in m²
- p_1, p_2 = pressure of fluid at section 1 and 2 in N/m^2
- v_{s1} , v_{s2} = specific volume of fluid at section 1 and 2 in m^3/kg
- u_1, u_2 = specific internal energy at section 1 and 2 in J/kg



Fig. 4.6

- V_1 , V_2 = velocity of fluid at section 1 and 2 in m/s
- Z_1, Z_2 = height above datum level at section 1 and 2 in m
- q_{1-2} = heat supplied to the system in J/kg of fluid w_{1-2} = work delivered by the system in J/kg of fluid.

Consider 1 kg of mass of the fluid (working substance). We know that total energy (J/kg) entering the system per kg of working fluid at section 1 is e_1

So,
$$e_1 = u_1 + p_1 v_{s1} + \frac{V_1^2}{2} + gZ_1 + q_{1-2}$$
 (4.24)

Here flow energy or displacement energy or flow work is the energy required to flow or move the working fluid against its pressure. Flow work or energy

$$FE = (p_1 A_1) x_1 = p_1 (A_1 x_1) = p_1 v_{s1}$$
(4.25)

Similarly, total energy (J/kg) leaving the system per kg of working fluid at section 2 is e_2

So,
$$e_2 = u_2 + p_2 v_{s2} + \frac{V_2^2}{2} + gZ_2 + w_{1-2}$$
 (4.26)

Assuming no loss of energy during flow, then according to law of conservation of energy

$$\Rightarrow \qquad u_1 + p_1 v_{s1} + \frac{V_1^2}{2} + gZ_1 + q_{1-2} = u_2 + p_2 v_{s2} + \frac{V_2^2}{2} + gZ_2 + w_{1-2}$$
(4.27)

Again we know that, specific enthalpy of working fluid $h = u + pv_s$ Thus, equation (4.27) may be re-written as

$$h_{1} + \frac{V_{1}^{2}}{2} + gZ_{1} + q_{1-2} = h_{2} + \frac{V_{2}^{2}}{2} + gZ_{2} + w_{1-2}$$
(4.28)

Otherwise,

 $e_1 = e_2$

$$h_1 + ke_1 + pe_1 + q_{1-2} = h_2 + ke_2 + pe_2 + w_{1-2}$$
(4.29)

It may be noted that all the terms in equation (4.28) represent the energy flow per unit mass of the working fluid, *i.e.*, in J/kg. If m_1 , m_2 be the mass flow rate of working fluid in kg/s then, for a steady flow, $m_1 = m_2 = m$, *i.e.*, mass flow rate at section 1 equals with mass flow rate at section 2. When the equation (4.28) is multiplied through by the mass (*m*) in kg/s, then all the term will represent the energy flow per unit time, *i.e.*, in J/s.

Thus the equation (4.28) may be re-written as

$$m\left(h_{1} + \frac{V_{1}^{2}}{2} + gZ_{1} + q_{1-2}\right) = m\left(h_{2} + \frac{V_{2}^{2}}{2} + gZ_{2} + w_{1-2}\right)$$
(4.30)

Both the equations (4.28) and (4.30) are known as steady flow energy equation (S.F.E.E.). The steady flow energy equation (4.28) for unit mass flow may be rearranged as

$$q_{1-2} - w_{1-2} = (h_2 - h_1) + \left(\frac{V_2^2}{2} - \frac{V_1^2}{2}\right) + g(Z_2 - Z_1)$$
$$= (h_2 - h_1) + (ke_2 - ke_1) + (pe_2 - pe_1)$$
(4.31)

In differential form, S.F.E.E. takes the shape as

$$\delta q - \delta w = dh + d(ke) + d(pe) \tag{4.32}$$

In thermodynamics, the effect of gravity is generally neglected. Therefore equation (4.31) may be re-written as

$$q_{1-2} - w_{1-2} = \left(h_2 - h_1\right) + \left(\frac{V_2^2}{2} - \frac{V_1^2}{2}\right)$$
(4.33)

Again if $V_1 = V_2$, then equation (4.33) will be reduced to

$$q_{1-2} - w_{1-2} = (h_2 - h_1) \tag{4.34}$$

In non-flow process the flow work at inlet and outlet is zero, *i.e.*, $p_1v_{s1} = 0 = p_2v_{s2}$. Therefore, $h_1 = u_1$ and $h_1 = u_2$.

Thus, the equation (4.34) will be
$$q_{1-2} - w_{1-2} = (u_2 - u_1)$$
 (4.35)

4.8 MASS BALANCE (CONTINUITY) EQUATION

In steady flow the mass flow rate (m) of the working fluid entering and leaving the system at sections 1 and 2 is given by,

$$m = \frac{A_1 V_1}{v_{s1}} = \frac{A_2 V_2}{v_{s2}}$$
(4.36)

This equation is known as equation of continuity.

4.9 WORK DONE IN A STEADY FLOW PROCESS

Specific enthalpy of working fluid of specific volume v will be h = u + pv. Differentiating this expression, we get

$$dh = du + d(pv) = du + pdv + vdp$$
(4.37)

According to first law of thermodynamics for closed system, we know that $\delta q = du + pdv$. Now equation (4.37) can be written as, $dh = \delta q + vdp$.

Substituting this in equation (4.32), we obtain

$$\delta q - \delta w = (\delta q + vdp) + d(ke) + d(pe)$$

$$\Rightarrow -\delta w = vdp + d(ke) + d(pe) \qquad (4.38)$$

If the change in kinetic and potential energies are negligible, *i.e.*, d(pe) = 0 and d(pe) = 0 then equation (4.38) may be written as

$$-\delta w = v dp$$
 or $\delta w = -v dp$

On integrating,
$$w_{1-2} = \int_{1}^{2} \delta w = -\int_{1}^{2} v dp$$
 (4.39)



Figures 4.7 (a) & (b) shows the difference between in a non-flow process and steady flow process respectively. In steady flow process, the work done $w_{1-2} = -\int_{1}^{2} v dp$ instead of $\int_{1}^{2} p dv$ in a

non-flow process. The (-ve) in (4.39) makes the integral positive quantity and represents the work done by the system.

4.10 WORK DONE VARIOUS STEADY FLOW PROCESS

The expression for work done in various steady flow processes are obtained as following:

(a) Constant volume process:

Work done
$$W_{1-2} = -\int_{1}^{2} \forall dp = \forall (p_1 - p_2)$$
 (4.40)

(b) Constant pressure process:

Work done
$$W_{1-2} = -\int_{1}^{2} \forall dp = \forall (p_1 - p_2) = 0$$
 (4.41)

(c) Constant temperature process:

For perfect gas, when temperature remains constant, we have

 $p \forall = p_1 \forall_1 = p_2 \forall_2 = Constant = RT$ Work done

$$W_{1-2} = -\int_{1}^{2} \forall dp = -\int_{1}^{2} \frac{p_{1} \forall_{1}}{p} dp = -p_{1} \forall_{1} \left[\ln p_{2} - \ln p_{1} \right]$$
$$= p_{1} \forall_{1} \ln \left(\frac{p_{1}}{p_{2}} \right) = RT \ln \left(\frac{p_{1}}{p_{2}} \right)$$
$$= p_{1} \forall_{1} \ln \left(\frac{\forall_{2}}{\forall_{1}} \right) = RT \ln \left(\frac{\forall_{2}}{\forall_{1}} \right)$$
(4.42)

(d) Adiabatic process

For adiabatic process

$$p \forall^{\gamma} = p_1 \forall_1^{\gamma} = p_2 \forall_2^{\gamma} = Constant$$
$$\forall = \forall_1 \left(\frac{p_1}{p}\right)^{1/\gamma}$$

 \Rightarrow

Work done

$$W_{1-2} = -\int_{1}^{2} \forall dp = -\int_{1}^{2} \forall_{1} \left(\frac{p_{1}}{p}\right)^{1/\gamma} dp$$

$$= -\forall_{1} p_{1}^{1/\gamma} \left[\frac{p_{2}^{1-1/\gamma}}{1-1/\gamma} - \frac{p_{1}^{1-1/\gamma}}{1-1/\gamma} \right]$$
$$= \frac{\gamma}{\gamma - 1} (p_{1} \forall_{1} - p_{2} \forall_{2})$$
(4.43)

(e) Polytropic process:

For polytropic process, $p \forall^n = p_1 \forall_1^n = p_2 \forall_2^n = Constant$ Proceed identically as done in adiabatic process.

Work done
$$W_{1-2} = \frac{n}{n-1} (p_1 \forall_1 - p_2 \forall_2)$$
 (4.44)

4.11 THROTTLING PROCESS

The throttling process is an irreversible steady flow expansion process in which a perfect gas is expanded through an orifice of minute dimension such as a narrow throat or a slightly opened valve as shown in figure 4.8. Due to the fall in pressure during expansion, the gas should come out with a large velocity, but due to high frictional resistance between the gas and the walls of





the aperture, there is no considerable change in velocity. The K.E of the gas is converted into heat which is utilised in warming the gas to its initial temperature. Here no heat is supplied or rejected during the throttling process and no work is done for free expansion, Therefore $q_{1-2} = 0$ and $w_{1-2} = 0$. Since there is no considerable change in velocity and inlet and outlet are at the same level, therefore, $V_1 = V_2$ and $Z_1 = Z_2$. Now steady flow energy equation (4.28) for unit mass will be reduced to $h_1 = h_2$ (4.45)

4.12 APPLICATION OF STEADY FLOW ENERGY EQUATION TO ENGINEERING SYSTEM

The application of steady flow energy equation to some of the engineering system such as nozzles, diffusers, boilers, turbine, compressor, condenser, and evaporator are discussed below.

4.12.1 Nozzle and Diffuser

Referring to figure 4.9, a nozzle is a device which increases the velocity as kinetic energy of the working fluid at the expense of its pressure drop. The nozzle is insulated so that no heat enters into or leaves from the system. The flow



through nozzle is considered as adiabatic process, so $q_{1-2} = 0$ and the system does not deliver any work, *i.e.*, $w_{1-2} = 0$. There is no change in potential energy, $pe_2 - pe_1 = 0$. Following equation (4.29) we know that the steady flow energy equation for unit mass flow energy equation is

$$q_{1-2} - w_{1-2} = (h_2 - h_1) + (ke_2 - ke_1) + (pe_2 - pe_1)$$

Thus for nozzle it will be reduced to

$$0 = (h_2 - h_1) + (ke_2 - ke_1)$$

$$\Rightarrow \qquad \frac{V_2^2}{2} - \frac{V_1^2}{2} = h_1 - h_2$$

This shows that the increase in K.E. will result in decrease in enthalpy. If the process is reversed, it is obvious that the decrease in K.E. will result in increase of enthalpy. Such a system is known as

Diffuser. From the above expression, we have $V_2 = \sqrt{V_1^2 + 2(h_1 - h_2)}$. If the initial velocity or velocity of approach V_1 is very small as compared to outlet velocity

 V_2 , then neglecting V_1 , the expression for V_2 will be $V_2 = \sqrt{2(h_1 - h_2)}$. Here, for continuous steady flow,

mass flow rate $m = \frac{A_1 V_1}{v_{s1}} = \frac{A_2 V_2}{v_{s2}}$

4.12.2 Boiler

It is a steam generating device which supplies heat to water and generates steam at certain pressure, as shown in figure 4.10. In this system, there is no change in kinetic energy and potential energy, also there is no work done by the system. So, $w_{1-2} = 0$, $pe_2 - pe_1 = 0$, $ke_2 - ke_1 = 0$. Hence the steady flow energy equation for a unit mass flow for a boiler is reduced to $q_{1-2} = h_2 - h_1$. This shows that the heat supplied to the system in a boiler increases the enthalpy of the system.

4.12.3 Turbine

A turbine as shown in figure 4.11 is a device which converts energy (K.E) of the working fluid into work in the turbine. The turbine is insulated so that there is no transfer of heat. In other words, the flow through a turbine is considered adiabatic. For turbine $q_{1-2} = 0$ and the change in kinetic energy and potential energy is negligible, *i.e.*, $pe_2 - pe_1 = 0$, $ke_2 - ke_1 = 0$. Then the steady flow energy equation for a unit mass flow for a turbine



Fig. 4.11

can be written as $w_{1-2} = h_2 - h_1$. This shows that the work done by the system is due to decrease in enthalpy of the working fluid.

4.12.4 Compressor

It is a device in which the fluid is compressed by the shaft work. It can be treated as turbine for the purpose of steady flow process. There are two types of compressors:

(a) Rotary Compressor: A rotary compressor, as shown in figure 4.12 is a device which compresses air and supplies the same at moderate pressure and in large quantities. In rotary compressor, the process can be treated as adiabatic as it is insulated so that no heat transfer takes place, *i.e.*, $q_{1-2} = 0$. In rotary compressor changes in kinetic and potential energies are negligible *i.e.* $pe_2 - pe_1 = 0$, $ke_2 - ke_1 = 0$. Then the steady flow energy equation for a unit mass flow for a rotary compressor will be reduced to $-(w_{1-2}) = h_2 - h_1$ or $w_{1-2} = h_2 - h_1$, or. Section 1 The (-ve) sign is used because work is Air (in)

done on the system. This equation shows that the work is done due to increase in enthalpy.

(b) Reciprocating Compressor: A reciprocating compressor, as shown in figure 4.13 is device supplying compressed air at a considerable high







Fig. 4.13

pressure and in small quantities. The reciprocating compressor is considered as a steady flow system provided it includes the receiver which reduces the fluctuation of flow considerably. Here the change in K.E. and P.E. are negligible, *i.e.*, $pe_2 - pe_1 = 0$, $ke_2 - ke_1 = 0$. So, for reciprocating compressor, the steady flow energy equation for a unit

mass flow will be reduced to $\left[-q_{1-2} - (-w_{1-2})\right] = h_2 - h_1$ or $w_{1-2} = q_{1-2} + (h_1 - h_2)$. The (-ve) sign is used because the heat is rejected and work is done on the system.

4.12.5 Condenser

A condenser, as shown in figure 4.14 is a device which is used to condense steam in case of steam power plant using water as the cooling medium, whereas in refrigeration system, it is used to condense refrigerant vapour using air as the cooling medium. In this system there is no change in K.E. and P.E., *i.e.*, $pe_2 - pe_1 = 0$, $ke_2 - ke_1 = 0$. Also there is no work done by the system, *i.e.*,



 $w_{1-2} = 0$. So the steady flow energy equation for unit mass will be reduced to $(-q_{1-2}) = h_2 - h_1$ or

 $q_{1-2} = h_1 - h_2$. The (-ve) sign is taken because the heat is lost by the coolant while passing through the condenser.

4.12.6 Evaporator

The evaporator, as shown in figure 4.15 is a device used in refrigeration systems in which the liquid refrigerant passes, receives heat and leaves as vapour refrigerant. For this system, the change in K.E. and P.E. is negligible. Also there is no work done by the system. So, $pe_2 - pe_1 = 0$, $ke_2 - ke_1 = 0$ and $w_{1-2} = 0$. Hence, the steady flow energy equation for unit mass in a evaporator will be $q_{1-2} = h_2 - h_1$. The process occurring in a evaporator is the reverse of that of a condenser.



Fig. 4.15

4.12.7 Heat Exchanger

The heat exchanger, as shown in figure 4.16 is a device in which heat is transferred from one fluid to another fluid, where steam is being condensed with the help of cooling water. Here the cooling water flows through the pipes where the steam flows over the tubes. The exchanger is well insulated from its out side. Therefore, heat does not cross the control volume surface; rather it is confined only between the two fluids. The change of K.E. and P.E. are also negligible and there is no central work done. So, for this system $pe_2 - pe_1 = 0$, $ke_2 - ke_1 = 0$ and $w_{1-2} = 0$. Hence, the steady flow energy equation for unit mass in a heat exchanger will be reduced to $(-q_{1-2}) = h_2 - h_1$ or $q_{1-2} = h_1 - h_2$. The (-ve) sign is taken because the heat is lost by the coolant while passing through the heat exchanger.



Fig. 4.16

Multiple Choice Questions

- 1. Joule's statement establishes that (when expressed in some units) during a cycle
 - (a) heat transfer is equal to work transfer
 - (b) work transfer is only a fraction of the heat transfer
 - (c) heat transfer is only a fraction of work transfer
 - (d) There is no relationship between work transfer and heat transfer
- 2. A system undergoes a process. The energy transfers are tabulated as follows

Н	eat transfer	Work transfer	Change in internal energy
	Q	W	U
(a)	52.7 kJ	20.7 kJ	32 kJ
(b)	32 kJ	52.7 kJ	20.7 kJ
(c)	30 kJ	20.4 kJ	52.7 kJ
(d)	20.7 kJ	32 kJ	52.7 kJ

3. A system undergoes a cycle of process AB, BC, CD and DA. The internal energy changes for various processes in kJ are

	AB	BC	CA	DA
(a)	32.5	-18.5	-30	16
(b)	18.5	-32.5	-30	16
(c)	30	-18.5	-32.5	16
(d)	16	-18.5	-32.5	30

(a) temperature and volume

(a)

- 4. Internal energy of ideal gas is a function of
- (b) pressure and volume
- (c) pressure and energy (d) temperature alone
- 5. Internal energy for gas in general can be written as

(a)
$$du = \left(\frac{du}{dt}\right)_{v} dt + \left(\frac{du}{dv}\right)_{T} dv$$

(b) $du = \left(\frac{du}{dt}\right)_{P} dt + \left(\frac{du}{dv}\right)_{T} dv$
(c) $du = \left(\frac{du}{dt}\right)_{v} dt + \left(\frac{du}{dv}\right)_{P} dv$
(d) $du = \left(\frac{du}{dt}\right)_{P} dt + \left(\frac{du}{dv}\right)_{P} dv$

6. Internal energy change of an ideal gas is expressed as

$$du = C_{\nu}dT \tag{b} \quad du = C_{p}dT$$

(c)
$$du = (C_P - C_v) dT$$
 (d) $du = C_v dT + \left(\frac{d_u}{d_v}\right)_T dv$

- 7. A closed system undergoes a process 1–2 for which the value of W_{1-2} and Q_{1-2} are 50 kJ and +20 kJ respectively. If the system is returned to state 1 and Q_{1-2} is –10 kJ, the work W_{1-2} is
 - (a) -80 kJ (b) +40 kJ
 - (c) -20 kJ (d) -40 kJ

b) pressure and volume

system is returned to state 1 in a record process for which W = 80 kJ. The heat transfer for the record process will be (a) 20 kJ (b) Zero (c) -20 kJ(d) intermediate 9. Figure shows three processes P, Q and R on $p - \forall$ plane namely isothermal compression $p \forall = C$, isentropic compression $p \forall^{\gamma} = C$, polytropic compression $p \forall^n = C$ for a perfect gas p Pressure where $1 < n < \gamma$. The correct identification is \cap (a) *P*(Isothermal), *Q*(Isentropic), *R*(Polytropic) (b) P(Polytrophic), Q(Isothermal), R(Isentropic) (c) P(Isentropic), Q(Polytropic), R(Isothermal) (d) *P*(Isothermal), *Q*(Polytropic), *R*(Isentropic) Volume 10. The work done in a free expansion process is (a) zero (b) minimum (c) maximum (d) positive 11. In a steady flow process (a) the mass flow rate is constant (b) the heat transfer rate is constant (c) the work transfer rate is constant (d) all of the above. 12. The work done in steady flow process is given by (b) $-\int_{1}^{2} pd \forall$ (a) $\int pd \forall$ (d) $-\int^2 \forall dp$ (c) $\int \forall dp$ 13. The throttling process is a (a) non-flow process (b) steady flow process (c) non-steady flow process (d) unsteady flow process Answers 6. (a) 2. (a) 5. (a) 7. (d) 1. (a) 3. (a) 4. (d) 8. (c) 9. (d) 10. (a) 11. (d) 12. (d) 13. (b)

8. A closed system goes from state 1 to state 2 in a process for which Q = 0 and W = 100 kJ. Then the

NUMERICAL EXAMPLES

EXAMPLE 1

A closed system receives a heat transfer of 120 kJ and delivers a work transfer of 150 kJ. Compute the change of internal energy.

SOLUTION

Here, heat transfer Q = 120 kJ, work transfer W = 150 kJ

Q = W + UWe know

 \Rightarrow

U = Q - W = 120 - 150 = -30 kJ

So, reduction of internal energy = 30 kJ.

EXAMPLE 2

300 tonne per second of steam is expanded in a turbine from an initial pressure of 90 bar. The specific enthalpies of steam at inlet and exit of the turbine are respectively 3300 kJ/kg and 2200 kJ/kg. Neglecting potential energy and kinetic energy terms and loss due to heat transfer, determine the output of the turbine in MW.

SOLUTION

Mass flow rate of steam (m) = 300 t/sec, initial pressure $p_1 = 90$ bar, pressure at exit = 0.1 bar,

specific enthalpy at inlet $h_1 = 3300 \text{ kJ/kg}$, specific enthalpy at exit $h_2 = 2200 \text{ kJ/kg}$

Work output of the turbine

$$W = m(h_1 - h_2) = 300 \times 10^3 (3300 - 2200) 10^{-3} = 33 \times 10^4 \text{MW}.$$

EXAMPLE 3

In an air compressor air flows steadily at the rate of 15 kg per minute. The air enters the compressor at 5 m/s with a pressure of 1 bar and a specific volume of 0.5 m3/kg. It leaves the compressor at 7.5 m/s with a pressure of 7 bar and a specific volume of $0.15 \text{ m}^3/\text{kg}$. The internal energy of the air leaving the compressor is 165 kJ/kg greater than that of the air entering. The cooling water in the compressor jackets absorbs heat from the air at the rate of 125 kJ/s. Determine: (i) power required to drive the compressor (ii) ratio of the inlet pipe diameter to outlet pipe diameter.

SOLUTION

Here,

 $m = 15 \text{ kg/min} = 0.25 \text{ kg/s}; V_1 = 5 \text{ m/s}; p_1 = 1 \text{ bar} = 0.1 \times 10^6 \text{ N/m}^2,$ $v_{s1} = 0.5 \text{ m}^3/\text{kg}; V_2 = 7.5 \text{ m/s}; p_2 = 7 \text{ bar} = 0.7 \times 10^6 \text{ N/m}^2;$ $v_{s2} = 0.15 \text{ m}^3/\text{kg}; u_2 - u_1 = 165 \text{ kJ/kg} = 165 \times 10^3 \text{ J/kg},$ $Q_{1-2} = 125 \text{ kJ/s}, q_{1-2} = 125/0.25 = 500 \times 10^3 \text{ J/kg}$ (i) Flow energy at inlet $p_1 v_{s1} = 0.1 \times 10^6 \times 0.5 = 50 \times 10^3 \text{ J/kg}$

 $p_2 v_{s2} = 0.7 \times 10^6 \times 0.15 = 105 \times 10^3 \text{ J/kg}$ Flow energy at outlet

Kinetic energy at inlet $ke_1 = \frac{V_1^2}{2} = \frac{5^2}{2} = 12.5 \text{ J/kg.}$

Kinetic energy at outlet $ke_2 = \frac{V_2^2}{2} = \frac{7.5^2}{2} = 28.1 \text{ J/kg}$

From the steady flow energy equation for a unit mass flow

$$u_1 + p_1 v_{s1} + k e_1 + p e_1 - q_{1-2} = u_2 + p_2 v_{s2} + k e_2 + p e_2 - w_{1-2}$$

Neglecting the potential energy, the steady flow energy equation for unit mass may be written as

$$u_{1} + p_{1}v_{s1} + ke_{1} - q_{1-2} = u_{2} + p_{2}v_{s2} + ke_{2} - w_{1-2}$$

$$w_{1-2} - q_{1-2} = (u_{2} - u_{1}) + (p_{2}v_{s2} - p_{1}v_{s1}) + (ke_{2} - ke_{1})$$

 \Rightarrow

 $w_{1-2} - 500 \times 10^3 = 165 - 10^3 + (105 \times 50)10^3 + (28.1 - 12.5)$ \Rightarrow

So,

 $w_{1-2} = 720 \times 10^3 \,\text{J/kg}$

Power required to drive the compressor

$$W_{1-2} = mW_{1-2} = 0.25 \times 720 \times 10^3 \text{ J/s} = 180 \times 10^3 \text{ J/s} = 180 \text{ kW}$$

(ii) Let, $D_1 =$ inlet pipe diameter, $D_2 =$ outlet pipe diameter

We know,

$$\frac{A_{1}V_{1}}{v_{s1}} = \frac{A_{2}V_{2}}{v_{s2}}$$
$$\frac{A_{1}}{A_{2}} = \frac{V_{2}}{V_{1}} \cdot \frac{v_{s1}}{v_{s2}}$$

$$D_1^2 - 7.5 \times 0$$

$$\Rightarrow \qquad \qquad \frac{D_1^2}{D_2^2} = \frac{7.5}{5} \times \frac{0.5}{0.15}$$
So,
$$\frac{D_1}{D_2} = 2.236$$

So,

 \Rightarrow

EXAMPLE 4

In a gas turbine, the gases flow at the rate of 5 kg/s. The gases enter the turbine at a pressure 7 bar with a velocity 120 m/s and leaves at a pressure 2 bar with velocity 250m/s. The turbine is insulated. If the enthalpy of the gas at inlet is 900 kJ/kg and at outlet 600 kJ/kg, determine the capacity of the turbine.

SOLUTION

Here,

$$m = 5 \text{kg/s}, p_1 = 0.7 \times 10^6 \text{ N/m}^2, V_1 = 120 \text{ m/s}$$

$$p_2 = 0.2 \times 10^6 \text{ N/m}^2, V_2 = 250 \text{ m/s},$$

$$h_1 = 900 \times 10^3 \text{ J/kg},$$

$$h_2 = 600 \times 10^3 \text{ J/kg}$$

From the steady flow energy equation for a unit mass flow is

$$h_1 + ke_1 + pe_1 - q_{1-2} = h_2 + ke_2 + pe_2 + w_{1-2}$$

Neglecting the potential energy,

$$w_{1-2} = (h_1 - h_2) + (ke_1 - ke_2) = (h_1 - h_2) + \left(\frac{V_1^2}{2} - \frac{V_2^2}{2}\right)$$

=
$$(900 - 600) 10^3 + \frac{1}{2} (120^2 - 250^2) = 275.95 \times 10^3 \text{ J/kg}$$

: Capacity of the turbine

$$= mw_{1-2} = 5 \times 275.95 \times 10^3 \text{ J/s} = 1379.75 \times 10^3 \text{ J/s} = 1379.75 \text{ kW}$$

EXAMPLE 5

A non-flow system has a mass of working fluid of 5 kg. The system undergoes a process in which 140 kJ heat is rejected to the surrounding and the system does 80kJ of work .Assuming that the initial specific internal energy of the system is 500 kJ/kg, determine the final specific internal energy.

SOLUTION

Here, $Q = -140 \text{ kJ}, W = 80 \text{ kJ}, u_1 = 500 \text{ kJ/kg}$ As, $Q - W = U_2 - U_1$, so, $U_2 - U_1 = -140 - 80 = -220 \text{ kJ}$ Difference in specific internal energy $u_2 - u_1 = \frac{U_2 - U_1}{m} = \frac{-220}{5} = -44 \text{ kJ/kg}$ So, $u_2 = -44 + 500 = 456 \text{ kJ/kg}$

EXAMPLE 6

A system is taken through a series of processes as a result of which it is restored to the initial state. The work and heat interaction for some of the process are measured and given below. Complete the table and determine the net work done and net heat interaction.

Process	W(kJ)	Q(kJ)	$\Delta U(kJ)$
1-2	100	200	?
2-3	?	-150	200
3-4	-250	?	?
4-1	300	?	50

SOLUTION

From 1st law of thermodynamics we know, $Q - W = U_2 - U_1$. For process 1 - 2,

$$\Delta U_{1-2} = Q_{1-2} - W_{1-2} = 200 - 100 = 100 \text{ kJ}$$

For process 2-3,

$$W_{2-3} = Q_{2-3} - \Delta U_{2-3} = -150 - 200 = -350 \text{ kJ}$$

For process 4 - 1,

$$Q_{4-1} = W_{4-1} + \Delta U_{4-1} = 300 + 50 = 350 \text{ kJ}$$

For the entire cycle, $\Delta U = 0$

So, $\Delta U_{1-2} + \Delta U_{2-3} + \Delta U_{3-4} + \Delta U_{4-1} = 0$

$$\Rightarrow \qquad 100 + 200 + \Delta U_{3-4} + 50 = 0$$

 $\Rightarrow \Delta U_{3-4} = -350$

For process 3 - 4,

$$Q_{3-4} = W_{3-4} + \Delta U_{3-4} = -250 - 350 = -600 \text{ kJ}$$

Net work done

$$W = W_{1-2} + W_{2-3} + W_{3-4} + W_{4-1}$$

$$= 100 - 350 - 250 + 300 = -200 \text{ kJ}$$

From 1^{st} law, it can be said net heat transfer = -200 kJ.

EXAMPLE 7

A stationary system consisting of 2 kg of fluid whose properties are related as u = (196 + 0.718t), $pv_s = [0.287(t + 273)]$, where u is the specific internal energy (kJ/kg), t is in °C, p is pressure (kPa) and v_s is specific volume (m³/kg). This system expands in an adiabatic process following $p \forall^{1.2} = constant$. The initial conditions are $p_1 = 1$ MPa, $t_1 = 200$ °C, final pressure $p_2 = 0.1$ MPa. Determine work done and change in energy for the process.

SOLUTION

Here,

$$U_1 = mu_1 = 2 (196 + 0.718 \times 200) = 679.2 \text{ kJ}$$

$$\forall_1 = mv_{s1} = 2 \frac{0.287(200 + 273)}{1 \times 10^3} = 0.2715 \text{ m}^3$$

Now,

$$\forall_2 = \forall_1 \left(\frac{p_1}{p_2}\right)^{\frac{1}{1.2}} = 0.2715 \left(\frac{1 \times 10^3}{0.1 \times 10^3}\right)^{\frac{1}{1.2}} = 1.849 \text{ m}^3$$

So,

$$v_{s2} = \left(\frac{\forall_2}{m}\right) = \frac{1.849}{2} = 0.9245 \text{ m}^3/\text{kg}$$

Hence,

$$t_2 = \frac{p_2 v_{s2}}{0.287} - 273 = \frac{0.1 \times 10^3 \times 0.9245}{0.287} - 273 = 49.12^{\circ}\text{C}$$

So,

$$U_2 = mu_2 = 2 (196 + 0.718 \times 49.12) = 462.54 \text{ kJ}$$

Here, $dE = dU = U_2 - U_1 = 462.54 - 679.2 = -216.66 \text{ kJ}$

From 1st law of thermodynamics

 $\delta W = \delta Q - dE = 0 - (-216.66) = 216.66 \text{ kJ}$

EXAMPLE 8

A gas undergoes a thermodynamic cycle consisting of the following processes: (i) *Process* 1–2: constant pressure at p = 1.4 bar, $\forall_1 = 0.28\text{m}^3$, $W_{1-2} = 10.5\text{kJ}$. (ii) *Process* 2–3: compression with at $p\forall$ = constant, and $U_3 = U_2$. (iii) *Process* 3–1: constant volume, $U_1 - U_3 = -26.4\text{kJ}$. There are no significant changes in KE and PE. (a) sketch the cycle on a $p - \forall$ diagram (b) calculate the net work done for the cycle in kJ (c) calculate the heat transfer.

 $\forall_2 = 0.355 \text{ m}^3$

SOLUTION

For process 1–2:

$$p_1 = p_2 = 1.4$$
 bar = 140 kPa, $\forall_1 = 0.28$ m³, $W_{1-2} = 10.5$ kJ

So,

 \Rightarrow

 \Rightarrow

 $W_{1-2} = p_1 \left(\forall_2 - \forall_1 \right)$ $10.5 = 140(\forall_2 - 0.28)$

For process 2–3: $\forall_3 = 0.28 \text{ m}^3$

So,
$$W_{2-3} = p_2 \forall_2 \ln\left(\frac{\forall_3}{\forall_2}\right) = 140 \times 0.355 \ln\left(\frac{0.28}{0.355}\right) = -11.79 \text{ kJ}$$

From 1st law of thermodynamics,

$$Q_{2-3} = W_{2-3} + (U_3 - U_2) = W_{2-3}$$

For process 3–1: being a constant volume process, $W_{3-1} = 0$

$$W_{\text{net}} = W_{1-2} + W_{2-3} + W_{3-1} = 10.5 - 11.79 = -1.29 \text{ kJ}$$

 $Q_{1-2} = W_{1-2} + (U_2 - U_1) = 10.5 + 26.4 = 36.9 \text{ kJ}$

EXAMPLE 9

When a system is taken form state *a* to state *b* in the figure shown along path a - c - b, 84 kJ of heat flows into the system and the system does 32 kJ of work. (a) how much will the heat flow into the system along path a-d-b, if the work done is 10.5 kJ? (b) when the system is returned from b to a along the curved path, the work done on the system is 21 kJ. Does the system absorb or liberate heat and how much of the heat is absorbed or liberated? Again if $U_a = 0$ and $U_d = 42$ kJ, find the heat absorbed in the process a - d and d - b.

SOLUTION

Here,

$$Q_{a-c-b} = 84 \text{ kJ}, W_{a-c-b} = 32 \text{ kJ}, W_{a-d-b} = 10.5 \text{ kJ}$$

We have,
$$Q_{a-c-b} = (U_b - U_a) + W_{a-c-b}$$

 \Rightarrow

$$(U_h - U_a) = 84 - 32 = 52 \text{ kJ}$$

So,

$$Q_{a-d-b} = (U_b - U_a) + W_{a-d-b} = 52 + 10.5 = 62.5 \text{ kJ}$$
$$Q_{b-a} = (U_a - U_b) + W_{b-a} = -52 - 21 = -73 \text{ kJ}$$



3

Pressure



125

Hence, the system liberates 73 kJ of heat.

Now, $W_{a-d-b} = W_{a-d} + W_{d-b} = W_{a-d} = 10.5 \text{ kJ}$ And $Q_{a-d} = (U_d - U_a) + W_{a-d} = 42 - 0 + 10.5 = 52.5 \text{ kJ}$ So, $Q_{d-b} = Q_{a-d-b} - Q_{a-d} = 62.5 - 52.5 = 10 \text{ kJ}$

EXAMPLE 10

A fluid system undergoes a non-flow frictionless process following the pressure-volume relation as $\left(p = \frac{5}{\forall} + 1.5\right)$ where p is in bar and \forall in m³. During the process the volume changes from 0.15 m³ to 0.05 m³ and the system rejects 45 kJ of heat. Determine (i) change in internal energy (ii) change in enthalpy. **SOLUTION**

Here,

$$p_1 = \frac{5}{\forall_1} + 1.5 = \frac{5}{0.15} + 1.5 = 34.83$$
 bar and
 $p_2 = \frac{5}{1.5} + 1.5 = \frac{5}{1.5} + 1.5 = 101.5$ bar

And

Now,

$$P_{2} = \frac{\forall}{\forall}_{2} = 1.3 = \frac{1}{0.05} + 1.3 = 101.5 \text{ bar}$$
$$W = \int_{\forall}_{\forall}_{1} p d\forall = \int_{0.15}^{0.05} \left(\frac{5}{\forall} + 1.5\right) \times 10^{5} d\forall$$
$$= \left[5\ln\frac{0.05}{0.15} + 1.5(0.05 - 0.15)\right] \times 10^{5}$$

 $= -5.64 \times 10^5 J = -564 \text{ kJ}$ From 1st law of thermodynamics,

$$\Delta U = Q - W = -45 - (-564) = 519 \text{ kJ}$$

So, internal energy is increased. Change in enthalpy,

$$\Delta H = \Delta U + (p_2 \forall_2 - p_1 \forall_1)$$

= 519 × 10³ + (101.5 × 0.05 - 34.83 × 0.15)10⁵
= 504 × 10³ J = 504 kJ

EXAMPLE 11

The properties of a system during a reversible constant pressure non-flow process p = 1.6 bar change from $v_{s1} = 0.3 \text{ m}^3/\text{kg}$, $t_1 = 20^{\circ}\text{C}$ to $v_{s2} = 0.55 \text{ m}^3/\text{kg}$, $t_2 = 260^{\circ}\text{C}$. The specific heat of the fluid is given by $C_p = \left(1.5 + \frac{75}{t+45}\right)\text{kJ/kg}^{\circ}\text{C}$. Determine (i) heat added/kg (ii) work done/kg (iii) change in internal energy/kg (iv) change in enthalpy/kg.

SOLUTION

(i) Heat added/kg of the fluid

$$q = \int_{t_1}^{t_2} C_p dt = \int_{20}^{260} \left(1.5 + \frac{75}{t+45} \right) dt$$
$$= 1.5 \left[260 - 20 \right] + 75 \ln \frac{260 + 45}{20 + 45}$$
$$= 475.94 \text{ kJ/kg}$$

(ii) Work done/kg of the fluid

$$w = \int_{v_1}^{v_2} p \, d\forall = p (v_{s2} - v_{s1})$$

= 1.6 ×10⁵ (0.55 - 0.3) Nm/kg
= 40 ×10³ J/kg = 40 kJ/kg

(iii) Change in internal energy/kg of the fluid

$$\Delta u = q - w = 475.94 - 40 = 435.94 \text{ kJ/kg}$$

$$h = u + w = q = 475.94 \text{ kJ/kg}$$

EXAMPLE 12

Air at 300°C and 1000 kPa (10 bar) expands to 300 kPa (3 bar) respectively following the law $pv^{1.35} = C$. Determine the work done per kg of air, and heat transfer if $C_p = 1$ kJ/kgK and $C_v = 0.714$ kJ/kgK. Solution

Here,

$$T_1 = 273 + 300 = 573$$
 K, $R = C_p - C_v = 1 - 0.714 = 0.286$ kJ/kg,

Hence,

$$\gamma = \frac{C_p}{C_v} = \frac{1}{0.714} = 1.4$$
$$T = T \left(\frac{p_2}{n} \right)^{\frac{n-1}{n}} = 572 \left(\frac{300}{1.35} \right)^{\frac{1.35-1}{1.35}} = 4$$

and

$$T_2 = T_1 \left(\frac{p_2}{p_1}\right)^{-n} = 573 \left(\frac{300}{1000}\right)^{\frac{1.05}{1.35}} = 419.37 \text{ K}$$

Work done in reversible polytropic process,

$$w_{1-2} = \frac{p_1 v_1 - p_2 v_2}{n-1} = \frac{R(T_1 - T_2)}{n-1}$$
$$= \frac{0.286(573 - 419.37)}{1.35 - 1} = 125.54 \text{ kJ/kg}$$



Heat transferred

$$q = \frac{\gamma - n}{\gamma - 1} \int_{1}^{2} p dv = \frac{1.4 - 1.25}{1.4 - 1} \times 125.54 = 47.08 \text{ kJ/kg}$$
EXAMPLE 13

At the inlet to a certain nozzle the specific enthalpy of fluid passing is 2800 kJ/kg and the velocity is 50 m/sec. At the discharge end the specific enthalpy is 2600 kJ/kg. The nozzle is horizontal and there is negligible heat loss from it, (i) find the velocity at exit of the nozzle (ii) if the inlet area is 900 cm^2 and specific volume at inlet is $0.187 \text{ cm}^3/\text{kg}$ find the mass flow rate (iii) if the specific volume at the nozzle exit is $0.498 \text{ m}^3/\text{kg}$, find the exit area of the nozzle.

SOLUTION

Here, $h_1 = 2800 \text{ kJ/kg}$, $h_2 = 2600 \text{ kJ/kg}$, $v_{s1} = 0.187 \text{ m}^3/\text{kg}$, $v_{s2} = 0.498 \text{ m}^3/\text{kg}$, $A_1 = 900 \text{ cm}^2$, $V_1 = 50 \text{ m/s}$

Applying energy equation at section 1 and 2,

$$h_1 + \frac{V_1^2}{2} + gZ_1 + \frac{\delta Q}{dm} = h_2 + \frac{V_2^2}{2} + gZ_2 + \frac{\delta W}{dm}$$

Depending on the given condition, above equation gets reduced to

$$h_1 + \frac{V_1^2}{2} = h_2 + \frac{V_2^2}{2}$$

$$\Rightarrow \qquad V_2 = \sqrt{V_1^2 + 2(h_1 - h_2)}$$

$$= \sqrt{50^2 + 2(2800 - 2600)} \times 10^3 = 634.4 \text{ m/s}$$

Mass flow rate $m = \frac{A_1 V_1}{v_{s1}} = \frac{900 \times 10^{-4} \times 50}{0.187} = 24.06 \text{ kg/s}$

Again,

$$m = \frac{A_1 V_1}{V_{s1}} = \frac{A_2 V_2}{V_{s2}}$$

_

 \Rightarrow

$$A_2 = m \cdot \frac{v_{s2}}{V_2} = 24.06 \times \frac{0.498}{634.4}$$

$$= 0.018887 \text{ m}^2 = 188.87 \text{ cm}^2$$

EXAMPLE 14

A turbine operates under steady flow condition receiving steam at the following state:

pressure 1.2 MPa, temperature 188°C, enthalpy 2785 kJ/kg, velocity 33.3 m/sec and elevation 3 m. The steam leaves the turbine at the following state:

pressure 20 kPa, enthalpy 2512 kJ/kg, velocity 100 m/sec and elevation 0 m. If heat is lost to the surroundings at the rate of 0.29 kJ/sec and the rate of steam flow through the turbine is 0.42 kg/sec, what is the power output of the turbine in kW?

SOLUTION

Here,

$$p_1 = 1.2 \text{ MPa}, p_2 = 20 \text{ kPa}, h_1 = 2785 \text{ kJ/kg}, h_2 = 2512 \text{ kJ/kg}$$

 $V_1 = 33.3 \text{ m/s}, V_2 = 100 \text{ m/s}, Z_1 = 3 \text{ m}, Z_2 = 0 \text{ m}$



: S.E.F.E. for control volume,

$$h_1 + \frac{V_1^2}{2} + Z_1g + q_{1-2} = h_2 + \frac{V_2^2}{2} + Z_2g + w_{1-2}$$

 \Rightarrow

$$w_{1-2} = (h_1 - h_2) + \frac{1}{2} (V_1^2 - V_2^2) + (Z_1 - Z_2)g + q_{1-2}$$

= $(2785 - 2512) + \frac{1}{2} (33.3^2 - 100^2) 10^{-3} + 9.81 \times 3 \times 10^{-3} + \frac{0.29}{0.42} = 269.27 \text{ kJ/kg}$

:. Power output $= mw_{1-2} = 0.42 \times 269.27 = 113.09 \text{ kW}$

EXAMPLE 15

At the Intel to a certain nozzle, the specific enthalpy of a fluid passing is 3000kJ/kg and the velocity is 60m/sec. At the discharge end, the specific enthalpy is 2762 kJ/kg. The nozzle is horizontal and there is negligible heat loss from it. (a) find the velocity at exit from the nozzle (b) if the inlet area 0.1m² and specific volume at inlet is 0.187 m³/kg, find the mass flow rate (c) if the specific volume at the nozzle exit is 0.498m³/kg, find exit area of the nozzle.

SOLUTION

Here, $h_1 = 3000 \text{ kJ/kg}, h_2 = 2762 \text{ kJ/kg}, V_1 = 60 \text{ m/s}, \text{ negligible } q_{1-2} = \text{negligible}, w_{1-2} = 0, A_1 = 0.1 \text{ m}^2,$ $v_{s1} = 0.187 \text{ m}^3/\text{kg}, v_{s2} = 0.498 \text{ m}^3/\text{kg}$ From steady flow energy equation (S.F.E.E.),

$$h_{1} + \frac{V_{1}^{2}}{2} + Z_{1}g + q_{1-2} = h_{2} + \frac{V_{2}^{2}}{2} + Z_{2}g + w_{1-2}$$
$$\left(V_{2}^{2} - V_{1}^{2}\right) = 2\left(h_{1} - h_{2}\right)$$

 \Rightarrow

 \Rightarrow

Mass flow rate

$$m = \frac{A_1 V_1}{v_{s1}} = \frac{0.1 \times 60}{0.187} = 32.08 \text{ kg/s}$$
$$A_2 = \left(\frac{A_1 V_1}{v_{s1}}\right) \cdot \left(\frac{v_{s1}}{V_2}\right) = \frac{32.08 \times 0.498}{692.5} = 0.023 \text{ m}^2$$

 $V_2 = \sqrt{60^2 + 2(3000 - 2762)10^3} = 692.5 \text{ m/s}$

Exit area of the nozzle

EXAMPLE 16

During a certain flow process the internal energy per kg of the gas decreases by 160 kJ/kg and the flow work increases by 5000 kNm for 10 kg of the gas. Determine the change in total enthalpy, specific enthalpy change in flow work per kg.

SOLUTION

As,
$$h_1 = u_1 + p_1 v_{s1}$$
 and $h_2 = u_2 + p_2 v_{s1}$

So,

:. Total enthalpy increases = $m(h_2 - h_1) = 10 \times 340 = 3400 \text{ kJ}$

EXAMPLE 17

A boiler is a steady flow system. Assuming that changes in K.E and P.E are negligible, determine the heat transferred per kg of steam if the water enters the boiler with an enthalpy of 160 kJ/kg and the steam leaves the boiler with 3128 kJ/kg.

SOLUTION

Now,

$$h_1 + \frac{V_1^2}{2} + Z_1g + q = h_2 + \frac{V_2^2}{2} + Z_2g + w_s$$

Assuming the changes in K.E and P.E negligible,

$$(q - w_s) = (h_2 - h_1) = 3182 - 160 = 3022 \text{ kJ/kg}$$

For boiler there will be no shaft work, so $w_s = 0$ Hence, heat transfer 3022 kJ/kg.

EXERCISE

- 1. Explain the first law of thermodynamics for the closed systems undergoing a cyclic change.
- 2. Define internal energy and prove that it is a property of a steam.
- 3. State the first law of thermodynamics and prove that for a non-flow process it leads to the energy equation $Q = \Delta E + Q$.
- 4. What is the mechanical equivalent of heat? Write down its value when heat is expressed in kJ and work is expressed in Nm.
- 5. What do you mean by PMM-I?
- 6. Why only in constant pressure non-flow process, the enthalpy change is equal to heat transfer?
- 7. Explain clearly the difference between a non-flow and steady flow process.
- 8. For isothermal flow and non-flow steady processes, prove that, $\int_{1}^{2} p d \forall = -\int_{1}^{2} \forall dp$, also state the

assumptions made.

- 9. Write down the general energy equation for steady flow system and simplify when applied for the following systems: (i) centrifugal water pump (ii) reciprocating air compressor (iii) steam nozzle (iv) gas turbine.
- 10. 1.35 kg of air at 179°C expands adiabatically and reversibly to three times its original volume and during the process, there is fall in temperature to 15.5°C. The mechanical work done during the process is 150 kNm. Determine C_p and C_y .
- 11. The energy of a system increases by 50 kJ while the system is receiving 43 kJ of work. How much heat is transferred and in which direction?

- 12. 2 kg of a gas at 10 bar expands adiabatically and reversibly till the pressure falls to 5 bar. During the process 0.1224 MNm non-flow work is done by the system and the temperature falls from 378°C to 257°C. Calculate the value of the index of expansion and characteristic gas constant.
- 13. Determine the characteristic gas constant if specific volume at 27°C and a pressure of 1 bar is 12 m³.
- 14. During a non-flow process, work/°C increase in dW/dt = 800 Nm/°C and the internal energy is given by U = (100 + 5t) Nm/°C. Determine the heat transferred from the system if temperature changes from 50°C to 140°C.
- 15. A gas is contained in a closed rigid tank fitted with paddle wheel. The paddle wheel stirs the gas for 20 minutes with the power varying with time W = -10t where W in watts, t in minutes. Heat transfer from the gas to the surroundings takes place at a constant rate of 50W. Determine (i) rate of change of energy of the gas at t = 10 minutes. (ii) net change in the energy of the gas.
- 16. A quantity of gas in a closed container has a pressure of 10 bar and volume of 0.925 m³ and temperature of 27°C. The characteristic gas constant is 287 Nm/kgK. Determine the mass of the gas in the container. If the pressure of the gas is decreased to 5 bar, how will the temperature of the gas be affected?
- 17. A gas occupies 9 m³ at a pressure of 2 bar. If it is compressed isothermally to 1.5 m³ determine the final pressure. 5 kg of gas at 150°C and 1.5 bar is heated at constant pressure until the volume is doubled. Determine the heat added and the final temperature reached.

This page intentionally left blank



SECOND LAW OF THERMODYNAMICS

5.1 INTRODUCTION

The first law of thermodynamics gives a quantitative estimate of heat and work interaction between a system and surroundings. If the system undergoes a thermodynamic process or a cycle, the first law cannot indicate whether the process or cycle in a particular direction would occur at all or not. It is the second law of thermodynamics which provides the answer to the questions of limitations in the first law of thermodynamics. Before we state the second law, it is necessary to clearly understand the meaning of the terms like thermal reservoir, source, sink, heat engine, refrigerators and heat pumps.

- (a) Thermal Reservoir: A thermal reservoir is a large system to which a finite quantity of energy as heat can be added or from which a finite amount of energy as heat can be extracted without changing its temperature. The ambient atmosphere is an example of a thermal reservoir. The environment constitutes the longest heat reservoir operating without any change in temperature. It is used as a heat sink and sometimes as a heat source.
- (b) Source: A source is a thermal reservoir at higher temperature from which energy in the form of heat is received by a heat engine.
- (c) Sink: A sink is a thermal reservoir at lower temperature to which energy as heat is rejected by a heat engine.
- (d) Heat Engine: A heat engine is a thermodynamic system operating in a cycle to which net positive amount of heat is added, and from which net positive amount of work is obtained. The heat engine operated in a cycle is known a heat engine cycle. A schematic diagram of heat engine is shown in figure 5.1. Thus for the satisfactory operation of a heat engine there should be at least two reservoirs of heat, one at a higher temperature and the other at a lower temperature, as shown in figure 5.1. Heat energy Q_H is received from the high temperature reservoir (source) at temperature reservoir(sink) at temperature T_L . Then, the remaining heat energy *i.e.*, $(Q_H Q_L)$ is converted into mechanical work (*W*). The ratio of the maximum mechanical work obtained to the total heat supplied to the engine is known as maximum thermal efficiency (η_{max}) of the engine.

Mathematically

$$\eta_{\text{max}} = \frac{\text{maximum work obtained}}{\text{total heat supplied}}$$
$$= \frac{W}{Q_H} = \frac{Q_H - Q_L}{Q_H}$$
$$= 1 - \frac{Q_L}{Q_H} = 1 - \frac{T_L}{T_H}$$

For a reversible engine we have, $\frac{Q_L}{T_L} = \frac{Q_H}{T_H}$.

(e) Refrigerator and Heat Pump (Reversible Heat Engine): It is a device that removes heat energy from a low temperature body (or reservoir) and rejects it to a high temperature body. It is a thermodynamic system operating in a cycle from which a net negative amount of heat is transferred, and on which a net negative amount of work is done.



Fig. 5.1

Schematic diagram of refrigerator and heat pump are shown in figure 5.2. Though there is no difference between the cycle of operations of refrigerator and a heat pump and achieve the same overall objective, but the basic purpose of each is quite different.

The device is called a refrigerator when the objective is to produce cold. It removes heat from a cool body at low temperature below atmospheric temperature, and rejects it to the surrounding *i.e.*,

atmospheric temperature. Thus, a refrigerator R operates between the temperature of the surrounding, and a temperature below that of the surrounding as shown in figure 5.2(a).

On the other hand, a heat pump is a device which is operating in a cyclic process, when the objective is to maintain the temperature of a hot body at a temperature higher that the temperature of surrounding. Thus, a heat pump P operates between the temperature of the surrounding, and a temperature above that of the surrounding as shown in figure 5.2(b).

In case of refrigerator, the atmosphere acts as a hot body while in case of a heat pump, the atmosphere acts as a cold body.

The performance of the refrigerator and heat pump is measured in terms of coefficient of performance (*COP*) which is defined as the ratio of the maximum heat transferred *i.e.*, heat taken from the cold body Q_L (in case of refrigerator) to the amount of work required W_R to produce the desired effect. From figure 5.2(a) we have,

$$Q_L + W_R = Q_H$$
 or, $W_R = Q_H - Q_L$



So, mathematically coefficient of performance for a refrigerator

$$(COP)_R = \frac{Q_L}{W_R} = \frac{Q_L}{Q_H - Q_L} = \frac{T_L}{T_H - T_L}$$

For heat pump coefficient of performance is defined as the ratio of the maximum heat transfer to the hot body Q_H to the amount of work (W_p) required to produce the desired effect. From figure 5.2(b) we have,

 $W_P + Q_L = Q_H$ or, $W_P = Q_H - Q_L$

Mathematically, maximum co-efficient of performance for heat pump

$$(COP)_{P} = \frac{Q_{H}}{W_{P}}$$
$$= \frac{Q_{H}}{Q_{H} - Q_{L}} = \frac{Q_{H} - Q_{L} + Q_{L}}{Q_{H} - Q_{L}}$$
$$= 1 + \frac{Q_{L}}{Q_{H} - Q_{L}} = 1 + \frac{T_{L}}{T_{H} - T_{L}}$$
$$= 1 + (COP)_{P}$$



Fig. 5.2 (b)

So, COP of a heat pump is greater than COP of a refrigerator by unity.

5.2 THE SECOND LAW OF THERMODYNAMICS

There are several statements of the second law of thermodynamics. Out of these, we will discuss two classical statements, known as the *Kelvin-Planck* statement and the *Clausius* statement. The Kelvin-Planck statement refers to a heat engine. The Clausius statement refers to a refrigerator or heat pump (reversible heat engine).

5.2.1 Kelvin-Planck Statement of the Second Law of Thermodynamics

According to Kelvin-Planck, *it is impossible to construct an engine working on a cyclic process, whose sole purpose is to convert heat energy from a single thermal reservoir into an equivalent amount of work.* In other words, no actual heat engine, working on a cyclic process, can convert whole of the heat supplied to it, into mechanical work. It means that there is a dissipation egradation of energy in the process of producing mechanical work from the heat supplied. Thus, the Kelvin-Planck statement of the second law of thermodynamics is sometimes known *a law of degradation of energy*.

(a) Perpetual Motion Machine of the Second Kind (briefly written a PMM-II or PMMSK): A Perpetual Motion Machine of the Second Kind (PMMSK) is a hypothetical

device (heat engine) which operates in cycle, absorbing energy as heat from a single thermal reservoir and delivering an equivalent amount of work. A schematic diagram of a perpetual motion machine of the second kind is shown in figure 5.3. The Kelvin-Planck statement of the second law of thermodynamics never supports possibility of devising a PMMSK. So, PMMSK violates the second law of thermodynamics. But the PMMSK does not violate the first law of thermodynamics, because, the first law of thermodynamics does not put any restriction as to the fraction of heat drawn from a heat reservoir that can be converted to work. 100% efficient heat engine does not violate the first of thermodynamics.

(b) Actual Heat Engine: The statement implies that a heat engine E of the type shown in figure 5.3 which will receive heat (Q_H) from a high temperature (T_H) reservoir, and convert it completely into work, is impossible. The only alternative is that there must be at least one other low temperature (T_L) reservoir to which heat Q_L must be rejected by the engine, because the second law of thermodynamics expresses the fact that the heat engine must reject heat to the low temperature reservoir and this also implies that two heat reservoirs are an essential minimum for a heat engine. Accordingly, we obtain a schematic representation of a heat engine as shown in figure 5.1. All heat engines must conform to this representation.

5.2.2 Clausius Statement of the Second Law of Thermodynamics

According to Clausius, it is impossible to construct a device (or self acting machine) working in a cyclic process, to transfer heat from a body at a lower temperature (T_L) to a body at a higher temperature (T_H) without the aid of an external agency. In other

words, heat cannot flow itself from a cold body to a hot body without the help of an external agency *i.e.*, without the expenditure of Mechanical work.

- (a) **PMM-II:** The device (reversible heat engine *i.e.*, refrigerator or heat pump) which violates the Clausius statement because no input work is supplied to the device to transfer heat from a cold body to a hot body is shown in figure 5.4. Such a device is called perpetual motion machine of the second kind.
- (b) 0 Reversible Heat Engine (Refrigerator/Heat Pump): The statement implies that refrigerator R or heat pump P of the type shown in figure 5.4 which will receive heat Q_L from a low temperature reservoir and transfer it to a high temperature reservoir is impossible. The only alternative is that there must be some work input (W). Accordingly, we obtain a schematic representation of refrigerator and heat







Е

Heat

Engine

High Temperature RESERVOIR (T_{H})

 $Q_{\rm H}$

 $W = Q_{\mu}$

pump as shown in figure 5.2(a) and (b). It is worth pointing out here that the devices in figure 5.2 represents a reversed heat engine since it comprises of all reversible processes.

5.3 EQUIVALENCE OF KELVIN-PLANCK AND CLAUSIUS OF SECOND LAW OF THERMODYNAMICS

The Kelvin-Planck and Clausius statement of the second law of thermodynamics are, though appear to be different from each other, but these two statements are virtually equivalent in all respects. The equivalence of Kelvin-Planck and Clausius statements can be proved, if it can be shown that the violation of Kelvin-Planck statement implies the violation of

Clausius statement and vice versa.

5.3.1 A Violation of Kelvin-Planck Statement Leads to a Violation of Clausius Statement of the Second Law of Thermodynamics

Consider a system as shown in figure 5.5. In this system, a heat engine having 100% thermal efficiency is violating the Kelvin-Planck statement as it converts the heat energy (Q_L) from a single high temperature reservoir T_H , into a equivalent amount of work *i.e.*, $W = Q_H$. This work output of the heat engine can be used to drive a heat pump (reversible heat engine) which $(W + Q_L) = Q_H + Q_L$ to a high temperature reservoir at T_H . If the combination of heat engine and a reversible heat engine is considered as a single system, as show in figure 5.5, then the result is a device that operates in a cycle and has no effect on the surroundings other than the transfer of heat Q_L from a low temperature reservoir to a high temperature reservoir, thus violating the Clausius statement. Hence, a violation of Kelvin-Planck statement leads to a violation of Clausius statement.

5.3.2 A Violation of Clausius Statement Leads to a Violation of Kelvin-Planck Statement of the Second Law of Thermodynamics

Consider a system as shown in figure 5.6. In this system, a heat pump (PMM-II) is violating the Clausius statement as it transfers heat from a low temperature reservoir at T_L to a high temperature reservoir at T_H without any expenditure of work. Now let a heat

engine, operating between the same heat reservoirs, receives an amount of heat Q_H (as discharged by heat pump) from the high temperature reservoir at T_H , does work $W_E = Q_H - Q_L$ and rejects an amount of heat Q_L to the low temperature reservoir at T_I . If the combination of the heat pump and the



Fig. 5.5



Fig. 5.6

heat engine is considered as a single system, as shown in figure 5.6, then the result is a device that operates in a cycle whose sole effect is to remove heat at the rate of $(Q_H - Q_L)$ and convert it completely into a equivalent amount of work, thus violating the Kelvin-Planck statement. Hence, a violation of Clausius statement leads to a violation of Kelvin-Planck statement of the second law of thermodynamics.

5.4 REVERSIBLE CYCLE

A thermodynamically reversible cycle consists of reversible processes. The process will be reversible when it is performed in such fashion that the system remains at all time infinitesimally near a state of thermodynamic equilibrium and at the end of process, both the system and the surroundings may be restored to their initial state. Therefore, reversible processes are purely ideal. In reversible process there are no dissipative effect, all the work done by the system during the performance of a process in one direction can be returned to the system during the reverse process. For example, consider a process in which the system (gas) is expanded from state 1 to state 2 following the path 1-2 as shown in figure 5.7. If during the thermodynamic process 1-2 the work done on the system by compressing is W_{1-2} and heat absorbed is Q_{1-2} , then by extracting heat Q_{1-2} from the system, we can bring the



system and surrounding, back from state 2 to state 1 following the same path 2–1, then the process is said to be reversible process.

5.5 IRREVERSIBLE CYCLE

A thermodynamic cycle will be irreversible, if any one of the process, constituting the cycle is an irreversible process. When heat and work are not completely restored back by reversing the process, then the process is known as irreversible or natural or real process. In an irreversible process, there is a loss of heat due to friction, radiation or conduction. In actual practice most of the process are irreversible to some degree. In irreversible cycle, the initial condition are not restored at the end of the cycle.

5.6 REVERSIBILITY AND IRREVERSIBILITY OF THERMODYNAMIC PROCESS

(a) Reversibility of Thermodynamic Process

Following are the conditions for reversibility of a cycle.

- 1. The departure from equilibrium is infinitesimal i.e., $dT \rightarrow 0$, $dp \rightarrow 0$. Accordingly, all the processes, taking place in the cycle of operation must be infinitely slowly, i.e., in a quasi-equilibrium manner.
- 2. The pressure and temperature of the working substance must not differ, appreciably from those of the surroundings at any stage in the process.

138

- 3. The working parts of the engine must be friction free, because, there should not be any loss of heat due to friction.
- 4. There should be no loss of energy during the cycle of operation.

(b) Irreversibility of Thermodynamic Process

The main causes for the irreversibility are

- 1. Due to lack of equilibrium during the process.
- 2. Due to involvement of dissipative effects.
- 3. Due to mechanical and fluid friction.
- 4. Due to unrestricted expansion.
- 5. Due to heat transfer with a finite temperature difference.

It is performed at a finite rate with finite potential differences. If there is some friction involved in the process, it becomes irreversible, because friction converts the mechanical work into heat. This heat cannot supply back the same amount of mechanical work, which was consumed for its production.

5.7 CARNOT CYCLE

Carnot cycle is a reversible cycle. This cycle was devised by *Carnot*, in 1828, who was the first scientist to analyse the problem of the efficiency of a heat engine, disregarding its mechanical details. This cycle, operating between two heat reservoirs, consists of an alternate series of two reversible isothermal and two reversible adiabatic or isentropic processes. These cycle can be operated either (i) as a vapour cycle (ii) or as a gas cycle.

(i) Vapour Cycle: A cycle is defined as a vapour cycle when the processes involve change of phase of working substance from liquid to vapour, and from vapour to liquid as shown in figure 5.8 on *p*-∀ diagram.

The Carnot vapour power cycle can be realised by a combination of flow processes as shown in figure 5.9. The cycle of operations is performed in different units having a defined function. The substance, *i.e.*, fluid, flows from one unit to the other unit in a steady flow. In the *boiler* isothermal process occurs, at temperature



Fig. 5.8



Fig. 5.9

 $T_{\mu\nu}$ and water is transferred into steam. In the *condenser* steam is transferred into water at temperature T_{L} , the reversible adiabatic expansion operation takes place in the *turbine* or *prime mover* and reversible adiabatic compression takes place in the pump.

The Carnot efficiency is given by

$$\eta_{\text{Carnot}} = \frac{T_H - T_L}{T_H}$$
$$= 1 - \frac{T_L}{T_H}$$

(ii) Gas Cycle: A cycle, in which the working substance remains a gas throughout, is referred to as a gas cycle, which is shown in figure 5.10.

In the Carnot cycle, the working substance (air supposed to behave like a perfect gas) is subjected to a cyclic operation, consisting of *two isothermal* and *two reversible adiabatic* operations.





Here, air is enclosed in a cylinder in which frictionless piston *A* moves. The walls of the cylinder and piston are perfectly non-conductor of heat. The bottom *B* of the cylinder can be covered by an insulating cap *i.e.*, the engine is assumed to work between two sources of infinite heat capacity, one at a higher temperature and other at a lower temperature. Four stages of the Carnot's cycle are

considered. The engine cylinder contains *m* kg of air at its original condition represented by point 1 on p- \forall diagram. At this point p_1 , T_1 and \forall_1 be the pressure, temperature and volume of the air.

1. First Stage (Isothermal expansion)

The source (hot body H.B.) at a higher temperature is brought in contact with the bottom *B* of the cylinder. The air expands at constant temperature from \forall_1 to \forall_2 *i.e.*, temperature T_2 at point 2 is equal to T_1 . Isothermal expansion is represented by the curve 1–2 on *p*- \forall diagram. Heat supplied by the hot body is totally absorbed by the air and is doing external work.

.. Heat supplied = Work done

So,

$$Q_{1-2} = p_1 \forall_1 \ln\left(\frac{\forall_2}{\forall_1}\right)$$
$$= mRT_1 \cdot \ln\left(\frac{\forall_2}{\forall_1}\right)$$
$$= 2.3 mRT_1 \log\left(\frac{\forall_2}{\forall_1}\right)$$
$$= 2.3 mRT_1 \log_r$$

Let $\frac{\forall_2}{\forall_1} = r = \text{expansion ratio.}$

2. Second Stage (Reversible adiabatic expansion)

The hot body is removed from the bottom of the cylinder and the insulating cap *i.e.*, is brought in contact. The air is now allowed to expand reversibly and adiabatically. Curve is represented by (2–3) on p- \forall diagram from \forall_2 to \forall_3 . Temperature falls from T_2 to T_3 . No heat is absorbed or rejected by the air.

:. Decrease internal energy = Work done by adiabatic exp.

$$= \frac{p_2 \forall_2 - p_3 \forall_3}{\gamma - 1} = \frac{mR(T_2 - T_3)}{\gamma - 1} = \frac{mR(T_1 - T_4)}{\gamma - 1}$$

3. Third Stage (Isothermal compression)

Removing the insulating cap from the bottom and is brought the cold body (C.B) in its contact. The air is compressed at constant temperature T_3 from \forall_3 to \forall_4 . Isothermal compression is represented by the curve (3–4) on p- \forall diagram. Heat is rejected to the cold body and is equal to the work done on the air.

Heat rejected
$$(Q_{3-4}) = p_3 \forall_3 \ln\left(\frac{\forall_3}{\forall_4}\right) = 2.3 \, mRT_3 \log r$$

$$\begin{bmatrix} \because & \frac{T_2}{T_3} = \left(\frac{\forall_3}{\forall_2}\right)^{\gamma-1} & \text{and} & \frac{T_1}{T_4} = \left(\frac{\forall_4}{\forall_1}\right)^{\gamma-1} \\ \therefore & \frac{\forall_3}{\forall_2} = \left(\frac{T_2}{T_3}\right)^{\frac{1}{\gamma-1}} & \text{and} & \frac{\forall_4}{\forall_1} = \left(\frac{T_1}{T_4}\right)^{\frac{1}{\gamma-1}} \\ \because & T_1 = T_2 \text{ and } T_3 = T_4 & \therefore & \frac{\forall_3}{\forall_2} = \frac{\forall_4}{\forall_1} \\ \therefore & \frac{\forall_3}{\forall_4} = \frac{\forall_2}{\forall_1} = r \end{bmatrix}$$

(4) Fourth Stage (Reversible adiabatic compression)

Now the insulated cap is brought in contact with the bottom of the cylinder and the air is allowed to be compressed reversibly and adiabatically represented by the curve (4-1) on p- \forall diagram. Temperature increases from T_4 to T_1 . Since no heat is absorbed or rejected,

: Increase in internal energy = Work done

$$=\frac{p_1\forall_1-p_4\forall_4}{\gamma-1}=\frac{mR(T_1-T_4)}{\gamma-1}$$

,

So, decrease in internal energy during reversible adiabatic expansion (2-3) = increase in internal energy during reversible adiabatic compression (4–1)

Work done = Heat supplied – Heat rejected

$$W = 2.3mRT_1 \log r - 2.3mRT_3 \log r$$

$$= 2.3mR \log r (T_1 - T_3)$$

$$\eta = \frac{\text{work done}}{\text{heat supplied}}$$

$$=\frac{T_1 - T_3}{T_1} = \left(1 - \frac{T_3}{T_1}\right)$$

5.8 **REVERSE CARNOT CYCLE**

The Carnot cycle is a reversible cycle. If the processes of the Carnot power cycle are carried out in the reverse order, it will become a refrigerator or heat pump cycle which is discussed before. It is then referred to as reverse Carnot cycle.

5.9 **CARNOT THEOREM**

It states that, no heat engine working between two temperature can be more efficient than the reversible engine working between the same two temperature or Carnot engine (hypothetical), among all engines operating between two fixed temperature, is the most efficient.

Multiple Choice Questions

- 1. In a power plant, Turbine work = 10,000 kJ, pump work = 10 kJ, heat supplied by boiler = 30,000 kJ. Thermal efficiency of the plant will be
 - (a) 27% (b) 33.333%
 - (c) 33.33% (d) 33.30%
- 2. A refrigerator and heat pump operate between the same temperature limits. If COP of the refrigerator is 4, the COP of the pump would be
 - (a) 3 (b) 5
 - (c) 4 (d) cannot predict
- 3. Energy requirement to heat an enclosure fully insulated will be less
 - (a) if heat pump is used for heating
 - (c) if steam is used for heating
- 4. It is impossible to construct an engine which while operating in a cycle produces no other effect to extract heat from a single reservoir and do equivalent amount of work.
 - (a) this refers to Clausius statement
 - (c) this refers to Carnot theorem (d) this does not refer to any of the above
- 5. Carnot cycle operates between temperature of 1000 K and 500 K. The Carnot efficiency will be 50% of working substance is
 - (a) air
 - (c) ammonia
- 6. Carnot cycle comprises
 - (a) two isothermal processes and two isentropic processes
 - (b) two constant volume processes and two isentropic processes
 - (c) one constant volume, one constant pressure and two isentropic processes
 - (d) two constant pressure and two isentropic processes
- 7. A heat engine is supplied with 800 kJ/sec of heat at 600 K and rejection take place at 300 K. Which of the following results report a reversible cycle.
 - (a) 200 kJ/sec are rejected (b) 400 kJ/sec are rejected
 - (c) 100 kJ/sec are rejected (d) 500 kJ/sec are rejected
- 8. A heat engine is supplied with 800 kJ/sec of heat at 600 K and rejects 100 kJ/sec at 300 K. The data refers to
 - (a) reversible cycle (b) irreversible cycle
 - (c) impossible cycle (d) none of the above
- 9. Equation $TdS = dU + pd\forall$ can be applied to processes which are
 - (a) only reversible (b) only irreversible
 - (c) reversible or irreversible (d) none of the above

(b) if electric strip heater is used for heating

(b) this refers to Kelvin-Planck statement

(d) if hot water is used for heating

143

- (b) nitrogen
- (d) any substance

10. Change in entropy for isothermal process carried on a gas whose specific volume changes from V_1 to V_2 and pressure changes from p_1 to p_2 is given by

(a)
$$\Delta S = -R \ln \frac{\forall_2}{\forall_1}$$
 (b) $\Delta S = R \ln \frac{p_2}{p_1}$

(c)
$$\Delta S = -R \ln \frac{p_2}{p_1}$$
 (d) $\Delta S = -R \ln \frac{\nabla_2}{\nabla_1} + R \ln \frac{p_2}{p_1}$

11. Change in entropy for polytropic process is expressed as

(a)
$$\Delta S = C_{\nu} \left(\frac{\gamma - n}{\gamma - 1}\right) \ln \frac{T_2}{T_1}$$

(b) $\Delta S = C_{\nu} \left(\frac{\gamma - n}{\gamma - 1}\right) \ln \frac{T_2}{T_1}$
(c) $\Delta S = \frac{\gamma - n}{\gamma - 1} \frac{R}{J} \ln \frac{T_2}{T_1}$
(c) $\Delta S = C_{\nu} \left(\frac{\gamma - n}{\gamma - 1}\right) \ln \frac{\forall_2}{\forall_1}$

12. The polytropic index of expansion *n* in the equation $pv^n = c$ for constant volume process is

(c)
$$\alpha$$
 (d) 0 (zero)

- 13. The change of entropy of a closed system
 - (a) is same for every process between two specified states
 - (b) is not the same for every process between two specified states
 - (c) is same only for isothermal process between two specified states
 - (d) is the same only for reversible adiabatic process between two specified states.
- 14. The entropy of fixed amount of an ideal gas
 - (a) decreases in every isothermal compression
 - (b) increase in every isothermal compression
 - (c) remains same in every isothermal compression
 - (d) may increase or decrease in every isothermal compression.
- 15. The specific internal energy, enthalpy amd entropy of an ideal gas are
 - (a) each function of temperature alone
 - (b) each function of pressure alone
 - (c) each function of volume alone
 - (d) each of the above statements (a), (b) and (c) are false.
- 16. One of the *Tds* equation has the form

(a)	Tds = dh + vdp	(b)	Tds = dh - vdp
(c)	Tds = du - pdv	(d)	Tds = dh + pdv

- 17. The entropy of a fixed amount incompressible substance
 - (a) decreases in every process in which temperature increases
 - (b) remains same in every process in which temperature increases
 - (c) increases in every process in which temperature increases
 - (d) is not affected in the process by increase or decrease of temperature.

18.	18. A closed system undergoes a process in which the work done on the system in 5 kJ and heat to occurs only at temperature T, the change in entropy for $Q = 5$ kJ and internally reversible pro						t trans	fer Q s is:									
	(a)	positive	·····p		- b,			(b)	negat	ive			<i>j</i> ·			
	(c)	zero						(d)	indete	ermina	ite.					
19	Sar	ne as 18 Fe	or $O =$	= 0 an	d inte	ernally	v revei	rsible	pro	cess T	he ch	ange in	entro	onv is			
	(a)	positive	- L					(b)	negat	ive			FJ			
	(c)	zero						((d)	indete	ermina	ite					
20.	Sar	ne as 18. Fe	or $O =$	=−5 k	J and	interr	nallv r	eversi	ble	the ch	ange i	n entroi	ov is				
	(a)	positive	2				J	(b)	negat	ive	1	5				
	(c)	zero						(d)	indete	ermina	ite					
21.	Same as 18. For $Q = +5$ kJ and internal irreversibilities present the entropy change is																
	(a)	positive	~					(b)	negat	ive		15	U			
	(c)	zero						(d)	indete	ermina	ite					
22.	Sar	ne as 18. Fe	or Q =	= 0 an	d inte	ernal i	rever	sibilit	ies	presen	t the e	ntropy	chnag	ge is			
	(a)	positive	-					(b)	negat	ive		-	-			
	(c)	zero						((d)	indete	ermina	ite					
23.	Sar	ne as 18. Fe	or Q =	=-5 k	J and	interr	nal irre	eversi	bili	ties pre	esent t	he entro	opy cl	hange i	S		
	(a)	positive						(b)	negat	ive						
	(c)	zero						(d)	indete	ermina	ite					
24.	Wh	nich of the f	follow	ving s	tatem	nent is	true										
	(a) no process is allowed in which the entropies of both the system and surroundings increase							se									
	(b) during a process, the entropy the system might decrease, while the entropy of surroundings																
	increases and conversely.																
	(c)	no proces unchange	s is a d.	llowe	ed in	which	the e	ntrop	ies	of botl	n the s	system	and t	he surr	oundi	ngs re	main
	(d)	a process	can c	occur	in wł	nich tł	ne entr	ropies	of	both th	ne syst	tem and	the	surrou	nding	decrea	ase.
A																	
Allowers																	
1. (c	1) 2 1) 12	. (b) 3.	(a) (a)	4. 14	(b) (a)	5. 15	(d)	6. 16	(a)) 7.) 17	(b)	8. 18	(c) (a)	9. 10	(c) (c)	10. 20	(c) (b)
21. (a	a) 22	. (c) 13. . (a) 23.	(a) (d)	24.	(a) (b)	13.	(u)	10.	(U	,	(0)	10.	(a)	19.		20.	(0)

NUMERICAL EXAMPLES

EXAMPLE 1

An engine works between the temperature limits of 1775 K and 375 K. What can be the maximum thermal efficiency of the engine?

SOLUTION

Here,

$$T_H = 1775 \text{ K}, T_L = 375 \text{ K}$$

 $\eta_{\text{max}} = \frac{T_H - T_L}{T_H} = \frac{1775 - 375}{1775} = 0.1887$

EXAMPLE 2

A heat engine receives heat @ 1500 kJ/min and gives an output of 8.2 kW. Determine thermal efficiency and the rate of heat rejection.

SOLUTION

Here,



Thermal efficiency
$$\eta_{\text{th}} = \frac{W}{Q_H} = \frac{8.2}{25} = 0.328$$

Rate of heat rejection

$$Q_L = Q_H - W = (25 - 8.2) \times 10^3 = 16.8 \times 10^3 \text{ J/s} = 16.8 \text{ kJ/s}$$

EXAMPLE 3

Find the coefficient of performance of heat transfer rate in the condenser of a refrigerator in kJ/hr, which has a refrigeration capacity of 12,000 kJ/hr, when power input is 0.75 kW.

SOLUTION

Here,

$$Q_L = 12000 \text{ kJ/hr}$$

$$W_R = 0.75 \times 60 \times 60 \text{ kJ/hr}$$

 $(\text{COP})_R = \frac{Q_L}{W_R} = \frac{12000}{0.75 \times 60 \times 60} = 4.44$

Hence,



Condenser

146

Heat transfer rate

 $Q_H = Q_L + W = 12000 + 0.75 \times 60 \times 60 = 14700 \text{ kJ/hr}$

EXAMPLE 4

A house requires 2×10^5 kJ/hr for heating in winter. Heat pump is used to absorb heat from a cold air house. Work required to operate heat pump is 3×10^4 kJ/hr. Determine heat abstracted from outside and coefficient of performance.

SOLUTION

Heat requirement Work required operating heat pump $W = 3 \times 10^5 \text{ kJ/hr}$ So, heat abstracted from outside $Q_L = Q_H - W = 2 \times 10^5 - 3 \times 10^4 = 17 \times 10^4 \text{ kJ/hr}$

$$(COP)_{H} = \frac{Q_{H}}{Q_{H} - Q_{L}} = \frac{2 \times 10^{5}}{2 \times 10^{5} - 17 \times 10^{4}} = 6.66$$

EXAMPLE 5

A cyclic heat engine operates between a source temperature of 1000°C and a sink temperature of 40°C. Find the least rate of heat rejection per kW net output of the engine.

SOLUTION

Here $T_H = 1000 + 273 = 1273$ K, $T_L = 40 + 273 = 313$ K For a reversible heat engine, the rate of heat rejection will be minimum for

$$\eta_{\text{max}} = 1 - \frac{T_L}{T_H} = 1 - \frac{313}{1273} = 0.754$$
$$Q_H = \frac{W_{\text{net}}}{\eta_{\text{max}}} = \frac{1}{0.754} = 1.326 \text{ kW}$$



Again,

EXAMPLE 6

A reversible heat engine operates between two reservoirs at 600°C and 40°C. The engine drives a reversible refrigerator, which operates, between reservoirs at 40°C and -20°C. The heat transfer to the engine 2 MJ and the net work output of the combined engine and refrigerator plant is 360 kJ. Find the heat transfer to the refrigerant and the net heat transfer at the reservoir at 40°C. Also find these values if the efficiency of the heat engine and COP of the refrigerator are each 40% of their maximum possible values.

SOLUTION

(i) Here $T_H = 600 + 273 = 873$ K, $T_{H1} = -20 + 273 = 253$ K, $T_L = T_{L1} = 40 + 273 = 313$ K.

Let $Q_{H} = 2000$ kJ is the heat transfer to the engine and Q_{H1} is the heat transfer to the refrigerant. So, maximum

efficiency of the heat engine $\eta_{\text{max}} = 1 - \frac{T_L}{T_H} = 1 - \frac{313}{873} = 0.6415$



Again, work done
$$W_E = Q_H \times \eta_{max} = 2000 \times 0.6415 = 1283 \text{ kJ}.$$

Since the net work output of the combined heat engine and refrigerator plant is $W = W_E - W_R = 360 \text{ kJ}$, work required for the refrigerator $W_R = W_E - W = 1283 - 360 = 923 \text{ kJ}$. We know maximum COP of the refrigerator

$$(\text{COP})_{\text{max}} = \frac{T_{H1}}{T_{H1} - T_{L1}} = \frac{253}{313 - 253} = 4.217$$

Again,

$$(\text{COP})_{\text{max}} = \frac{Q_{H1}}{Q_{H1} - Q_{L1}} = \frac{Q_{H1}}{W_R}$$

 \Rightarrow

$$Q_{H1} = 4.217 \times 923 = 3892.3 \text{ kJ}$$

 $Q_{L1} = Q_{H1} + W_R = 3892.3 + 923 = 4815.3 \text{ kJ}$

Here, And

$$Q_L = Q_H - W_E = 2000 - 1283 = 717 \text{ kJ}$$

Hence, net heat transfer to the reservoir at 40°C is $Q_L + Q_{L1} = 717 + 4815.3 = 5532.3$ kJ

(ii) When the efficiency of the heat engine and COP of the refrigerator are each 40% of their maximum possible values,

$$\eta_{actual} = 0.4 \times 0.6415 = 0.2566$$

Work done by the engine

$$W_E = Q_H \times \eta_{\text{actual}} = 2000 \times 0.2566 = 513.2 \text{ kJ}$$

Work required for the refrigerator

$$W_R = W_E - W = 513.2 - 360 = 153.2 \text{ kJ}$$

Now
$$(COP)_{actual} = 0.4 \times 4.217 = 1.6868$$

So, heat transfer to the refrigerant

Here,

$$Q_{H1} = 1.6868 \times 153.2 = 258.4 \text{ kJ}$$

 $Q_{L1} = Q_{H1} + W_R = 258.4 + 153.2 = 411.6 \text{ kJ} \text{ and}$
 $Q_L = Q_H - W_E = 2000 - 411.6 = 1588.4 \text{ kJ}$

Hence, net heat transfer to the reservoir at 40°C is

 $Q_L + Q_{L1} = 1588.4 + 411.6 = 2000 \text{ kJ}$

EXAMPLE 7

A Carnot engine working between 650 K and 310 K produce 150 kJ of work. Find the thermal efficiency and heat added during the process.

SOLUTION

Efficiency

$$\eta = 1 - \frac{T_L}{T_H} = 1 - \frac{310}{650} = 0.523$$

Heat added

 $Q = \frac{W}{\eta} = \frac{150}{0.523} = 286.8 \text{ kJ}$

EXAMPLE 8

A Carnot engine is operated between two reservoirs at T_{H} and T_{H1} . The work output of the engine is 0.6 times the heat rejected. The difference in temperature between the source and the sink is 200°C. Calculate the thermal efficiency, source temperature and sink temperature.

SOLUTION

Here,	$W = 0.6Q_{H1-L1}$ and $T_H - T_{H1} = 200^{\circ}$ C
Thermal efficiency	$\eta = \frac{W}{W + Q_{H1-L1}} = \frac{0.6Q_{H1-L1}}{0.6Q_{H1-L1} + Q_{H1-L1}} = \frac{0.6}{1.6} = 0.375$
Again,	$\eta = \frac{T_H - T_{H1}}{T_H} = \frac{200}{T_H}$
⇒	$T_H = \frac{200}{0.375} = 533.3 \text{ K} = 260.3^{\circ}\text{C}$
\Rightarrow	$T_{H1} = 260.3 - 200 = 60.3^{\circ}\mathrm{C}$
Example 9	

A Carnot cycle has an efficiency of 32%. Assuming that the lower temperature is kept constant, determine the % increase on the upper temperature of the cycle if the cycle efficiency is raised to 48%. **SOLUTION**

Here, $\eta = 1 - \frac{T_L}{T_H} = 0.32$

$$\Rightarrow$$

$$\frac{T_L}{T_H} = 1 - 0.32 = 0.68$$

When the efficiency is raised to 48%, $\frac{T_L}{T} = 1 - 0.48 = 0.52$

So,
$$\frac{T}{T_H} = \frac{0.68}{0.52} = 1.3077$$

or,

$$\frac{T_H + x}{T_H} = \frac{1 + 0.3077}{1}$$

 \Rightarrow

The % increase on the upper temperature of the cycle is 30.77

x = 0.3077

EXAMPLE 10

A vessel of 2.5 m³ capacity contains 1 kg mole of N_2 at 100°C. If the gas is cooled to 30°C, calculate the change in specific entropy.

 $T_H = 100 + 273 = 373 \text{ K}, T_L = 30 + 273 = 303 \text{ K}$

SOLUTION

The ratio of specific heats is 1.4 and 1 kg-mole of N_2 is 28 kg.

Here Now,

 \Rightarrow

$$C_p - C_v = R = \frac{R_u}{M} = \frac{8.314}{28} = 0.297$$

$$\Rightarrow \qquad 1.4C_v - C_v = 0.297$$

$$C_v = 0.74 \text{ kJ/kgK}$$

So,
$$S_H - S_L = mC_v \ln \frac{T_L}{T_H} = 1.0 \times 0.74 \ln \left(\frac{303}{373}\right) = -0.1536 \text{ kJ/kgK}$$

The negative sign indicates decrease in entropy.

EXAMPLE 11

300 kJ/s heat is supplied from a constant fixed source of 290°C to a heat engine. The heat rejection takes place at 8.5°C. The following results were obtained: (i) 215 kJ/s heat rejected (ii) 150 kJ/s heat rejected (iii) 75 kJ/s heat rejected. Determine these cycles are reversible or irreversible.

SOLUTION

Applying Clausius inequality to all these processes

(i)
$$\sum \frac{\delta Q}{T} = \frac{300}{290 + 273} - \frac{215}{8.5 + 273} = -0.2309 < 0$$

Hence, the cycle is irreversible.

(ii)
$$\sum \frac{\delta Q}{T} = \frac{300}{290 + 273} - \frac{150}{8.5 + 273} = 0$$

Hence, the cycle is reversible.

(iii)
$$\sum \frac{\delta Q}{T} = \frac{300}{290 + 273} - \frac{75}{8.5 + 273} = 0.2664 > 0$$

Hence, the cycle is impossible.

EXAMPLE 12

1 kg of air is heated on constant volume from 100°C to 400°C. If $C_v = 0.7186$ kJ/kgK, determine the change in entropy in the air.

SOLUTION

Change in entropy of air

$$\Delta S = \int_{T_L}^{T_H} \frac{dQ}{T} = \int_{100+273}^{400+273} mC_v \frac{dT}{T} = 1 \times 0.7186 \ln\left(\frac{673}{373}\right) = 0.4241 \text{ kJ/K}$$

EXAMPLE 13

A Carnot power cycle operates between 1200 K and 300 K. If 4.8 MJ of heat is supplied to the cycle, determine the heat rejected and the change in entropy for this rejection process.

SOLUTION

For Carnot cycle,

Hence,

Change in entropy for the process of heat rejection

EXAMPLE 14

A heat engine receives reversible 420 kJ/cycle heat from a source at 327°C and rejects heat reversible to a sink at 27°C. There are no other heat transfers. For each of the three hypothetical amount heat

rejection as (i) 210 kJ/cycle (ii) 105 kJ/cycle (iii) 315 kJ/cycle compute the cyclic integral $\frac{\delta Q}{T}$ in each case. Show which one is reversible, irreversible or impossible.

SOLUTION

Applying Clausius inequality to all these processes

(i)
$$\sum \frac{\delta Q}{T} = \frac{420}{327 + 273} - \frac{210}{27 + 273} = 0$$

Hence, the cycle is reversible.

(ii)
$$\sum \frac{\delta Q}{T} = \frac{420}{327 + 273} - \frac{105}{27 + 273} = 0.35 > 0$$

Hence, the cycle is impossible.

(iii)
$$\sum \frac{\delta Q}{T} = \frac{420}{327 + 273} - \frac{315}{27 + 273} = -0.35 < 0$$

Hence, the cycle is irreversible.



EXAMPLE 15

Two bodies each of equal mass *m* and equal heat capacity *C* are at temperature T_1 and T_2 respectively and $T_1 > T_2$. If the first body is used a source of heat for reversible engine and the second as the sink, show that the maximum work obtainable from such arrangement is $mC(\sqrt{T_1} - \sqrt{T_2})^2$.

SOLUTION

Let the common temperature of the two bodies be T_c . Maximum work is obtained when the engine operates along a reversible cycle. In that case, entropy change of the composite system formed by the two bodies is zero. Thus,

 $mC \ln \frac{T_c}{T_1} + mC \ln \frac{T_c}{T_2} = 0$ $mC \ln \frac{T_c^2}{T_1T_2} = 0$

 \Rightarrow

 \Rightarrow

 $\frac{T_c^2}{T_1 T_2} = 1$ $T_c = \sqrt{T_1 T_2}$

 \Rightarrow

Now the work done

$$W_{\text{net}} = mC(T_1 - T_c) - mC(T_c - T_2)$$

= $mC[(T_1 - \sqrt{T_1T_2}) - (\sqrt{T_1T_2 - T_2})]$
= $mC[T_1 + T_2 - 2\sqrt{T_1T_2}] = mC(\sqrt{T_1} - \sqrt{T_2})^2$

EXAMPLE 16

A Carnot heat engine draws heat from a reservoir at temperature T_A and rejects heat to another

reservoir at temperature T_{B} . The Carnot forward cycle again drives a Carnot reversed cycle engine or Carnot refrigerator, which absorbs heat from reservoir at temperature T_{C} and rejects heat to reservoir at temperature T_{B} . Derive an expression for the ratio of heat absorbed from reservoir at temperature T_{C} to heat drawn from reservoir at temperature T_{A} . If T_{A} is 500 K and T_{C} is 250 K, determine the temperature T_{B} , such that heat supplied to engine Q_{A} is equal to heat absorbed by refrigerator Q_{C} . Determine efficiency and COP of Carnot refrigerator.

SOLUTION

$$\eta_{\text{engine}} = \frac{Q_A - Q_B}{Q_A} = \frac{T_A - T_B}{T_A} = \frac{W_{\text{carnot}}}{Q_A}$$



$$W_{\text{carnot}} = \frac{T_A - T_B}{T_A} \cdot Q_A \tag{a}$$

Now,
$$(COP)_{\text{refrigerator}} = \frac{Q_C}{Q_{2B} - Q_C} = \frac{T_C}{T_B - T_C} = \frac{Q_C}{W_{\text{carnot}}}$$

$$W_{\text{carnot}} = \frac{T_B - T_C}{T_C} Q_C \tag{b}$$

Equating (a) and (b),
$$\frac{Q_C}{Q_A} = \frac{\frac{T_A - T_B}{T_A}}{\frac{T_B - T_C}{T_C}} = \frac{T_C}{T_A} \left[\frac{T_A - T_B}{T_B - T_C} \right]$$

$$\Rightarrow \qquad 1 = \frac{250}{500} \left(\frac{500 - T_B}{T_B - 250} \right)$$

$$\Rightarrow$$
 $T_B = 333.33 \text{ K}$

Hence,

 \Rightarrow

 \Rightarrow

$$\eta_{\text{engine}} = \frac{T_A - T_B}{T_A} = 0.334$$

And
$$(COP)_{\text{refrigerator}} = \frac{T_C}{T_B - T_C} = 3$$
.

EXAMPLE 17

A mass *m* of a fluid at temperature T_1 is mixed with an equal mass of the same fluid at temperature T_2 . Prove that the resultant change of entropy is $2mC \ln \frac{0.5(T_1 + T_2)}{\sqrt{T_1T_2}}$ and also prove that it is always positive.

SOLUTION

The thermal equilibrium is obtained at common temperature $T_c = \frac{T_1 + T_2}{2}$. Thus, change in entropy is given by

$$\Delta S = mC \int_{T_1}^{T_c} \frac{dT}{T} - mC \int_{T_c}^{T_2} \frac{dT}{T}$$
$$= mC \left[\ln \frac{T_c}{2T_1} - \ln \frac{2T_2}{T_c} \right]$$
$$= mC \left[\ln \frac{T_c}{2T_1} + \ln \frac{T_c}{2T_2} \right]$$

$$= mC \ln \frac{T_c^2}{4T_1T_2}$$
$$= mC \ln \left[\frac{T_1 + T_2}{2\sqrt{T_1T_2}}\right]^2 = 2mC \ln \left[\frac{T_1 + T_2}{2\sqrt{T_1T_2}}\right]$$

Now the arithmetic mean $\frac{T_1 + T_2}{2}$ is always greater than the geometric mean $\sqrt{T_1T_2}$. Hence the above result is positive.

EXAMPLE 18

A hot iron forging of specific heat 0.12 kcal/kgK weighing 30 kg and at a temperature of 500°C is dropped in 200 kg of oil at 20°C and having a specific heat of 0.6 kcal/kgK for quenching. The iron forging has surroundings. Determine entropy change of forging, entropy change of oil and entropy change of universe. **SOLUTION**

SOLUTIO

Here,
$$T_{\text{forging}} = 500 + 273 = 773 \text{ K} \text{ and } T_{\text{oil}} = 20 + 273 = 293 \text{ K}$$

By energy balance, a common equilibrium temperature (T_c) is obtained after quenching.

So,
$$m_{\text{forging}}C_{\text{forging}}\left(T_{\text{forging}}-T_{c}\right) = m_{\text{oil}}C_{\text{oil}}\left(T_{c}-T_{\text{oil}}\right)$$

$$\Rightarrow \qquad 30 \times 0.12 \times (773 - T_c) = 200 \times 0.6 \times (T_c - 293)$$

$$\Rightarrow$$
 $T_c = 307 \text{ K}$

Thus,

$$\Delta S_{\text{forging}} = 30 \times 0.12 \ln \left(\frac{307}{773}\right) = -3.32 \text{ kcal/K}$$

and

$$\Delta S_{\text{oil}} = 200 \times 0.6 \ln\left(\frac{307}{293}\right) = 5.60 \text{ kcal/K}$$

Hence,

$$\Delta S_{\text{universe}} = \Delta S_{\text{forging}} + \Delta S_{\text{oil}}$$

= -3.32 + 5.60 = 2.28 kcal/K

Thus, the process is irreversible.

EXERCISE

- 1. State the limitations of first law of thermodynamics.
- 2. What is the difference between a heat engine and a reversed heat engine?
- 3. Enumerate the conditions that must be fulfilled by a reversible process. Give some examples of ideal reversible processes.
- 4. What is an irreversible process? Give some examples of irreversible processes.

- 5. Narrate the propositions of Clausius statement and Kelvin-Planck statement.
- 6. Define heat engine, refrigerator and heat pump.
- 7. What is the perpetual motion machine of the second kind?
- 8. What is the efficiency of heat engine?
- 9. What is COP of the reversible heat engine?
- 10. Derive the expression for COP of heat pump and refrigerator. Establish the relation between COP of pump and COP of refrigerator.
- A heat engine is supplied with 278 kJ/s of heat at a constant fixed temperature of 283°C and the heat rejection takes places at 5°C. Then results reported as: (i) 208 kJ/s rejected (ii) 139 kJ/s rejected (iii) 70 kJ/s rejected. Classify which of the results report a reversible cycle or irreversible cycle or impossible cycle.
- 12. Source *A* can supply energy @11 MJ/min at 320°C. A second source *B* can supply energy @110 MJ/min at 68°C. Which source would you choose to supply energy to an ideal reversible engine that is to produce large amount of power if the temperature of the surroundings is 40°C.
- A domestic food freezer maintains a temperature of -15°C. The ambient air temperature is 30°C. If heat leaks into the freezer continuously @1.75 kJ/s, what will be the least power necessary to pump this heat out continuously.
- 14. A heat engine supplied heat @1700 kJ/min and gives an output of 9 kW. Determine the thermal efficiency and rate of heat rejection.
- 15. A fish freezing plant requires 50 tons of refrigerant. The freezing temperature is -40°C while the ambient temperature is 35°C. If the performance of the plant is 15% of the theoretical reversed Carnot cycle working within the same temperature limits, calculate the power required.

This page intentionally left blank



ENTROPY

6.1 INTRODUCTION

When first law of thermodynamics was applied for thermodynamic processes, the existence of a property, internal energy, was found. This is called the *corollary* 1 of the First Law of Thermodynamics. Similarly, when second law of thermodynamics applied to thermodynamic processes, the second law also leads to the definition of a new property, known as entropy. So, if the first law is said to be the law of internal energy, then second law may be stated to be the law of entropy. The term 'entropy' was first introduced by Clausius. Entropy means transformation. It is an important thermodynamic property of a working substance, which increases with addition of heat and decreases with its removal. As a matter of fact, it is comparatively more easy to define change of entropy than to define the term entropy. When change of entropy (dS) is multiplied by the absolute temperature (T), gives the heat (δQ) absorbed or rejected by the working substance.

Mathematically, heat absorbed by the working substance

$$\delta Q = TdS$$
 $\therefore dS = \frac{\delta Q}{T}$ (6.1)

So change of entropy is the change of heat per unit absolute temperature.

Entropy may be defined as following:

Entropy is function of a quantity of heat which shows the possibility of conversion of that heat into work. The increase in entropy is small when heat is added at high temperature and is greater when heat addition is at low temperature. Thus for minimum entropy there is maximum availability for conversion into work, for maximum entropy, there is a minimum availability for conversion into work.

It may also be defined as the thermodynamic property of a substance which remains constant when substance is expanded or expressed reversible adiabatically in a cylinder.

6.2 IMPORTANCE OF ENTROPY

In a reversible Carnot cycle, the maximum possible efficiency obtainable is given by

$$\eta_{\max} = \frac{T_H - T_L}{T_H} \tag{a}$$

where T_{H} is highest absolute temperature and T_{L} is lowest absolute temperature. Again, in general efficiency is given by

$$\eta_{\text{max}} = \frac{\text{maximum work obtained}}{\text{heat applied or absorbed}}$$

$$= \frac{\delta W}{\delta Q}$$

$$\delta W = \eta_{\text{max}} \times \delta Q$$

$$= \left(\frac{T_H - T_L}{T_H}\right) \times \delta Q \qquad (6.2)$$
The above expression may be written as

or,

For one degree temperature drop, the above expression may be written as

$$\delta W = \frac{\delta Q}{T} \quad \text{where } T_H - T_L = 1$$

= dS
= change in entropy (6.3)

It can be understood from above equation that

- 1. The maximum amount of work obtained per degree drop in temperature is equal to change of entropy i.e., change of entropy is equal to the maximum work for one degree temperature drop.
- 2. The change in entropy may be regarded as a measure of the rate of availability or unavailability of heat for temperature into work.
- 3. The increase in entropy is obtained from a given quantity of heat at low temperature.

6.3 UNITS OF ENTROPY

We know

$$dS = \frac{\delta Q}{T}$$

So unit of entropy depends upon the unit of heat and the absolute temperature. If δQ is expressed in kJ and T in K, then unit of entropy kJ/K.

6.4 CLAUSIUS THEOREM

Clausius theorem state that "the cyclic integral $\left(\frac{\delta Q}{T}\right)$ for a reversible cycle is equal to zero"

Mathematically

$$\oint_R \frac{\delta Q}{T} = 0$$

The letter R emphasizes the fact that the equation is valid only for reversible cycle. It can be proved as below.

Entropy

Consider a system undergoing a reversible process 1-b-2 from a equilibrium state 1 to another equilibrium state point 2, by the reversible path is shown in figure 6.1. Let a reversible adiabatic 1-a be drawn through 1 and another reversible adiabatic 2-c be drawn through 2, then reversible isothermal a-b-c is drawn in such a way that the area enclosed by 1-a-b and b-c-2 are equal on a $p-\forall$ diagram.

Then the actual reversible process 1-b-2 can be replaced by the combination of the above processes, resulting in the overall process 1-a-b-c-2. The process 1-b-2 and 1-a-b-c-2 are equivalent if and only if the heat and work interaction are identical for both the processes. Applying the first law of thermodynamics for the two processes,

1. For the process 1-b-2,

$$U_2 - U_1 = Q_{1-b-2} - W_{1-b-2}$$
(a)

2. For the process 1-a-b-c-2,

$$U_{2} - U_{1} = Q_{1-a-b-c-2} - W_{1-a-b-c-2}$$

= $Q_{1-a} + Q_{a-b-c} + Q_{c-2} - W_{1-a-b-c-2}$
= $Q_{a-b-c} - W_{1-a-b-c-2}$ (b)

Since 1-a and c-2 are reversible adiabatic process, $Q_{1-a} = 0$ and $Q_{c-2} = 0$

Now consider the cycle 1-a-b-c-2-b-1. The net work done in this cycle is zero since the area enclosed by 1-a-b is equal to the area enclosed by b-c-2 on a $p-\forall$ diagram in figure 6.1

i.e.,
$$\oint \delta W = W_{1-a-b-c-2} + W_{2-b-1} = 0$$

or, $W_{1-a-b-c-2} = -W_{2-b-1} = W_{1-b-2}$ (c)

Comparing the equation (a) and (c), we have

$$Q_{1-a-b-c-2} - W_{1-a-b-c-2} = Q_{a-b-c} - W_{1-b-2}$$
$$= Q_{1-b-2} - W_{1-b-2}$$
$$\therefore \qquad Q_{1-a-b-c-2} = Q_{a-b-c} = Q_{1-b-2}$$

That is, heat transferred along in the processes 1-b-2 is equal to the heat transferred in the processes 1-a-b-c-2. So, the process 1-b-2 can be replaced by the process 1-a-b-c-2.

Any reversible path may be substituted by a reversible zigzag path, between the same end states, consisting of a reversible adiabatic followed by a reversible isothermal and then by a reversible adiabatic, such that the heat transferred during the isothermal processes is same as that transferred during the original process.



A

Fig. 6.1



Consider a system which undergoes the cyclic change as shown in figure 6.2. Now draw a family of closely spaced adiabatic line similar to a-3 and 2-f covering the entire cycle. Then join the adjacent adiabatic lines d-e-f as shown in figure 6.2, such that the area of 1-a-b = area of b-c-2 and area of 3-d-e = area of e-f-4.

Thus the segments 1-b-2 and 3-e-4 of the reversible cycle are replaced by 1-a-b-c-2 and 3-d-e-f-4, respectively. Now the differential cycle a-b-c-4-f-e-d-a is a Carnot cycle. Thus the original cycle can be transformed into a combination of several differential Carnot cycles. If the adiabatics are close to one another and number of Carnot cycles are large, the saw-toothed zig-zag line will coincide with the original cycle.





For the elemental cycle a-d-f-c, dQ_1 heat is absorbed reversibly at T_1 and δQ_2 , heat is rejected reversibly at T_2

$$\therefore \qquad \frac{\delta Q_1}{T_1} = \frac{\delta Q_2}{T_2} \tag{6.5}$$

Heat supply is taken as positive and heat rejected as negative

00

00

So,
$$\frac{\partial Q_1}{T_1} + \frac{\partial Q_2}{T_2} = 0$$

If similar equations are written from all the elemental Carnot cycles, then for whole original cycle

$$\frac{\delta Q_1}{T_1} + \frac{\delta Q_2}{T_2} + \frac{\delta Q_3}{T_3} + \dots = 0$$
$$\sum \frac{\delta Q}{T} = 0$$

or for a differential Carnot cycles

$$\oint \frac{\delta Q}{T} = 0 \quad [\text{ for reversible cycle}]$$

This is known as *Clausius theorem*.

For reversible cycle the thermal efficiency is given by

$$\eta_{\rm th} = \frac{Q_1 - Q_2}{Q_1} = \frac{T_1 - T_2}{T_1} \tag{6.6}$$

But for an irreversible engine according to Carnot theorem, the thermal efficiency is

$$\frac{Q_1 - Q_2}{Q_1} < \frac{T_1 - T_2}{T_1}$$
$$1 - \frac{Q_2}{Q_1} < 1 - \frac{T_2}{T_1}$$

or

160

or	$\frac{Q_2}{Q_1} > \frac{T_2}{T_1}$
or	$\frac{Q_2}{T_2} > \frac{Q_1}{T_1}$
or	$\frac{Q_1}{T_1} < \frac{Q_2}{T_2}$
or	$\frac{Q_1}{T_1} - \frac{Q_2}{T_2} < 0$

Therefore, in differential form, for irreversible engine

$$\oint \frac{\delta Q}{T} < 0. \tag{6.7}$$

6.5 ENTROPY— A POINT FUNCTION OR A PROPERTY OF A SYSTEM

Let a system shown in Figure 6.3 be taken from an initial equilibrium state 1 to a final equilibrium state 2 by the reversible path A and the system is brought back from 2 to 1 by another reversible path B. Then the two path Aand B together constitute a reversible cycle. From Clasius theorem, we have

$$\oint_{A,B} \frac{\delta Q}{T} = 0$$

[It can be replaced as the sum of two integrals one for path A and other for path B]

> $\oint_{1A}^{2} \frac{\delta Q}{T} + \oint_{2B}^{1} \frac{\delta Q}{T} = 0$ $\oint_{1.4}^{2} \frac{\delta Q}{T} = -\oint_{2B}^{1} \frac{\delta Q}{T}$

or

:.

Path *B* is a reversible path,

$$\oint_{1A}^{2} \frac{\delta Q}{T} = \oint_{1B}^{2} \frac{\delta Q}{T}$$

Since path A and B represent any two reversible paths, $\oint \frac{\delta Q}{T}$ is independent of the path.

So $\oint \frac{\delta Q}{T}$ is a property, is a point function, and its values depend on initial and final points only, not on the path







$$\therefore \qquad \oint_{1}^{2} \frac{\delta Q}{T} = S_{2} - S_{1} = dS \text{ (change of entropy)}$$
(6.8)

6.6 CLAUSIUS INEQUALITY

The Clausius inequality states that when a closed system undergoes a cyclic process, the cyclic integral of $\frac{\delta Q}{T}$ is less than zero (negative) for irreversible cyclic process and equal to zero for a reversible cyclic process.

$$\oint \frac{\delta Q}{T} < 0 \text{ (for irreversible process)}$$
$$\oint \frac{\delta Q}{T} = 0 \text{ (for reversible process)}$$

Combining the above two equation, we have

$$\oint \frac{\delta Q}{T} \le 0 \tag{6.9}$$

which is called *Clausius inequality*. It is already proved in art.6.4. The Clausius inequality not only gives mechanical expression to the 2nd law of thermodynamics, but it also gives the quantitative measure of irreversibility of the system. For example,

For irreversible cyclic process $\oint \frac{\delta Q}{T} < 0$ So, it may be written as $\oint \frac{\delta Q}{T} + m = 0$

where m represents the amount by which the given cyclic process is irreversible.

When *m* is equal to zero, that is $\oint \frac{\delta Q}{T} > 0$ is impossible, because it voilates the 2nd law of thermodynamics.

6.7 PRINCIPLE OF INCREASE OF ENTROPY

The equation for the Clausius inequality is written as

$$\oint \frac{\delta Q}{T} \le 0 \tag{a}$$

Again we know the change in entropy

$$dS = \frac{\delta Q}{T}$$

Being a thermodynamic property, the cyclic-integral of a entropy is zero, that is $\oint dS = 0$

Therefore equation (a) may be written as

$$\oint \frac{\delta Q}{T} \le \oint dS$$

:.

:.

 $\therefore \qquad \qquad \frac{\delta Q}{T} \le dS$

$$dS \ge \frac{\delta Q}{T}$$

When the process is reversible, then

$$dS = \frac{\delta Q}{T} \tag{6.10}$$

and when the process is irreversible, then

$$dS > \frac{\delta Q}{T} \tag{6.11}$$

In case of isolated system $\delta Q = 0$, if we apply the equation (b) in isolated system the equation (b) will be

 $dS \ge 0$,

when dS = 0, for reversible cyclic process, S =Constant

Now for an irreversible cyclic process

dS > 0

Since, in actual practice, all processes are irreversible, therefore the entropy of such system like universe goes on increasing. This is known as the *principle of increase of entropy*.

6.8 ENTROPY AND TEMPERATURE RELATION

Entropy is the thermodynamic property of a working substance that increases with the increase of temperature and decreases with decrease of temperature.

if

 δQ = heat absorbed, dS = increase in entropy

T = absolute temperature

 $\therefore \qquad \delta Q = T \cdot dS$

Consider the heating of a working substance by a reversible process as shown in the figure 6.4 by a curve from 1 to 2 in temperature-entropy diagram.

At a small quantity of heat is supplied δQ ,

$$\delta Q = T \cdot dS$$

= area under curve during change of entropy

 $\int_{1}^{2} \delta Q = \int_{1}^{2} T \cdot dS$ $\int_{1}^{2} dS = \int_{1}^{2} \frac{\delta Q}{T}$



(6.12)



(b)
So, total area under the *T-S* diagram of any thermodynamic process represents the heat absorbed or heat rejected.

As $\int \frac{\delta Q}{T}$ is same for all reversible paths between states 1 and 2 it is independent of the path and

is a function of end states only.

6.9 GENERAL EXPRESSION FOR CHANGE OF ENTROPY OF A PERFECT GAS

Consider a certain quantity of perfect gas heated by any thermodynamic process.

Let, m = mass of the gas $p_1 \text{ and } p_2 = \text{initial and final pressure}$ $\forall_1 \text{ and } \forall_2 = \text{initial and final volume}$ $T_1 \text{ and } T_2 = \text{initial and final absolute temperature}$

(i) In terms of volume and absolute temperature

$$\delta Q = dU + \delta W$$

= $m \cdot C_V \cdot dT + p d \forall$
$$\frac{\delta Q}{T} = m \cdot C_V \cdot \frac{dT}{T} + p \frac{d \forall}{T}$$

 $p \forall = mRT$
$$\frac{p}{T} = \frac{mR}{\forall}$$

:.

∴ As,

 $dS = mC_{v} \cdot \frac{dT}{T} + \frac{mR}{\forall} \cdot d\forall$

 $\left[\because \frac{\delta Q}{T} = dS\right]$

Integrating within limits,

$$\int_{S_1}^{S_2} dS = mC_v \int_{T_1}^{T_2} \frac{dT}{T} + mR \cdot \int_{v_1}^{v_2} \frac{d\forall}{\forall}$$

$$(S_2 - S_1) = mC_v (\ln T_2 - \ln T_1) + mR (\ln \forall_2 - \ln \forall_1)$$

$$= 2.3 m \left[C_v \cdot \log\left(\frac{T_2}{T_1}\right) + R \cdot \log\left(\frac{\forall_2}{\forall_1}\right) \right]$$

$$= 2.3 m \left[C_v \cdot \log\left(\frac{T_2}{T_1}\right) + (C_p - C_v) \log\left(\frac{\forall_2}{\forall_1}\right) \right]$$

or,

(ii) In terms of pressure and absolute temperature

We know
$$\frac{p_1 \forall_1}{T_1} = \frac{p_2 \forall_2}{T_2}$$
$$\therefore \qquad \frac{\forall_2}{\forall_1} = \frac{p_1}{p_2} \cdot \frac{T_2}{T_1}$$

Proceeding in the similar way,

$$(S_2 - S_1) = mC_v \cdot \ln\left(\frac{T_2}{T_1}\right) + mR \cdot \ln\left(\frac{p_1}{p_2} \cdot \frac{T_2}{T_1}\right)$$
$$= mC_v \cdot \ln\left(\frac{T_2}{T_1}\right) + mR \cdot \ln\left(\frac{p_1}{p_2}\right) + mR \ln\left(\frac{T_2}{T_1}\right)$$
$$= m\ln\left(\frac{T_2}{T_1}\right) [C_v + R] + mR \ln\left(\frac{p_1}{p_2}\right)$$
$$= 2.3 m \left[C_p \cdot \log\left(\frac{T_2}{T_1}\right) + R \cdot \log\left(\frac{p_1}{p_2}\right)\right]$$
$$= 2.3 m \left[C_p \cdot \log\left(\frac{T_2}{T_1}\right) + R \cdot \log\left(\frac{p_1}{p_2}\right)\right]$$

(iii) In terms of pressure and volume

We know,
$$\frac{p_1 \forall_1}{T_1} = \frac{p_2 \forall_2}{T_2}$$

 $\therefore \qquad \frac{T_2}{T_1} = \frac{p_2}{p_1} \cdot \frac{\forall_2}{\forall_1}$

Proceeding in the similar fashion

$$(S_2 - S_1) = mC_v \cdot \ln\left(\frac{p_2}{p_1} \cdot \frac{\forall_2}{\forall_1}\right) + mR \cdot \ln\left(\frac{\forall_2}{\forall_1}\right)$$
$$= mC_v \cdot \ln\left(\frac{p_2}{p_1}\right) + mC_v \cdot \ln\left(\frac{\forall_2}{\forall_1}\right) + mR \ln\left(\frac{\forall_2}{\forall_1}\right)$$
$$= mC_v \cdot \ln\left(\frac{p_2}{p_1}\right) + m\ln\left(\frac{\forall_2}{\forall_1}\right) \cdot (C_v + R)$$
$$= mC_v \cdot \ln\left(\frac{p_2}{p_1}\right) + m\ln\left(\frac{\forall_2}{\forall_1}\right) \cdot C_p$$
$$= 2.3 m \left[C_v \cdot \log\left(\frac{p_2}{p_1}\right) + C_p \cdot \log\left(\frac{\forall_2}{\forall_1}\right)\right]$$

6.10 CHANGE OF ENTROPY OF PERFECT GAS DURING VARIOUS THERMODYNAMIC PROCESS

(a) Constant volume process (Isochoric process): A certain quantity of gas is heated at constant volume, where

 p_1 = Initial pressure, T_1 = Initial temperature



(b) Constant pressure processes (Isobaric process):

A certain quantity of a perfect gas heated at constant pressure, where

 $m = \text{mass of gas}, \forall_1 = \text{initial volume}, T_1 = \text{initial temperature}$

 \forall_2, T_2 = Corresponding values of volume and temperature for the final conditions for a small change of temperature

,	$\sum_{p} p$	A
	$\frac{\partial Q}{T} = mC_p \cdot \frac{dT}{T}$	
<i>.</i>	$dS = mC_p \cdot \frac{dT}{T}$	p = c
Now,	$\int_{S_1}^{S_2} dS = mC_p \cdot \int_{T_1}^{T_2} \frac{dT}{T}$	T ₁
.: .	$(S_2 - S_1) = mC_p \cdot [\ln T]_{T_1}^{T_2}$	
	T.	S ₁ Entropy S ₂
	$= mC_p \cdot \ln \frac{T_2}{T_1}$	Fig. 6.6

Entropy

$$= 2.3 \ mC_p \cdot \log\left(\frac{T_2}{T_1}\right)$$
$$= 2.3 \ mC_p \cdot \log\left(\frac{\forall_2}{\forall_1}\right) \qquad \qquad \left[\because \frac{T_2}{T_1} = \frac{\forall_2}{\forall_1}\right]$$

6.11 IRREVERSIBILITY

The actual work which a system does is always less than the idealized reversible work, and the difference between the two is called the *irreversibility of the process*.

If, I = the irreversibility,

W = the actual work

 $W_{\rm max}$ = the idealized reversible worker maximum work.

Then, Irreversibility, $I = W_{\text{max}} - W$

This is also sometimes referred to a *degradation* or *dissipation*.

The irreversibility in (i) non-flow process and (ii) steady flow-process is discussed below

(i) For a non-flow process

In non-flow process between the equilibrium states, when the system exchanges heat only with environment, irreversibility for unit mass,

$$i = w_{\max} - w$$

$$i = [(u_1 - u_2) - T_0 (S_1 - S_2)] - [(u - u_2) + q]$$

$$= T_0 (S_2 - S_1) - q$$

$$= T_0 (\Delta S)_{\text{system}} + T_0 (\Delta S)_{\text{surrounding}}$$

So,

 $i \ge 0$

(ii) For steady flow-process

$$i = w_{\text{max}} - w$$

$$= \left[\left\{ b_1 + \frac{C_1^2}{2} + gz_1 \right\} - \left\{ b_2 + \frac{C_2^2}{2} + gz_2 \right\} \right]$$

$$- \left[\left\{ h_1 + \frac{C_1^2}{2} + gz_1 \right\} - \left\{ h_2 + \frac{C_2^2}{2} + gz_2 \right\} + q \right]$$

$$= T_0 \left(S_2 - S_1 \right) - q$$

$$= T_0 \left(\Delta S \right)_{\text{system}} + T_0 \left(\Delta S \right)_{\text{surrounding}}$$

The same expression for irreversibility applies to both flow and non-flow process.

		Multi	ple C	Choice Questions		
1.	Sec	ond law of thermodynamics de	efines			
	(a)	heat	(b)	work	(c)	enthalpy
	(d)	entropy	(e)	internal energy		
2.	For	a reversible adiabatic process	, the c	hange in entropy is		
	(a)	zero	(b)	minimum	(c)	maximum
	(d)	infinite	(c)	unity		
3.	For	any reversible process, the ch	ange	in entropy of the system and	surro	unding is
	(a)	zero	(b)	unity	(c)	negative
	(d)	positive	(e)	infinite		
4.	For	any irreversible process the n	et enti	ropy change is		
	(a)	zero	(b)	positive	(c)	negative
	(d)	infinite	(e)	unity		
5.	The	processes of a Carnot cycle a	re			
	(a)	two adiabatics and two const	ant vo	blume		
	(D)	one constant volume and one	e cons	tant pressure and two isentr	opics	
	(d)	two constant volumes and two	inais	hermals		
	(e)	two isothermals and two isen	tropic	s		
6	Isot	ropic flow is	opie			
0.	(a)	irreversible adiabatic flow	(b)	ideal fluid flow	(c)	perfect gas flow
	(d)	frictionless reversible flow	(e)	reversible adiabatic flow		
7.	In a	Carnot engine, when the wor	king s	ubstance gives heat to the si	nk	
	(a)	the temperature of the sink in	crease	es		
	(b)	the temperature of the sink rem	mains	same		
	(c)	the temperature of the source	decre	eases		
	(d)	the temperature of both the s	ink an	d the source decrease		
	(e)	changes depend on the operation	ating c	conditions		
8.	If th	ne temperature of the source is	increa	ased, the efficiency of a Carr	ot eng	ine
	(a)	decreases	(b)	increases	(c)	does not change
	(d)	will be equal to the efficiency	(e)	depends on other factors		
0	T 1		· · · ·			
9.	1 ne	working substance	engine	e depends on		
	(a)	on the temperature of the sou	rce or	nlv		
	(c)	on the temperature of the sin	c only	,		
	(d)	on the temperature of both th	e soui	rce and the sink		
	(e)	on the construction of engine	e			
		e				

10. The efficiency of a carnot engine using on ideal gas as the working substance is

(a)
$$\frac{T_1 - T_2}{T_1}$$
 (b) $\frac{T_1}{T_1 - T_2}$ (c) $\frac{T_1 T_2}{T_1 - T_2}$

(d)
$$\frac{T_1 - T_2}{T_1 T_2}$$
 (e) $\frac{T_2 (T_1 - T_2)}{T_1 (T_1 + T_2)}$

- 11. In a reversible cycle, the entropy of the system
 - (a) increases (b) decreases (c) does not change
 - (d) first increases and then decreases
 - (e) depends on the properties of working substance
- 12. A frictionless heat engine can be 100% efficient only if its exhaust temperature is:
 - (a) equal to its input temperature
 - (c) 0°C
 - (e) -100°C
- 13. Kelvin-Planck's law deals with
 - (a) conservation of energy
 - (c) conservation of mass
- (b) conservation of heat
- (d) conservation of heat into work

(b) less than its input temperature

- (e) conservation of work into heat
- 14. Which of the following statements is correct according to Clausius statements of second law of thermodynamics?
 - (a) it is impossible to transfer heat from a body at a lower temperature to a body at a higher temperature

(d) 0K

- (b) it is impossible to transfer heat from a body at a lower temperature to a body at a higher temperature without the aid of an external source
- (c) it is possible to transfer heat from a body at a lower temperature to a body at a higher temperature by using refrigeration cycle
- (d) none of the above
- 15. According to Kelvin Planck's statement of second law of temperature
 - (a) it is impossible to construct an engine working on a cyclic process, whose sole purpose is to convert heat energy into work
 - (b) it is possible to construct an engine working on cyclic process, whose sole purpose is to convert the heat energy into work
 - (c) it is impossible to construct a device which while working in a cyclic process produces as effect other than the transfer of heat from a colder body to a hotter body
 - (d) when two dissimilar metals are heated at one end and cooled at the other, the e.m.f. developed is proportional to the difference of their temperatures at the two end
 - (e) none of the above
- 16. The property of working substance which increases or decreases as the heat is supplied or removed in a reversible manner is known as
 - (a) enthalpy (b) internal energy
 - (c) entropy (d) external energy

17.	The	e entropy may be expressed as	a fun	ction of			
	(a)	pressure and temperature	(b)	temperatu	re and volume	(c)	heat and work
	(d)	all of the above	(e)	none of the	ie above		
18.	The	e change of entropy, when hea	it is ab	osorbed by	the gas is		
	(a)	positive	(b)	negative		(c)	positive or negative
19.	Wh	ich of the following statement	s is co	orrect?			
	(a)	the increase in entropy is obt	ained	from a give	on quantity of he	at at a l	ow temperature
	(b)	the change of entropy may b transformation into work.	e rega	arded as a r	neasure of the ra	ate of t	he availability of heat for
	(c)	the entropy represents the ma	iximui	m amount o	f work obtainab	le per d	egree drop in temperature
	(d)	all of the above					
20.	The	e condition for the reversibility	of cy	cle is			
	(a)	the pressure and temperature surroundings at any stage in	of wor the p	rking substa rocess	ince must not dif	fer, app	preciably from those of the
	(b)	all the processes taking place	in the	e cycle of op	peration, must be	e extren	nely slow
	(c)	the working parts of the engin	ne mu	st be frictio	n-free		
	(d)	there should be no loss of en	ergy d	luring the c	ycle of operation	n	
	(e)	all of the above					
21.	In a	in irreversible process there is	a				
	(a)	loss of heat		(b)	no loss of work	2	
	(c)	gain of heat		(d)	no gain of heat		
22.	The	e main cause for the irreversibi	lity is				
	(a)	mechanical and fluid friction					
	(b)	unrestricted expansion					
	(c)	heat transfer with a finite tem	peratu	ire differen	ce		
	(d)	all of the above					
23.	The	e efficiency of the Carnot cycle	e may	be increase	d by		
	(a)	increasing the highest tempe	rature				
	(b)	decreasing the highest tempe	erature	e			
	(c)	increasing the lowest temper	ature				
	(d)	decreasing the lowest temper	ature				
	(e)	keeping the lowest temperatu	are co	nstant.			
24.	Wh	ich of the following is the corr	rect st	atement?			
	(a)	all the reversible engines hav	e the s	same efficie	ency		
	(b)	all the reversible and irrevers	ible ei	ngines have	the same efficie	ency	
	(c)	irreversible engines have max	imum	efficiency			
	(d)	all engine are designed as rev	ersible	e in order to	obtain maximur	n effici	ency

Entr	opy																	171
								A	nsw	ers								
1. 11. 21.	(d) 2.(c) 12.(a) 22.	(a) (d) (d)	3. 13. 23.	(a) (d) (d)	4. 14. 24.	(b) (b) (a)	5. 15.	(e) (c)	6. 16.	(e) (c)	7. 17.	(b) (a)	8. 18.	(b) (a)	9. 19.	(d) (d)	10. 20.	(a) (e)
						N	UME	RIC	AL E	ΞΧΑ	MPL	.ES						

EXAMPLE 1

 0.05 m^3 of air at a pressure of 8 bar and temperature 280°C expands to eight times its original volume and the final temperature after expansion is 25°C. Calculate the change of entropy of air during the process. Assume $C_p = 1.005 \text{ kJ/kg K}$ and $C_v = 0.712 \text{ kJ/kg K}$.

SOLUTION

Given

Let Here

$$\forall_{2} = 8 \forall_{1} = 8 \times 0.05 = 4m^{3}$$

$$T_{2} = 25^{\circ}C = (25 + 273) = 298 K$$

$$C_{p} = 1.005 \text{ kJ/kg K}, C_{v} = 0.712 \text{ kJ/kg K}$$

$$m = \text{mass of air in kg}$$

$$R = C_{p} - C_{v} = 1.005 - 0.712 = 0.293 \text{ kJ/kg K}$$

$$= 293 \text{ J/kg K}$$

 $\forall_1 = 0.05 \, m^3, p_1 = 8 \, \text{bar} = 8 \times 10^5 \, \text{N/m}^2$

 $T_1 = 280^{\circ}\text{C} = (280 + 273) = 553 \text{ K}$

As,

$$p_1 \forall_1 = mRT_1$$

$$m = \frac{p_1 \forall_1}{RT_1} = \frac{8 \times 10^5 \times 0.05}{293 \times 553} = 0.247 \text{ kg}$$

Change of entropy

$$(S_2 - S_1) = 2.3 \text{ m} \left[C_v \cdot \log\left(\frac{T_2}{T_1}\right) + R \cdot \log\left(\frac{\forall_2}{\forall_1}\right) \right]$$

= 2.3 × 0.247 $\left[0.712 \log\left(\frac{298}{553}\right) + 0.293 \log\left(\frac{0.4}{0.05}\right) \right]$
= 0.568 $\left[-0.19 + 0.26 \right]$
= 0.04 kJ/K

EXAMPLE 2

A vessel of 2.5 m³ capacity contains 1 kg mole of N_2 at 100°C. If the gas is cooled to 30°C, calculate the change in specific entropy. The ratio of specific heats is 1.4 and one kg-mole nitrogen is 28 kg.

SOLUTION

Here,

 $\forall = 2.5 \text{ m}^3$, M = 1 kg mole = 28 kg

$$T_{1} = 100^{\circ}\text{C} = (100 + 273) = 373 \text{ K}$$

$$T_{2} = 30^{\circ}\text{C} = (30 + 273) = 303 \text{ K}$$

$$V = \frac{C_{p}}{C_{v}} = 1.4$$

$$R_{u} = \text{Universal gas constant} = 8.314 \text{ kJ/kg K}$$

$$R = \frac{R_{u}}{M} = \frac{8.314}{28} = 0.297 \text{ kJ/kg K}$$

$$\therefore \quad C_{p} - C_{v} = R \qquad \therefore 1.4C_{v} - C_{v} = 0.297 \text{ So}, C_{v} = \frac{0.297}{0.4} = 0.74 \text{ kJ/kg K}$$
Hence, $(S_{2} - S_{1}) = 2.3 \text{ m } C_{v} \cdot \log\left(\frac{T_{2}}{T_{1}}\right)$

$$= 2.3 \times 1 \times 0.74 \log\left(\frac{303}{373}\right)$$

$$= -0.1536 \text{ kJ/kg K}$$
(-) indicates decrease in entropy.

EXAMPLE 3

300 kJ/s of heat is supplied at a constant fixed temperature of 290°C to a heat engine. The heat rejection takes place at 8.5°C. The following results were obtained.

(i) 150 kJ/sec are rejected

(ii) 215 kJ/sec are rejected

(iii) 75 kJ/sec are rejected.

Classify which of the result report a reversible, cyclic or irreversible cycle or impossible results. Heat supplied at $290^{\circ}C = 300 \text{ kJ/sec}$ and heat rejected at $8.5^{\circ}C$ (i) 150 kJ/sec (ii) 215 kJ/sec (iii) 75 kJ/sec.

SOLUTION

Applying Clausius inequality to the cycle or process

(i)
$$\sum_{\text{cyclic}} \frac{\delta Q}{T} = \frac{300}{290 + 273} - \frac{150}{8.5 + 273}$$

=0.5328 - 0.5328 = 0

: Cycle is reversible

(ii)
$$\sum_{\text{cyclic}} \frac{\delta Q}{T} = \frac{300}{290 + 273} - \frac{215}{8.5 + 273}$$

= 0.5328 - 0.7638 < 0

: Cycle is irreversible

(iii)
$$\sum_{\text{cyclic}} \frac{\delta Q}{T} = \frac{300}{290 + 273} - \frac{75}{8.5 + 273}$$

=0.5328 - 0.2664 = 0.2664 > 0

: Cycle is impossible by 2nd law of thermodynamics.

EXAMPLE 4

One kilogram of air is heated on constant volume from 100°C and 400°C. If $C_v = 0.7186$ kJ/kg K, determine the change in entropy of the air.

SOLUTION

Change of Entropy of air

$$(S_2 - S_1) = \int_{T_1}^{T_2} \frac{\delta Q}{T} = \int_{T_1}^{T_2} m C_v \frac{dT}{T}$$
$$= 2.3 m C_v \cdot \log\left(\frac{T_2}{T_1}\right)$$
$$= 2.3 \times 1 \times 0.7186 \times \log\left(\frac{400 + 273}{100 + 273}\right)$$
$$= 0.4236 \text{ kJ/K}$$

EXAMPLE 5

A Carnot power cycle is operated between 1200 K and 300 K. If 4.8 MJ of heat is supplied to the cycle, determine the heat rejected and the change in entropy for the process of heat rejection.

SOLUTION

For Carnot Cycle

$$\frac{Q_H}{T_H} = \frac{Q_L}{T_L} \text{ (reversible cycle)}$$
$$Q_L = \left(\frac{T_L}{T_H}\right) Q_H$$

:..

$$= 4.8 \times \frac{300}{1200} = 1.2 \text{ MJ}$$

Change of entropy for the process heat rejection

$$\Delta S = \frac{Q_L}{T_L} = \frac{1.2 \times 10^3}{300} = 4.0 \text{ kJ/K}$$



E	XAMPLE 6	
rev	A heat engir ersibly to a sink at 27°C	he receives reversible 420 kJ/cycle of heat from a source at 327°C and rejects heat C. There are no other heat transfers. For each of the three hypothetical amount of heat
reje	ection in (a), (b) and (c)	below. Compute the cycle integral of $\frac{\delta Q}{T}$ from this result. Show which is irrevers-
ible	which is reversible and	d which is impossible. (a) 210 kJ/cycle (b) 105 kJ/cycle (c) 315 kJ/cycle.
Soi	LUTION	
	Here,	$T_1 = 327^{\circ}\text{C}, Q_1 = 420 \text{ kJ/cycle}$
		$T_2 = 27^{\circ}\mathrm{C}$
(a)	As	$Q_2 = 210 \text{kJ/cycle}$
		$\oint \frac{\delta Q}{T} = \frac{420}{327 + 273} - \frac{210}{27 + 273}$
		$=\frac{420}{600}-\frac{210}{300}=0$
÷	Case is reversible	
(b)	As	$Q_2 = 105 \text{ kJ/cycle}$
		$\oint \frac{\delta Q}{T} = \frac{420}{327 + 273} - \frac{105}{27 + 273} = 0.35 > 0$
	Case is impossible	
(c)	As	$Q_3 = 315 \text{kJ/cycle}$
		$\oint \frac{\delta Q}{T} = \frac{420}{327 + 273} - \frac{315}{27 + 273}$
		=-0.35<0
	Case is irreversible	

EXAMPLE 7

An iron cube at a temperature of 400°C is dropped into an insulated bath containing 10 kg water at 25°C. The water finally reaches a temperature of 50°C at steady state. Given that the specific heat of water is equal to 4186 J/kg K. Find the entropy changes for the iron cube and the water. Is the process reversible? If so why?

SOLUTION

Given temperature of iron cube	$=400^{\circ}C=400+273=673$ K
Temperature of water	$=25^{\circ}C=25+273=298$ K
Mass of water	=10 kg

Temperature of water and cube after equilibrium

 $= 50^{\circ}C = 50 + 273 = 323 \text{ K}$ $C_{_{DW}} = 4186 \,\text{J/kg K}$ Specific heat of water, heat lost by iron cube = Heat gained by water Now, So, $m_i C_{Di}(673 - 323) = m_w C_{DW}(323 - 298)$ $=10 \times 4186(323 - 298)$ $m_i C_{pi} = \frac{10 \times 4186 (323 - 298)}{(673 - 323)} = 2990$ *:*..

Where $m_i =$ Mass of iron in kg, and

 C_{pi} = Specific heat of iron in J/kg K.

Entropy of iron at 673 K = $m_i C_{pi} \ln\left(\frac{673}{273}\right)$ $=2990 \ln \left(\frac{673}{273}\right)$ = 2697.8 J/K [Taking 0°C as datum]

Entropy of water at 298 K = $m_w C_{pw} \ln \left(\frac{298}{273}\right)$

$$= 10 \times 4186 \ln\left(\frac{298}{273}\right)$$

= 3667.8 J/K [Taking 0°C as datum]

Entropy of iron at 323 K = 2990 × ln $\left(\frac{323}{273}\right)$ = 502.8 J/K

Entropy of water at 323 K = $10 \times 4186 \ln \left(\frac{323}{273}\right) = 7040.04 \text{ J/K}$

Changes in entropy of iron = 502.8 - 2697.8 = -2195 J/KChange in entropy of water = 7040.04 - 3667.8 = 3372.24 J/KNet change in entropy $(\Delta S) = 3372.24 - 2195 = 1177.24 \text{ J/K.} > 0$

Hence, the process is irreversible.

EXAMPLE 8

An ideal gas is heated from temperature T_1 to T_2 by keeping its volume constant. The gas is expanded back to its initial temperature according to the law $p \forall^n = \text{constant}$. If the entropy change in the two processes are equal, find the value of n in terms of the adiabatic index g.

SOLUTION

Change in entropy during constant volume process

$$= mC_{v} \ln\left(\frac{T_{2}}{T_{1}}\right) \tag{i}$$

Change in entropy during polytropic process ($p \forall^n = \text{constant}$)

$$= mC_{\nu} \left(\frac{\gamma - n}{n - 1}\right) \ln\left(\frac{T_2}{T_1}\right)$$
(ii)

For the same entropy, we have

$$\begin{pmatrix} \frac{\gamma - n}{n - 1} \end{pmatrix} = 1$$

(\(\gamma - n\)) = (n - 1) or, 2n = \(\gamma + 1\)
$$n = \frac{\gamma + 1}{2}$$

or,

.:.

EXAMPLE 9

Air at 20°C and 1.05 bar occupies 0.025 m³. The air is heated at constant volume until the pressure is 4.5 bar, and then cooled at constant pressure back to original temperature. Calculate:

(i) the net heat flow form the air.

(ii) the net entropy change

Sketch the process on *T-S* diagram

Assume, $C_p = 1.005 \text{ kJ/kgK}$, $C_v = 0.718 \text{ kJ/kgK}$

SOLUTION

Forair

$$T_{1} = 20 + 273 = 293 \text{ K}$$

$$\forall_{1} = \forall_{3} = 0.025 \text{ m}^{3}$$

$$p_{1} = 1.05 \text{ bar} = 1.05 \times 10^{5} \text{ N/m}^{2}$$

$$p_{2} = 4.5 \text{ bar} = 4.5 \times 10^{5} \text{ N/m}^{2}$$

(i) Net heat flow:

For a perfect gas (corresponding to point 1 of air)

$$m = \left(\frac{p_1 \forall_1}{RT_1}\right) = \frac{1.05 \times 10^5 \times 0.025}{0.287 \times 10^3 \times 293} = 0.0312 \text{ kg}$$

For a perfect gas at constant volume

$$\frac{p_1}{T_1} = \frac{p_2}{T_2}$$
$$\frac{1.05}{293} = \frac{4.5}{T_2}$$

or

$$\Rightarrow$$
 $T_2 = \frac{4.5 \times 293}{1.05} = 1255.7 \text{ kg}$



At constant volume,

$$Q_{1-2} = mC_v(T_2 - T_1) = 0.0312 \times 0.718 (1255.7 - 293)$$

= 21.56 kJ

Also at constant pressure,

$$Q_{2-3} = mC_p (T_1 - T_2) = 0.0312 \times 1.005 (293 - 1255.7)$$

= -30.18 kJ

Net heat flow = $Q_{1-2} + Q_{2-3} = 21.56 + (-30.18) = -8.62 \text{ kJ}$ Heat rejected = 8.62 kJ

(ii) Net entropy change:

Referring to figure, net decrease in entropy

$$S_1 - S_3 = (S_2 - S_3) - (S_2 - S_1)$$

At constant pressure $\delta Q = mC_p dT$,

So,
$$S_2 - S_3 = \int_{293}^{1255.7} \frac{mC_p dT}{T} = 0.0312 \times 1.005 \ln\left(\frac{1255.7}{293}\right)$$
$$= 0.0456 \text{ kJ/K}$$

 $(S_1 - S_3) = (S_2 - S_3) - (S_2 - S_1)$

At constant volume $\delta Q = mC_v dT$,

$$S_2 - S_1 = \int_{293}^{1255.7} \frac{mC_v dT}{T} = 0.0312 \times 0.718 \ln\left(\frac{1255.7}{293}\right)$$

So,

$$= 0.0456 - 0.0326 = 0.013$$
 kJ/K

= 0.0326 kJ/K

Hence, decrease in entropy = 0.013 kJ/K.

EXAMPLE 10

0.04 kg of carbondioxide (moecular weight = 44) is compressed from 1 bar, 20°C, until the pressure is 9 bar, and the volume is then 0.003 m³. Calculate the change of entropy. Take C_p for carbon dioxide as 0.88 kJ/kg K, and assume carbon dioxide to be a perfect gas.

SOLUTION

Mass of carbondioxide, m = 0.04 kg Molecular weight, M = 44Initial pressure, $p_1 = 1$ bar = 1×10^5 N/m² Initial temperature $T_1 = 20 + 273 = 293$ K

Final pressure $p_2 = 9$ bar



Final volume C_p for carbondioxide Characteristics gas constant

 $R = \frac{R_u}{M} = \frac{8314}{44} = 189 \,\mathrm{Nm/kgK}$

To find T_2 , using the relation $p_2 \forall_2 = mRT_2$

:.

$$T_2 = \frac{p_2 \forall_2}{mR} = \frac{9 \times 10^5 \times 0.003}{0.04 \times 189} = 357 \text{ K}$$

Now

 $S_A - S_2 = R \ln\left(\frac{p_2}{p_1}\right) = \left(\frac{189}{10^3}\right) \ln(9/1) = 0.4153 \text{ kJ/kg K}$

 $\forall_2 = 0.003 \text{ m}^3$

= 0.88 kJ/kg K

Also at constant pressure from 1 to A

$$S_A - S_1 = C_p \ln \frac{T_2}{T_1} = 0.88 \ln \left(\frac{357}{293}\right) = 0.1738 \text{ kJ/kg K}$$

Then $(S_1 - S_2) = (S_A - S_2) - (S_A - S_1) = 0.4153 - 0.1738$
 $= 0.2415 \text{ kJ/kg K}$

Hence for 0.04 kg of carbondioxide, decrease in entropy

$$= m(S_1 - S_2) = 0.04 \times 0.2415 = 0.00966 \text{ kJ/k}$$

EXAMPLE 11

Calculate the change of entropy of 1 kg of air expanding polytropically in a cylinder behind a piston from 7 bar and 600°C to 1.05 bar. The index of expansion is 1.25.

SOLUTION

Here,

 $p_1 = 7 \text{ bar} = 7 \times 10^5 \text{ N/m}^2$ $T_1 = 600 + 273 = 873 \text{ K}$ $p_2 = 1.05 \text{ bar} = 1.05 \times 10^5 \text{ N/m}^2$

Index of expansion, n = 1.25Mass of air = 1 kg

To find T_2 , using the relation

$$\frac{T_2}{T_1} = \left(\frac{p_2}{p_1}\right)^{\frac{n-1}{n}}$$
$$\frac{T_2}{873} = \left(\frac{1.05}{7}\right)^{\frac{1.25-1}{1.25}}$$



 \Rightarrow

:.
$$T_2 = 873 \times \left(\frac{1.05}{7}\right)^{\frac{1.25-1}{1.25}} = 597.3 \text{ K}$$

Now replace the process 1 to 2 by process 1 to A and A to 2.

Then at constant temperature from 1 to A

$$S_A - S_1 = R \ln\left(\frac{\forall_2}{\forall_1}\right) = R \ln\left(\frac{p_1}{p_2}\right) = 0.544 \text{ kJ/kg K}$$

At constant temperature from A to 2

$$S_A - S_2 = C_p \ln\left(\frac{T_2}{T_1}\right) = 1.005 \ln\left(\frac{873}{597.3}\right) = 0.3814 \text{kJ/kg K}$$

Then

 $S_2 - S_1 = 0.544 - 0.3814 = 0.1626 \text{ kJ/kg K}$ Increase in entropy = 0.1626 kJ/kg K

EXAMPLE 12

In an air turbine the air expands from 7 bar and 460°C to 1.012 bar and 160°C. The heat loss from the turbine can be assumed to be negligible.

(i) show that the process is irreversible.

(ii) calculate the change of entropy per kg of air.

SOLUTION

Here,

$$p_1 = 7 \text{ bar} = 7 \times 10^5 \text{ N/m}^2$$

 $T_1 = 460 + 273 = 733 \text{ K}$
 $p_2 = 1.012 \text{ bar} = 1.012 \times 10^5 \text{ N/m}^2$
 $T_2 = 160 + 273 = 433 \text{ K}$

 $7 h_{ev} = 7 \times 105 N l_{ev}^{2}$

(i) Since heat loss is negligible, the process is adiabatic. For a reversible adiabatic process for a perfect gas,

$$\frac{T_2}{T_1} = \left(\frac{p_2}{p_1}\right)^{\frac{\gamma-1}{\gamma}}$$



...

...

$$\frac{T_2}{733} = \left(\frac{1.012}{7}\right)^{\frac{1.4-1}{1.4}}$$
$$T_2 = 421.8 \,\mathrm{K} = 148.8^{\circ}\mathrm{C}$$

But the actual temperature 160°C at 1.012 bar. Hence the process is irreversible.

(ii) Change of entropy can be found by cosidering a reversible constant pressure process between 2 and 2'.

:
$$S_2' - S_2 = C_p \ln\left(\frac{T_2}{T_2'}\right) = 1.005 \ln\left(\frac{433}{421.6}\right) = 0.02681 \text{ kJ/kg K}$$

Increase in entropy = 0.02681 kJ/kg K

EXAMPLE 13

A fluid undergoes a reversible adiabatic compression from 4 bar, 0.3 m³ to 0.08 m³ according to $p\forall^{1.25}$ = constant. Determine (i) change in enthalpy (ii) change in internal energy (iii) change in entropy (iv) heat transfer (v) work transfer.

(i)

SOLUTION

$$p_1 = 4 \text{ bar} = 4 \times 10^5 \text{ N/m}^2$$

 $\forall_1 = 0.3 \text{ m}^3, \forall_2 = 0.08 \text{ m}^3$

 $p_2 = p_1 \left(\frac{\forall_1}{\forall_2}\right)^n = 4 \left(\frac{0.3}{0.08}\right)^{1.25} = 20.87 \text{ bar}$

For reversible adiabatic process,

 $p_1 \forall_1^n = p_2 \forall_2^n$

(i) We know
$$\int_{H_1}^{H_2} dH = \int_{p_1}^{p_2} \forall dp$$
 (ii)
Also, $p_1 \forall_1^n = p \forall^n$
$$\therefore \qquad \forall = \left(\frac{p_1 \forall_1^n}{p}\right)^{1/n}$$



From equation (ii)

$$\int_{H_{1}}^{H_{2}} dH = \int_{p_{1}}^{p_{2}} \left(\frac{p_{1}\forall_{1}^{n}}{p}\right)^{1/n} dp$$

$$H_{2}-H_{1} = \left(p_{1}\forall_{1}^{n}\right)^{1/n} \left[\frac{p^{1-1/n}}{1-1/n}\right]_{p_{1}}^{p_{2}}$$

$$= \frac{n\left(p_{2}\forall_{2}-p_{1}\forall_{1}\right)}{n-1} \qquad [From equation (i)]$$

$$= \frac{1.25}{(1.25-1)\times10^{3}} \left[20.87\times10^{5}\times0.08-4\times10^{5}\times0.3\right] kJ$$

$$= 234.8 kJ$$

or,

: Change in enthalpy = 234.8 kJ
(ii)
$$H_2 - H_1 = (U_2 + p_2 \forall_2) - (U_1 + p_2 \forall_3)$$

$$H_2 - H_1 = (U_2 + p_2 \forall_2) - (U_1 + p_1 \forall_1)$$

= (U - U) + (p \forall - p \forall)

$$= (U_2 - U_1) + (p_2 \nabla_2 - p_1 \nabla_1)$$

$$U_2 - U_1 = (H_2 - H_1) - (p_2 \nabla_2 - p_1 \nabla_1)$$

$$= 234.8 - \left(\frac{20.87 \times 10^5 \times 0.08 - 4 \times 10^5 \times 0.3}{10^3}\right)$$

$$= 187.84 \text{ kJ}$$

Entropy

∴ Change in internal energy = 187.84 kJ (iii) Change in entropy $(S_2 - S_1) = 0$ (iv) Heat transfer, $Q_{1-2} = 0$ (v) Work transfer, W_{1-2} ∴ $W_{1-2} = Q_{1-2} - (U_2 - U_1) = 0 - 187.84 = -187.84$ kJ

EXAMPLE 14

3 kg of water at 80°C is mixed with 4 kg of water at 15°C in an isolated system. Calculate the change of entropy due to mixing process. Assume specific heat of water at constant pressure in 4.187 kJ/kgK.

SOLUTION

Applying first law of thermodynamics to the isolated system: Total energy before mixing = Total energy after mixing

$$\therefore 3C_{pw}(80-0) + 4C_{pw}(15-0) = 7C_{pw}(t_m - 0)$$

$$[C_{pw} = Specific heat of water at constant pressure]$$
or, $240C_{pw} + 60C_{pw} = 7C_{pw}t_m$
or, $240 + 60 = 7t_m$

:.

Initial entropy of the system

$$= 3C_{pw} \ln\left(\frac{80+273}{273}\right) + 4C_{pw} \ln\left(\frac{15+273}{273}\right)$$

$$=0.7709C_{pw}+0.2139C_{pw}=0.9848C_{pw}$$

 $t_m = \frac{300}{7} = 42.85^{\circ}\text{C}$

Final entropy of the system

$$= (3+4)C_{pw} \ln\left(\frac{42.85+273}{273}\right) = 1.0205C_{pw}$$

Net change in entropy,

$$\Delta S = \text{Final entropy} - \text{Initial entropy} = 1.0205 C_{pw} - 0.9848 C_{pw} = 0.0357 C_{pw} = 0.1495 \text{ kJ/K} [C_{pw} = 4.187 \text{ kJ/kg K}]$$

:. Change in entropy = 0.1495 kJ/K.

EXERCISE

- 1. What do you mean by Clausius inequality?
- 2. What do you mean by the term 'entropy'? What are the characteristics of entropy?



- 3. Prove that entropy is a property of a system.
- 4. What is a temperature–entropy diagram?
- 5. Derive the expressions for entropy changes for a closed system in the following cases: (i) general case for change of entropy of a gas (ii) heating a gas at constant volume (iii) heating a gas at constant pressure (iv) polytropic process.
- 6. 1 kg of ice at -5° C is exposed to the atmosphere, which is at 20°C. The ice melts and comes into thermal equilibrium with the atmosphere. Determine the increase of entropy in the universe. What is the minimum amount of work necessary to convert the water back into ice at -5° C. Take C_p of the ice 2.093 kJ/kg C and the latent heat of fusion of ice 333.33 kJ/kg.
- 7. 1 m³ of air is heated reversibly at constant pressure from 15°C to 300°C, and is then cooled reversibly at constant volume back to the initial temperature. If initial pressure is 1.03 bar, calculate the net heat flow and overall change of entropy.
- 8. 1 kg of air is allowed to expand reversibly in a cylinder behind a piston in such a way that the temperature remains constant at 260°C while the volume is doubled. The piston is then moved in, and heat is rejected by the air reversibly at constant pressure until the volume is the same as it was initially. Calculate the net heat flow and the overall change of entropy.
- 9. In reversible process, the rate of heat transfer to the system per unit temperature rise is given by 0.5 kJ/C. Find the change in entropy of the system if its temperature rises from 500 K to 800 K.
- 10. 0.03 m³ of nitrogen contained in a cylinder behind a piston is initially at 1.05 bar and 15°C. The gas is compressed isothermally and reversibly until the pressure is 4.2 bar. Calculate the change of entropy, the heat flow, and the work done and sketch the processes on pressure-volume diagram and temperature-entropy diagram. Assume nitrogen to act as a perfect gas and its molecular weight is 28.



PROPERTIES OF PURE SUBSTANCES

7.1 INTRODUCTION

All the thermodynamic devices employ a fluid as a medium of energy transport between the system and the surrounding. This fluid is known as the *working substance*. Substances may exist in various forms. A *phase* is any homogeneous form of a substance, that is, solid, liquid and gas.

A pure substance is one that is homogeneous and invariable in chemical aggregate. It may exist in one or more phase, but the chemical composition remains the same in all phases. Thus, solid water (ice), liquid water or water vapour (steam) are pure substance since chemical composition is the same in each phase. In the steam power plant, the working substance is water. It is vaporised to steam in the boiler. The steam flows through and run the turbine. It is condensed to liquid water in the condenser. The water is then returned to the boiler by a pump. The mixtures of water and steam or vapour are pure substance since chemical composition is the same in each phase.

In this chapter, we will consider a system comprising of a single component, that as a system in which only one pure substance is present. Our aim here is to study the phase in which this pure substance may exist, and also the conditions or states as specified by its properties under which it may exist in a particular space. Also in this chapter, we shall consider methods (Table, charts, equations etc.) for the presentation of thermodynamic properties of pure substances.

7.2 PHASE EQUILIBRIUM OF A PURE SUBSTANCE ON T-V DIAGRAM

Consider 1 kg of water ice at -40° C (-40° F) contained in piston-cylinder arrangement as shown in figure 7.1. The piston and weights maintain a constant pressure in the cylinder. If we heat the system from outside, change of state of water that take place inside the cylinder as a result of this heating. The change in state of water on heating and cooling at constant pressure on temperature, *i.e.*, specific volume diagram is shown in figure 7.2 in stepwise manner.



Step-1

Process S–*S*: At temperature 40°C, ice is at solid state *S*. The volume of ice will increase slightly with the increase in temperature as shown in figure 7.1(b). It will cause the piston to move slightly upwards. At the end of this step, the ice is in state *S* as in figure 7.2. It is sensible heating, in which the temperature of the substance change but the phase remain the same, is known as melting of ice into water.



Step-2

Process $S-f_i$: On further heating,

temperature remaining constant at 0°C, ice in state S melts to water in state f_1 . There is a change of phase from solid to liquid. In this process, there is a nominal decrease in the specific volume, since liquid water at 0°C is heavier than ice. Until all the ice melted to water as shown in figure 7.1(f) in between such as at A in figure 7.2. There is a mixture of solid and liquid in equilibrium as shown in figure 7.1(c).

It is seen from figure 7.2 that water-ice decrease in volume when changing from solid to liquid phase. It is latent heating, in which the phase changes, with temperature remaining constant.

Step-3

Process $f_1 - f_2$: On further heating causes the temperature of liquid water at f_1 at 0°C to rise. The temperature continues to rise until point f_2 at 100°C is reached. In this process, the specific volume of water at first decreases as the temperature rises to about 4°C, and then increases. A state of liquid water in between this process of heating, say *F* at 60°C, is shown in figure 7.1(e). After reaching 100°C temperature the state is shown in figure 7.1(f). Process $f_1 - f_2$ is sensible heating.

Step-4

Process f_2 -g: If heating is continued, water in state f_2 starts evaporating, the temperature remaining constant at 100°C. As a result, there is a large increase in specific volume of water as it changes from liquid phase to vapour phase. Until all the liquid water has evaporated to vapour at g as in figure 7.1(h). In between such as at B in figure 7.2. There is a mixture of liquid and vapour in equilibrium as shown in figure 7.1(g). So the process f_2 -g is *latent heating*, called vaporisation of liquid water into water vapour.

Step-5

Process g-G: On further heating of vapour at g causes its temperature to rise above 100°C, say to 350°C at G. The vapour at G, as shown in figure 7.1(i), occupies a greater volume than the vapour at g.

We see that there are two kinds of change of state taking place on addition of heat at constant pressure. These are *sensible heating* and *latent heating*. Thus, in figure 7.2, we have melting of ice

in water, (process *S*– f_1 at 0°C) and vaporisation of liquid water into water vapour, (process f_2 –g at 100°C).

If the process is carried in the opposite direction by cooling is shown in figure 7.2 starting from G, process G-g sensible cooling of water vapour, process $g-f_2$ condensation of steam, process f_2-f_1 sensible cooling of water, process f_1 -S is fusion of water into ice and finally process S-S is sensible cooling of ice.

7.3 TEMPERATURE AND TOTAL HEAT GRAPH DURING STEAM FORMATION

The process of steam formation, as discussed before, may also be represented on a graph, whose abscissa represents the total heat and the vertical ordinate represents the temperature. The point *A* represents the initial condition of water at 0°C and pressure *p*, in bar is shown in figure 7.3. The line *ABCD* shows the relation between temperature and heat at a specific pressure of p(in bar).

For formation of super heated steam from water at freezing point the heat is absorbed in the three stages.

1. Heating of water up to boiling temperature or saturation temperature *t* shown by *AB*. Heat absorbed by water is known as *sensible heat*.



- 2. Change of state from liquid to steam is shown by *BC*. Heat absorbed during this stage is *PQ* known as *latent heat of vaporisation*.
- 3. *CD* is the super heating process. Heat absorbed during this process is given by *QR* is known as heat of superheating. If pressure increases the boiling point also increases.

The line joining the point A, B, E, K is known as *saturated liquid line* which forms the boundary between water and steam. Similarly a line passing through any steam points L, F, C is known as *dry saturated steam line*, which forms boundary lines between wet and superheated steam. When the pressure and saturation temperature increases the latent heat of vaporisation decreases, it becomes zero at N, where liquid and dry state lines meet. This point is known as *critical point* and the temperature corresponding to the critical point N is known as *critical temperature* and the pressure as *critical pressure*.

For steam critical temperature = 374.15° C and critical pressure = 220.9 bar ≈ 221 bar.

7.4 PHASE EQUILIBRIUM AT HIGHER PRESSURE

Now increase the weights on the piston in figure 7.1, thus increasing the pressure on ice. We observe

similar changes of state taking place, is shown in figure 7.2, with the following difference shown in figure 7.4.

- (a) As the saturation pressure is increased, the saturation temperature t_{sat} for boiling and condensation is also increased.
- (b) As the saturation pressure is increased, the saturation temperature t_{sat} for melting and freezing is slightly decreased.
- (c) If the specific volume of saturated liquid water at f is v_f and saturated water vapour at g is v_g then the increase in specific volume $(v_g v_f)$ becomes smaller and smaller at higher and higher pressures.



In figure 7.4 at a pressure 220.9 bar shown by line $p = p_c$, increase in specific volume $(v_g - v_f)$ becomes zero, *i.e.*, $v_g = v_f$. At this pressure $v_g = v_f = 0.003155 \text{ m}^3/\text{kg}$ and there is no constant temperature vaporisation process. There is a point of inflexion at *C* with zero slope at a temperature of 374.15°C or 647.15 K and is called *critical point*. At the critical point the saturated liquid and saturated vapour states are identical. The pressure, temperature and specific volume at this point are named as critical pressure, critical temperature and critical volume. Critical data for some substance are given in table 7.1, with critical temperatures in decreasing order.

Substance	Critical Temperature (°C)	Critical Pressure (bar)	Critical Volume (m³/kg)
Mercury	>1550	>200	_
Water	374.3	220.9	0.003155
Ammonia	132	112.8	_
Carbon-dioxide	31.05	73.9	0.002143
Oxygen	-118.35	50.8	0.002438
Nitrogen	-147	33.99	_
Hydrogen	-239.85	13.0	0.032192
Helium	-267.9	2.29	_

Table: 7.1Critical Data for Some Substance

It is seen that there are no separate or distinct liquid and vapour phases at a pressure critical and above. Below the critical temperature and above critical pressure the fluid is considered as *compressed liquid* and as *superheated vapour* if the temperature is higher than critical.

7.5 PHASE EQUILIBRIUM AT LOWER PRESSURES

If we decreases the weights in the experiment of the art 7.2, so that pressure on the piston is below atmospheric pressure the saturation temperature for vaporisation decreases. The specific volume during phase change from liquid to vapour is more. The melting point of ice is slightly higher than 0°C. The line of vaporisation and melting processes come closer as shown in figure 7.4. When we go to lower and lower pressure, at a saturation pressure of 0.006113 bar, the line for vaporisation and melting process conjoin at a temperature of 0.01°C. At this temperature and pressure, we see that all the three phases of water solid, liquid and vapour can exist in equilibrium. This known as the *triple point state*. The straight line such as $s - f_1 - f_2 - g$ in figure 7.4 is called the *triple point line*.

Substance	Temperature (0°C)	Pressure (bar)
Copper	1083	0.000079
Silver	961	0.01
Zinc	419	5.066
Water	0.01	0.006113
Mercury	-39	0.00000013
Ammonium	-78	6.1
Carbon-dioxide	-52	517
Oxygen	-210	12.53
Nitrogen	-219	0.15
Hydrogen	-259	7.194
Helium	-271	0.5

Table: 7.2Triple Point Data for Some Substance

7.6 THERMODYNAMIC SURFACE

Graphical representation of the state of a pure substance which must have two independent properties, and any third as dependent one, can only be provided by a three dimensional plot. Such plots are called *thermodynamic surfaces*. The thermodynamics surfaces of a substance such as water which expands on freezing $p-\forall -T$ diagram is shown in figure 7.5. The constant temperature line, one

drawn through the critical point is called *critical isotherm*, and one with a phase transition. Figure 7.6 shows a $p - \forall -T$ surface for a normal substance which contracts on freezing.



Fig. 7.5

7.7 $p \rightarrow \forall$ DIAGRAM OF A PURE SUBSTANCES

The $p-\forall$ diagram of a pure substance is an another projection of the thermodynamic surface. It is very useful in analysing problem of thermodynamics. Figure 7.6 shows this diagram for a substance like water that expands on freezing.



7.8 IMPORTANT TERMS OF STEAM

(a) Wet steam

When the steam contains moisture or particle of water in suspension it is called wet steam. It means evaporation of water not completed and the whole of the latent heat has not been absorbed.

(b) Dry saturated steam

When the wet steam is further heated and it does not contain any suspended particles of water known as dry saturated steam. The dry saturated steam has absorbed its full latent heat and behaves practically in the same way as a perfect gas.

(c) Superheated steam

When the dry steam is further heated at a constant pressure, thus rising its temperature it is said to be superheated steam. Since the press is constant the volume of super heated steam increases. The volume of 1 kg of superheated steam is considerably greater than the volume of 1 kg of dry saturated steam at the same pressure.

(d) Dryness fraction or quality of wet steam

It is the ratio of the mass of actual dry steam to the mass of same quantity of wet steam. It denoted by 'x'

$$x = \frac{m_g}{m_g + m_f} = \frac{m_g}{m}$$

 m_g = Mass of actual dry stream m_f = Mass of water in suspension m = Mass of wet steam $= m_g + m_f$

(e) Sensible heat of water

Since specific heat of water at constant pressure = 4.2 kJ/kg K

:. Heat absorbed by 1 kg of water from 0° C to t° C (sensible heat)

$$= 1 \times 4.2 \{(t + 273) - (0 + 273)\}$$

= 4.2 t kJ/kg

Note: Sensible heat of water is taken equal to the specific enthalpy $= h_{f}$.

(f) Latent heat of vaporisation

It is the amount of heat absorbed to evaporate 1 kg of water at its boiling point or saturation temperature without change of temperature. It is denoted by change of temperature h_{fg} and the value depends on its pressure. Latent heat of steam or heat of vaporisation or water = 2257 kJ/kg at atmospheric pressure. If the steam is wet and dryness fraction is x, heat absorbed by it during evaporation is xh_{fg} .

(g) Enthalpy or total heat of steam

Enthalpy = Sensible heat + Latent heat

It is amount of heat absorbed by water from freezing point to saturation temperature plus the heat absorbed during evaporation denoted by h_g and the value of dry saturation steam may be read directly from the steam table.

Enthalpy of (i) wet steam $h = h_f + xh_{fg.}$ (ii) dry steam $h = h_g = h_f + h_{fg}$ (iii) superheated steam

x = dryness fraction in case of dry steam x = 1

If we further add heat to the dry steam its temperature increases while the pressure remains constant. The increase in steam temperature shows the super heat stage of the steam

 $h_{sup} = h_f + h_{fg} + C_p (t_{sup} - t)$ where, t_{sup} = temperature other superheated steam t = saturation temperature at constant pressure $t_{sup} - t$ = degree of superheat C_p of steam lies between 1.67 kJ/kg K to 2.5 kJ/kg K.

(h) Specific volume of steam

It is the volume occupied by the steam per unit mass at a given temperature and pressure and is expressed in m^3/kg . It is reciprocal of the density of steam in kg/m^3 . Specific volume decreases as the increases in pressure.

(i) Wet steam

Consider 1 kg of wet steam of dryness fraction x, *i.e.*, this steam will have x kg of dry steam and (1-x) kg of water. Let v_r be the vol of 1 kg of water then, volume of one kg of wet steam

$$= xv_{g} + (1-x)v_{g}$$

 $\therefore v_f$ is very small as compared to v_a

 \therefore (1-x) v_f may be rejected

So, volume of 1 kg of wet steam = $xv_g m^3$

Specific volume of wet steam $v = xv_{\sigma} m^3/kg$.

(ii) Dry steam

We know in case of dry steam mass of water in suspension is zero and dryness fraction = 1

So, specific volume of dry steam = $v_o m^3/kg$

(iii) Superheated steam

When the dry saturated steam is further heated under a constant pressure, there is no increase in volume with the rise in temperature. The super heated steam behaves like a perfect gas.

: According to Charle's law

$$\frac{v_{\sup}}{T_{\sup}} = \frac{v_g}{T}$$
 So, $v_{\sup} = \frac{v_g \cdot T_{\sup}}{T}$

 v_{sup} = specific volume of super heated steam

 v_g = specific volume of dry steam at the pressure of steam formation

- T_{sup} = absolute temperature of super heated steam.
- T^{\dagger} = absolute temperature at the pressure of steam for motion.

emperature in °C	Absolute pressure	Specific in m ³	volume //kg	Sp	ecific enthalpy in kJ/kg		S	pecific entropy in kJ/kg		Temperature in °C
	in bar	Water	Steam	Water	Evaporation	Steam	Water	Evaporation	Steam	
(t)	(p)	(V _f)	(V _g)	(n _f)	(n _{fg})	(n _g)	(s _f)	(S _{fg})	(s _g)	(t)
0	0.00611	0.001000	206.31	0.0	2501.6	2501.6	0.000	9.158	9.158	0
1	0.00657	0.001000	192.61	4.2	2499.2	2503.4	0.015	9.116	9.131	1
2	0.00706	0.001000	176.92	8.4	2496.8	2505.2	0.031	9.074	9.105	2
3	0.00758	0.001000	168.17	12.6	2494.5	2507.1	0.046	9.033	9.079	3
4	0.00813	0.001000	157.27	16.8	2492.1	2508.9	0.061	8.992	9.053	4
5	0.00872	0.001000	147.16	21.0	2489.7	2510.7	0.076	8.951	9.027	5
6	0.00935	0.001000	137.78	25.2	2487.4	2512.6	0.091	8.911	9.002	6
7	0.01001	0.001000	129.06	29.4	2485.0	2514.4	0.106	8.870	8.976	7
8	0.01072	0.001000	120.97	33.6	2482.6	2516.2	0.121	8.830	8.951	8
9	0.01147	0.001000	113.44	37.8	2480.3	2518.1	0.136	8.791	8.927	9
10	0.01227	0.001000	106.43	42.0	2477.9	2519.9	0.151	8.751	8.902	10
11	0.01312	0.001000	99.909	46.2	2475.5	2521.7	0.166	8.712	8.878	11
12	0.01401	0.001000	93.835	50.4	2473.2	2523.6	0.181	8.673	8.854	12
13	0.01497	0.001001	88.176	54.6	2470.8	2525.4	0.195	8.635	8.830	14
14	0.01597	0.001001	82.900	58.7	2468.5	2527.2	0.210	8.597	8.806	14
15	0.01704	0.001001	77.978	62.9	2466.1	2529.1	0.224	8.559	8.783	15
16	0.01817	0.001001	73.384	67.1	2463.8	2530.9	0.239	8.520	8.759	16
17	0.01936	0.001001	69.095	71.3	2461.4	2532.7	0.253	8.483	8.736	17
18	0.02062	0.001001	65.087	75.5	2459.0	2534.5	0.268	8.446	8.714	18
19	0.02196	0.001002	61.341	79.7	2456.7	2536.4	0.282	8.409	8.691	19

(t)	(p)	(v _f)	(v _g)	(h _f)	(h _{fg})	(h _g)	(s _f)	(s_{fg})	(s _g)	(t)
20	0.02337	0.001002	57.838	83.9	2454.3	2538.2	0.296	8.372	8.668	20
21	0.02485	0.001002	54.561	88.0	2452.0	2540.0	0.310	8.336	8.646	21
22	0.02642	0.001002	51.492	92.2	2449.6	2541.8	0.325	8.299	8.624	22
23	0.02808	0.001002	48.619	96.4	2447.2	2543.6	0.339	8.263	8.602	23
24	0.02982	0.001002	45.926	100.6	2444.9	2545.5	0.353	8.228	8.581	24
25	0.03166	0.001003	43.402	104.8	2442.5	2547.3	0.367	8.192	8.559	25
26	0.03360	0.001003	41.034	108.9	2440.2	2549.1	0.381	8.157	8.538	26
27	0.03564	0.001003	38.813	113.1	2437.8	2550.9	0.395	8.122	8.517	27
28	0.03778	0.001004	36.728	117.3	2435.4	2552.7	0.409	8.087	8.496	28
29	0.04004	0.001004	34.769	121.5	2433.1	2554.5	0.423	8.052	8.475	29
30	0.04242	0.001004	32.929	125.7	2430.7	2556.4	0.437	8.018	8.455	30
31	0.04491	0.001005	31.199	129.8	2428.3	2558.2	0.450	7.984	8.434	31
32	0.04753	0.001005	29.572	134.0	2425.9	2560.0	0.464	7.950	8.414	32
33	0.05029	0.001005	28.042	138.2	2423.6	2561.8	0.478	7.916	8.394	33
34	0.05318	0.001006	26.601	142.4	2421.2	2563.6	0.491	7.883	8.374	34
35	0.05622	0.001006	25.245	146.6	2418.8	2565.4	0.505	7.849	8.354	35
36	0.05940	0.001006	23.967	150.7	2416.4	2567.2	0.518	7.817	8.335	36
37	0.06274	0.001007	22.763	154.9	2414.1	2569.0	0.532	7.783	8.315	37
38	0.06624	0.001007	21.627	159.1	2411.7	2570.8	0.545	7.751	8.296	38
39	0.06991	0.001007	20.557	163.3	2409.3	2572.6	0.559	7.718	8.277	39
40	0.07375	0.001008	19.546	167.5	2406.9	2574.4	0.572	7.686	8.258	40
41	0.07777	0.001008	18.592	171.6	2404.5	2576.2	0.585	7.654	8.239	41
42	0.08199	0.001009	17.692	175.8	2402.1	2577.9	0.599	7.622	8.221	42
43	0.08639	0.001009	16.841	180.0	2399.7	2579.7	0.612	7.591	8.203	43
44	0.09100	0.001009	16.036	184.2	2397.3	2581.5	0.625	7.559	8.184	44
45	0.09582	0.001010	15.276	188.4	2394.9	2583.3	0.638	7.528	8.166	45
46	0.10080	0.001010	14.557	192.5	2392.5	2585.1	0.651	7.497	8.148	46
47	0.10612	0.001011	13.877	196.7	2390.1	2586.9	0.664	7.466	8.130	47
48	0.11162	0.001011	13.233	200.9	2387.7	2588.6	0.678	7.435	8.113	48
49	0.11736	0.001012	12.623	205 1	2385 3	2590.4	0 691	7 404	8 095	49

Mechanical Science-II

							1			
(t)	(p)	(v _f)	(v _g)	(h _f)	$(\mathbf{h}_{\mathrm{fg}})$	(h _g)	(s _f)	(s _{fg})	(s _g)	(t)
50	0.12335	0.001012	12.046	209.3	2382.9	2592.2	0.704	7.374	8.078	50
51	0.12961	0.001013	11.499	213.4	2380.5	2593.9	0.716	7.344	8.060	51
52	0.13611	0.001013	10.980	217.6	2378.1	2595.7	0.729	7.314	8.043	52
53	0.14293	0.001014	10.488	221.8	2375.7	2597.5	0.742	7.284	8.026	53
54	0.15002	0.001014	10.022	226.0	2373.2	2599.2	0.755	7.254	8.009	54
55	0.15741	0.001015	9.5789	230.2	2370.8	2601.0	0.768	7.225	7.993	55
56	0.16511	0.001015	9.1587	234.3	2368.4	2602.7	0.780	7.196	7.976	56
57	0.17313	0.001016	8.7598	238.5	2366.0	2604.5	0.793	7.166	7.959	57
58	0.18147	0.001016	8.3808	242.7	2363.5	2606.2	0.806	7.137	7.943	58
59	0.19016	0.001017	8.0208	246.9	2361.1	2608.0	0.818	7.109	7.927	59
60	0.19920	0.001017	7.6785	251.1	2358.6	2609.7	0.831	7.080	7.911	60
61	0.20861	0.001018	7.3532	255.3	2356.1	2611.4	0.844	7.051	7.895	61
62	0.21838	0.001018	7.0437	259.5	2353.7	2613.2	0.856	7.023	7.879	62
63	0.22855	0.001019	6.7493	263.6	2351.3	2614.9	0.868	6.995	7.863	63
64	0.23912	0.001019	6.4690	267.8	2348.8	2616.6	0.881	6.967	7.848	64
65	0.25009	0.001020	6.2023	272.0	2346.4	2618.4	0.893	6.939	7.832	65
66	0.26150	0.001020	5.9482	276.2	2343.9	2620.1	0.906	6.911	7.817	66
67	0.27334	0.001021	5.7062	280.4	2341.4	2621.8	0.918	6.884	7.802	67
68	0.28563	0.001022	5.4756	284.6	2338.9	2623.5	0.930	6.856	7.786	68
69	0.29838	0.001022	5.2558	288.8	2336.4	2625.2	0.943	6.828	7.771	69
70	0.31162	0.001023	5.0463	293.0	2333.9	2626.9	0.955	6.802	7.757	70
71	0.32535	0.001024	4.8464	297.2	2331.4	2628.6	0.967	6.775	7.742	71
72	0.33958	0.001024	4.6557	301.3	2329.0	2630.3	0.979	6.748	7.727	72
73	0.35434	0.001025	4.4737	305.5	2326.5	2632.0	0.991	6.721	7.712	73
74	0.36964	0.001025	4.3000	309.7	2324.0	2633.7	1.003	6.695	7.698	74
75	0.38549	0.001026	4.1341	313.9	2321.5	2635.4	1.015	6.668	7.683	75
76	0.40191	0.001027	3.9757	318.1	2318.9	2637.0	1.027	6.642	7.669	76
77	0.41891	0.001027	3.8243	322.3	2316.4	2638.7	1.039	6.616	7.655	77
78	0.43652	0.001028	3.6796	326.5	2313.9	2640.4	1.051	6.590	7.641	78
79	0.45474	0.001029	3.5413	330.7	2311.4	2642.1	1.063	6.564	7.627	79

Properties of Pure Substances

(t)	(s _g)	(s _{fg})	(s _f)	(h _g)	(h _{fg})	(h _f)	(v _g)	(v _f)	(p)	(t)
80	7.613	6.538	1.075	2643.8	2308.9	334.9	3.4091	0.001029	0.47360	80
81	7.599	6.512	1.087	2645.4	2306.3	339.1	3.2826	0.001030	0.49311	81
82	7.586	6.487	1.099	2647.1	2303.8	343.3	3.1616	0.001031	0.51329	82
83	7.572	6.461	1.111	2648.7	2301.2	347.5	3.0458	0.001031	0.53416	83
84	7.559	6.436	1.123	2650.4	2298.7	351.7	2.9350	0.001032	0.55573	84
85	7.545	6.411	1.134	2652.0	2296.1	355.9	2.8288	0.001033	0.57803	85
86	7.532	6.386	1.146	2653.6	2293.5	360.1	2.7272	0.001033	0.60108	86
87	7.519	6.361	1.158	2655.3	2291.0	364.3	2.6298	0.001034	0.62489	87
88	7.506	6.337	1.169	2656.9	2288.4	368.5	2.5365	0.001035	0.64948	88
89	7.493	6.312	1.181	2658.5	2285.8	372.7	2.4470	0.001035	0.67487	89
90	7.480	6.287	1.193	2660.1	2283.2	376.9	2.3613	0.001036	0.70109	90
91	7.467	6.263	1.204	2661.7	2280.6	381.1	2.2791	0.001037	0.72815	91
92	7.454	6.238	1.216	2663.4	2278.0	385.4	2.2002	0.001038	0.75606	92
93	7.442	6.215	1.227	2665.0	2275.4	389.6	2.1245	0.001038	0.78489	93
94	7.429	6.190	1.239	2666.6	2272.8	393.8	2.0519	0.001039	0.81461	94
95	7.417	6.167	1.250	2668.1	2270.1	398.0	1.9822	0.001040	0.84526	95
96	7.404	6.143	1.261	2669.7	2267.5	402.2	1.9153	0.001041	0.87686	96
97	7.392	6.119	1.273	2671.3	2264.9	406.4	1.8510	0.001041	0.90944	97
98	7.380	6.096	1.284	2672.9	2262.3	410.6	1.7893	0.001042	0.94301	98
99	7.368	6.072	1.296	2674.4	2259.6	414.8	1.7300	0.001043	0.97761	99
100	7.355	6.048	1.307	2676.0	2256.9	419.1	1.6730	0.001044	1.0133	100
102	7.331	6.002	1.329	2679.1	2251.6	427.5	1.5655	0.001045	1.0876	102
104	7.308	5.956	1.352	2682.2	2246.3	435.9	1.4662	0.001047	1.1668	104
106	7.284	5.910	1.374	2685.3	2240.9	444.4	1.3742	0.001048	1.2504	106
108	7.261	5.865	1.396	2688.3	2235.4	452.9	1.2889	0.001050	1.3390	108
110	7.239	5.821	1.418	2691.3	2230.0	461.3	1.2099	0.001052	1.4327	110
112	7.216	5.776	1.440	2694.3	2224.5	469.8	1.1366	0.001054	1.5316	112
114	7.194	5.732	1.462	2697.2	2218.9	478.3	1.0685	0.001055	1.6362	114
116	7.172	5.688	1.484	2700.2	2213.5	486.7	1.0052	0.001057	1.7465	116
118	7.151	5.645	1.506	2703.1	2207.9	495.2	0.94634	0.001059	1.8628	118

(t)	(p)	(v _f)	(v _g)	(h _f)	(h _{fg})	(h _g)	(s _f)	(s_{fg})	(s _g)	(t)
120	1.9854	0.001061	0.89152	503.7	2202.3	2706.0	1.528	5.601	7.129	120
122	2.1145	0.001063	0.84045	512.2	2196.6	2708.8	1.549	5.559	7.108	122
124	2.2504	0.001064	0.79283	520.7	2190.9	2711.6	1.570	5.517	7.087	124
126	2.3933	0.001066	0.74840	529.2	2185.2	2714.4	1.592	5.475	7.067	126
128	2.5435	0.001068	0.70691	537.8	2179.4	2717.2	1.613	5.433	7.046	128
130	2.7013	0.001070	0.66814	546.3	2173.6	2719.9	1.634	5.392	7.026	130
132	2.8670	0.001072	0.63188	554.8	2167.8	2722.6	1.655	5.351	7.006	132
134	3.0407	0.001074	0.59795	563.4	2161.9	2725.3	1.676	5.310	6.986	134
136	3.2229	0.001076	0.56618	572.0	2155.9	2727.9	1.697	5.270	6.967	136
138	3.4138	0.001078	0.53641	580.5	2150.0	2730.5	1.718	5.229	6.947	138
140	3.6139	0.001080	0.50849	589.1	2144.0	2733.1	1.739	5.189	6.928	140
142	3.8231	0.001082	0.48230	597.7	2137.9	2735.6	1.760	5.150	6.910	142
144	4.0420	0.001084	0.45771	606.3	2131.8	2738.1	1.780	5.111	6.891	144
146	4.2709	0.001086	0.43460	614.9	2125.7	2740.6	1.801	5.071	6.872	146
148	4.5101	0.001089	0.41288	623.5	2119.5	2743.0	1.821	5.033	6.854	148
150	4.7600	0.001091	0.39245	632.2	2113.2	2745.4	1.842	4.994	6.836	150
155	5.4333	0.001096	0.34644	653.8	2097.4	2751.2	1.892	4.899	6.791	155
160	6.1806	0.001102	0.30676	675.5	2081.2	2756.7	1.943	4.805	6.748	160
165	7.0077	0.001108	0.27240	697.2	2064.8	2762.0	1.992	4.713	6.705	165
170	7.9202	0.001114	0.24255	719.1	2048.0	2767.1	2.042	4.621	6.663	170
175	8.9244	0.001121	0.21654	741.1	2030.7	2771.8	2.091	4.531	6.622	175
180	10.027	0.001128	0.19380	763.1	2013.2	2776.3	2.139	4.443	6.582	180
185	11.233	0.001135	0.17386	785.3	1995.1	2780.4	2.187	4.355	6.542	185
190	12.551	0.001142	0.15632	807.5	1976.8	2784.3	2.236	4.268	6.504	190
195	13.987	0.001149	0.14084	829.9	1957.9	2787.8	2.283	4.182	6.465	195
200	15.549	0.001156	0.12716	852.4	1938.5	2790.9	2.331	4.097	6.428	200
205	17.243	0.001164	0.11503	875.0	1918.8	2793.8	2.378	4.013	6.391	205
210	19.077	0.001172	0.10424	897.7	1898.5	2796.2	2.425	3.929	6.354	210
215	21.060	0.001181	0.094625	920.6	1 877.7	2798.3	2.471	3.846	6.317	215
220	23.198	0.001190	0.086038	943.7	1856.2	2799.9	2.518	3.764	6.282	220

(t)	(s _g)	(s _{fg})	(s _f)	(h _g)	(h _{fg})	(h _f)	(v _g)	(v _f)	(p)	(t)
225	6.246	3.682	2.564	2801.2	1834.3	966.9	0.078349	0.001199	25.501	225
230	6.211	3.601	2.610	2802.0	1811.7	990.3	0.071450	0.001209	27.976	230
235	6.175	3.519	2.656	2802.3	1788.5	1013.8	0.065245	0.001219	30.632	235
240	6.141	3.439	2.702	2802.2	1764.6	1037.6	0.059654	0.001229	33.478	240
245	6.106	3.358	2.748	2801.6	1740.0	1061.6	0.054606	0.001240	36.523	245
250	6.071	3.277	2.794	2800.4	1714.6	1085.8	0.050037	0.001251	39.776	250
255	6.036	3.197	2.839	2798.7	1688.5	1110.2	0.045896	0.001263	43.246	255
260	6.001	3.116	2.885	2796.4	1661.5	1134.9	0.042134	0.001276	46.943	260
265	5.966	3.035	2.931	2793.5	1633.6	1159.9	0.038710	0.001289	50.877	265
270	5.930	2.954	2.976	2789.9	1604.7	1185.2	0.035588	0.001303	55.058	270
275	5.895	2.873	3.022	2785.5	1574.7	1210.8	0.032736	0.001317	59.496	275
280	5.858	2.790	3.068	2780.4	1543.6	1236.8	0.030126	0.001332	64.202	280
285	5.822	2.707	3.115	2774.5	1511.3	1263.2	0.027733	0.001349	69.186	285
290	5.785	2.624	3.161	2767.6	1477.6	1290.0	0.025535	0.001366	74.461	290
295	5.747	2.539	3.208	2759.8	1442.5	1317.3	0.023513	0.001384	80.037	295
300	5.708	2.453	3.255	2751.0	1406.0	1345.0	0.021649	0.001404	85.927	300
305	5.669	2.366	3.303	2741.1	1367.7	1373.4	0.019927	0.001425	92.144	305
310	5.628	2.277	3.351	2730.0	1327.6	1402.4	0.018334	0.001448	98.700	310
315	5.586	2.186	3.400	2717.6	1285.5	1432.1	0.016856	0.001473	105.61	315
320	5.542	2.092	3.450	2703.7	1241.1	1462.6	0.015480	0.001500	112.89	320
325	5.497	1.996	3.501	2688.0	1194.0	1494.0	0.014195	0.001529	120.56	325
330	5.449	1.896	3.553	2670.2	1143.7	1526.5	0.012989	0.001562	128.63	330
335	5.398	1.792	3.606	2649.7	1089.5	1560.2	0.011854	0.001598	137.12	335
340	5.343	1.681	3.662	2626.2	1030.7	1595.5	0.010780	0.001639	146.05	340
345	5.283	1.564	3.719	2598.9	966.4	1632.5	0.097631	0.001686	155.45	345
350	5.218	1.438	3.780	2567.7	895.8	1671.9	0.087991	0.001741	165.35	350
355	5.144	1.295	3.849	2530.4	813.8	1716.6	0.078592	0.001809	175.77	355
360	5.060	1.139	3.921	2485.4	721.2	1764.2	0.069398	0.001896	186.75	360
365	4.958	0.956	4.002	2428.0	610.0	1818.0	0.060116	0.002016	198.33	365
370	4.814	0.703	4.111	2342.8	452.6	1890.2	0.049728	0.002214	210.54	370
374.15	4.443	0.000	4.443	2107.4	0.0	2107.4	0.003170	0.003170	221.20	374.15

Absolute pressure	e Temperature in °C	Specific in m	volume ³ /kg	Sp	Specific enthalpy in kJ/kg			Specific entropy in kJ/kg K			
(p)	(t)	Water (v _f)	Steam (v _g)	Water (h _f)	Evaporation (h _{fg})	Steam (h _g)	Water (s _f)	Evaporation (s _{fg})	Steam (s _g)	(p)	
0.0061	0.000	0.001000	206.31	0.0	2501.6	2501.6	0.000	9.158	9.158	0.0061	
0.010	6.983	0.001000	129.21	29.3	2485.1	2514.4	0.106	8.871	8.977	0.010	
0.015	13.04	0.001001	87.982	54.7	2470.8	2525.5	0.196	8.634	8.830	0.015	
0.020	17.51	0.001001	67.006	73.5	2460.1	2533.6	0.261	8.464	8.725	0.020	
0.025	21.10	0.001002	54.256	88.4	2451.8	2540.2	0.312	8.333	8.645	0.025	
0.030	24.10	0.001003	45.667	101.0	2444.6	2545.6	0.354	8.224	8.578	0.030	
0.035	26.69	0.001003	39.479	111.8	2438.6	2550.4	0.391	8.132	8.523	0.035	
0.040	28.98	0.001004	34.802	121.4	2433.1	2554.5	0.423	8.053	8.476	0.040	
0.045	31.03	0.001005	31.141	130.0	2428.2	2558.2	0.451	7.983	8.434	0.045	
0.050	32.90	0.001005	28.194	137.8	2423.8	2561.6	0.476	7.920	8.396	0.050	
0.060	36.18	0.001006	23.741	151.5	2416.0	2567.5	0.521	7.810	8.331	0.06	
0.070	39.03	0.001007	20.531	163.4	2409.2	2572.6	0.559	7.718	8.277	0.07	
0.080	41.53	0.001008	18.105	173.9	2403.2	2577.1	0.593	7.637	8.230	0.08	
0.090	43.79	0.001009	16.204	183.3	2397.8	2581.1	0.622	7.566	8.188	0.09	
0.100	45.83	0.001010	14.675	191.8	2392.9	2584.7	0.649	7.502	8.151	0.10	
0.11	47.71	0.001011	13.416	199.7	2388.4	2588.1	0.674	7.444	8.118	0.11	
0.12	49.45	0.001012	12.362	206.9	2384.3	2591.2	0.696	7.391	8.087	0.12	
0.13	51.06	0.001013	11.466	213.7	2380.3	2594.0	0.717	7.342	8.059	0.13	
0.14	52.57	0.001013	10.694	220.0	2376.7	2596.7	0.737	7.296	8.033	0.14	
0.15	54.00	0.001014	10.023	226.0	2373.2	2599.2	0.755	7.254	8.009	0.15	

(p)	(s _g)	$(\mathbf{s}_{\mathrm{fg}})$	(s _f)	(h _g)	(h _{fg})	(h _f)	(v _g)	(v _f)	(t)	(p)
0.16	7.987	7.215	0.772	2601.6	2370.0	231.6	9.4331	0.001015	55.34	0.16
0.17	7.966	7.178	0.788	2603.8	2366.9	236.9	8.9111	0.001015	56.62	0.17
0.18	7.946	7.142	0.804	2605.9	2363.9	242.0	8.4452	0.001016	57.83	0.18
0.19	7.927	7.109	0.818	2607.9	2361.1	246.8	8.0272	0.001017	58.98	0.19
0.20	7.909	7.077	0.832	2609.9	2358.4	251.5	7.6498	0.001017	60.09	0.20
0.21	7.892	7.047	0.845	2611.7	2355.8	255.9	7.3073	0.001018	61.15	0.21
0.22	7.876	7.018	0.858	2613.5	2353.4	260.1	6.9951	0.001018	62.16	0.22
0.23	7.861	6.991	0.870	2615.2	2351.0	264.2	6.7093	0.001019	63.14	0.23
0.24	7.846	6.964	0.882	2616.8	2348.6	268.2	6.4467	0.001019	64.08	0.24
0.25	7.832	6.939	0.893	2618.3	2346.3	272.0	6.2045	0.001020	64.99	0.25
0.26	7.819	6.915	0.904	2619.9	2344.2	275.7	5.9803	0.001020	65.87	0.26
0.27	7.806	6.891	0.915	2621.3	2342.1	279.2	5.7724	0.001021	66.72	0.27
0.28	7.793	6.868	0.925	2622.7	2340.0	282.7	5.5778	0.001021	67.55	0.28
0.29	7.781	6.847	0.935	2624.1	2338.1	286.0	5.3982	0.001022	68.35	0.29
0.30	7.769	6.825	0.944	2625.4	2336.1	289.3	5.2293	0.001022	69.12	0.30
0.32	7.747	6.785	0.962	2628.0	2332.4	295.6	4.9220	0.001023	70.62	0.32
0.34	7.727	6.747	0.980	2630.4	2328.9	301.5	4.6504	0.001024	72.03	0.34
0.36	7.707	6.711	0.996	2632.6	2325.5	307.1	4.4076	0.001025	73.37	0.36
0.38	7.688	6.677	1.011	2634.8	2322.3	312.5	4.1900	0.001026	74.66	0.38
0.40	7.671	6.645	1.026	2636.9	2319.2	317.7	3.9934	0.001027	75.89	0.40
0.42	7.654	6.614	1.040	2638.9	2316.3	322.6	3.8148	0.001027	77.06	0.42
0.44	7.638	6.584	1.054	2640.7	2313.4	327.3	3.6522	0.001028	78.19	0.44
0.46	7.623	6.556	1.067	2642.6	2310.7	331.9	3.5032	0.001029	79.28	0.46
0.48	7.609	6.530	1.079	2644.3	2308.0	336.3	3.3663	0.001029	80.33	0.48
0.50	7.595	6.504	1.091	2646.0	2305.4	340.6	3.2401	0.001030	81.35	0.50
0.52	7.581	6.478	1.103	2647.6	2302.9	344.7	3.1233	0.001031	82.33	0.52
0.54	7.569	6.455	1.114	2649.2	2300.5	348.7	3.0148	0.001031	83.28	0.54
0.56	7.556	6.431	1.125	2650.7	2298.2	352.5	2.9139	0.001032	84.19	0.56
0.58	7.544	6.409	1.135	2652.1	2295.8	356.3	2.8197	0.001033	85.09	0.58
0.60	7.533	6.388	1.145	2653.6	2293.7	359.9	2.7317	0.001033	85.95	0.60

		1					1			
(p)	(t)	(v _f)	(v _g)	(h _f)	(h _{fg})	(h _g)	(s _f)	(s _{fg})	(s _g)	(p)
0.62	86.80	0.001034	2.6491	363.5	2291.4	2654.9	1.155	6.367	7.522	0.62
0.64	87.62	0.001034	2.5715	366.9	2289.4	2656.3	1.165	6.346	7.511	0.64
0.66	88.42	0.001035	2.4985	370.3	2287.3	2657.6	1.174	6.326	7.500	0.66
0.68	89.20	0.001036	2.4297	373.6	2285.2	2658.8	1.183	6.307	7.490	0.68
0.70	89.96	0.001036	2.3647	376.8	2283.3	2660.1	1.192	6.288	7.480	0.70
0.72	90.70	0.001037	2.3031	379.9	2281.4	2661.3	1.201	6.270	7.471	0.72
0.74	91.43	0.001037	2.2448	382.9	2279.5	2662.4	1.209	6.253	7.462	0.74
0.76	92.14	0.001038	2.1895	385.9	2277.7	2663.6	1.217	6.235	7.452	0.76
0.78	92.83	0.001038	2.1369	388.9	2275.8	2664.7	1.225	6.219	7.444	0.78
0.80	93.51	0.001039	2.0869	391.7	2274.1	2665.8	1.233	6.202	7.435	0.80
0.85	95.15	0.001040	1.9721	398.6	2269.8	2668.4	1.252	6.163	7.415	0.85
0.90	96.71	0.001041	1.8691	405.2	2265.7	2670.9	1.270	6.125	7.395	0.90
0.95	98.20	0.001042	1.7771	411.5	2261.7	2673.2	1.287	6.091	7.378	0.95
1.00	99.63	0.001043	1.6938	417.5	2257.9	2675.4	1.303	6.057	7.360	1.00
1.01325	100.00	0.001044	1.6730	419.1	2256.9	2676.0	1.307	6.048	7.355	1.01325
1.05	101.0	0.001045	1.6181	423.3	2254.3	2677.6	1.308	6.025	7.343	1.05
1.10	102.3	0.001046	1.5492	428.8	2250.8	2679.6	1.333	5.995	7.328	1.10
1.15	103.6	0.001047	1.4861	434.2	2247.4	2681.6	1.347	5.966	7.313	1.15
1.20	104.8	0.001048	1.4281	439.3	2244.1	2683.4	1.361	5.937	7.298	1.20
1.25	106.0	0.001049	1.3746	444.4	2240.8	2685.2	1.374	5.911	7.285	1.25
1.30	107.1	0.001050	1.3250	449.2	2237.8	2687.0	1.387	5.885	7.272	1.30
1.35	108.2	0.001050	1.2791	453.4	2234.8	2688.7	1.399	5.860	7.259	1.35
1.40	109.3	0.001051	1.2363	458.4	2231.9	2690.3	1.411	5.836	7.247	1.40
1.45	110.4	0.001052	1.1963	462.8	2229.0	2691.8	1.423	5.812	7.235	1.45
1.50	111.4	0.001053	1.1590	467.1	2226.3	2693.4	1.433	5.790	7.223	1.50
1.60	113.3	0.001055	1.0911	475.4	2220.8	2696.2	1.455	5.747	7.202	1.60
1.70	115.2	0.001056	1.0309	483.2	2215.8	2699.0	1.475	5.706	7.181	1.70
1.80	116.9	0.001058	0.97718	490.7	2210.8	2701.5	1.494	5.668	7.162	1.80
1.90	118.6	0.001059	0.92895	497.9	2206.1	2704.0	1.513	5.631	7.144	1.90
2.00	120.2	0.001061	0.88540	504.7	2201.6	2706.3	1.530	5.597	7.127	2.00

Properties of Pure Substances
(p)	(s _g)	$(\mathbf{s}_{\mathrm{fg}})$	(s _f)	(h _g)	(h_{fg})	(h _f)	(v _g)	(v _f)	(t)	(p)
2.1	7.111	5.564	1.547	2708.5	2197.2	511.3	0.84586	0.001062	121.8	2.1
2.2	7.095	5.532	1.563	2710.6	2193.0	517.6	0.80980	0.001064	123.3	2.2
2.3	7.080	5.502	1.578	2712.6	2188.9	523.7	0.77677	0.001065	124.7	2.3
2.4	7.066	5.473	1.593	2714.5	2184.9	529.6	0.74641	0.001066	126.1	2.4
2.5	7.052	5.445	1.607	2716.4	2181.1	535.3	0.71840	0.001068	127.4	2.5
2.6	7.039	5.418	1.621	2718.2	2177.3	540.9	0.69247	0.001069	128.7	2.6
2.7	7.026	5.392	1.634	2719.9	2173.7	546.2	0.66840	0.001070	130.0	2.7
2.8	7.014	5.367	1.647	2721.5	2170.1	551.4	0.64600	0.001071	131.2	2.8
2.9	7.002	5.342	1.660	2723.1	2166.6	556.5	0.62509	0.001072	132.4	2.9
3.0	6.991	5.319	1.672	2724.7	2163.2	561.5	0.60553	0.001074	133.5	3.0
3.1	6.980	5.297	1.683	2726.1	2159.9	566.2	0.58718	0.001075	134.7	3.1
3.2	6.969	5.274	1.695	2727.6	2156.7	570.9	0.56995	0.001076	135.8	3.2
3.3	6.959	5.253	1.706	2729.0	2153.5	575.5	0.55373	0.001077	136.8	3.3
3.4	6.949	5.232	1.717	2730.3	2 150.4	579.9	0.53843	0.001078	137.9	3.4
3.5	6.939	5.212	1.727	2731.6	2147.3	584.3	0.52397	0.001079	138.9	3.5
3.6	6.930	5.192	1.738	2732.9	2 144.4	588.5	0.51029	0.001080	139.9	3.6
3.7	6.921	5.173	1.748	2734.1	2141.4	592.7	0.49733	0.001081	140.8	3.7
3.8	6.912	5.154	1.758	2735.3	2138.6	596.7	0.48502	0.001082	141.8	3.8
3.9	6.903	5.136	1.767	2736.5	2135.7	600.8	0.47333	0.001083	142.7	3.9
4.0	6.894	5.118	1.776	2737.6	2 132.9	604.7	0.46220	0.001084	143.6	4.0
4.1	6.886	5.100	1.786	2738.7	2 130.2	608.5	0.45159	0.001085	144.5	4.1
4.2	6.878	5.083	1.795	2739.8	2127.5	612.3	0.44147	0.001086	145.4	4.2
4.3	6.870	5.067	1.803	2740.9	2124.9	616.0	0.43181	0.001087	146.3	4.3
4.4	6.862	5.050	1.812	2741.9	2122.3	619.6	0.42257	0.001088	147.1	4.4
4.5	6.855	5.035	1.820	2742.9	2119.7	623.2	0.41373	0.001089	147.9	4.5
4.6	6.847	5.018	1.829	2743.9	2117.2	626.7	0.40526	0.001090	148.7	4.6
4.7	6.840	5.003	1.837	2744.8	2114.7	630.1	0.39714	0.001090	149.5	4.7
4.8	6.833	4.988	1.845	2745.7	2112.2	633.5	0.38934	0.001091	150.3	4.8
4.9	6.826	4.973	1.853	2746.6	2109.8	636.8	0.38186	0.001092	151.1	4.9
5.0	6.819	4.959	1.860	2747.5	2107.4	640.1	0.37466	0.001093	151.8	5.0

(p)	(s _g)	(s _{fg})	(s _f)	(h _g)	$(\mathbf{h}_{\mathrm{fg}})$	(h _f)	(v _g)	(v _f)	(t)	(p)
5.2	6.806	4.931	1.875	2749.2	2102.7	646.5	0.36106	0.001095	153.3	5.2
5.4	6.793	4.903	1.890	2750.9	2098.1	652.8	0.34844	0.001096	154.8	5.4
5.6	6.781	4.877	1.904	2752.5	2093.7	658.8	0.33669	0.001098	156.2	5.6
5.8	6.769	4.851	1.918	2754.0	2089.3	664.7	0.32572	0.001099	157.5	5.8
6.0	6.758	4.827	1.931	2755.5	2085.1	670.4	0.31546	0.001101	158.8	6.0
6.2	6.747	4.803	1.944	2756.9	2080.8	676.1	0.30584	0.001102	160.1	6.2
6.4	6.736	4.780	1.956	2758.2	2076.7	681.5	0.29680	0.001104	161.4	6.4
6.6	6.725	4.757	1.968	2759.5	2072.7	686.8	0.28829	0.001105	162.6	6.6
6.8	6.715	4.735	1.980	2760.8	2068.8	692.0	0.28026	0.001107	163.8	6.8
7.0	6.705	4.713	1.992	2762.0	2064.9	697.1	0.27268	0.001108	165.0	7.0
7.2	6.696	4.693	2.003	2763.2	2061.2	702.0	0.26550	0.001110	166.1	7.2
7.4	6.686	4.672	2.014	2764.3	2057.4	706.9	0.25870	0.001111	167.2	7.4
7.6	6.677	4.652	2.025	2765.4	1053.7	711.7	0.25224	0.001112	168.3	7.6
7.8	6.668	4.633	2.035	2766.4	2050.1	716.3	0.24610	0.001114	169.4	7.8
8.0	6.660	4.614	2.046	2767.4	2046.5	720.9	0.24026	0.001115	170.4	8.0
8.2	6.651	4.595	2.056	2768.4	2043.0	725.4	0.23469	0.001116	171.4	8.2
8.4	6.643	4.577	2.066	2769.4	2039.6	729.9	0.22938	0.001118	172.4	8.4
8.6	6.635	4.560	2.075	2770.4	2036.2	734.2	0.22431	0.001119	173.4	8.6
8.8	6.627	4.542	2.085	2771.3	2032.8	738.5	0.21946	0.001120	174.4	8.8
9.0	6.619	4.525	2.094	2772.1	2029.5	742.6	0.21482	0.001121	175.4	9.0
9.2	6.612	4.509	2.103	2773.0	2026.2	746.8	0.21037	0.001123	176.3	9.2
9.4	6.604	4.492	2.112	2773.8	2023.0	750.8	0.20610	0.001124	177.2	9.4
9.6	6.597	4.476	2.121	2774.6	2019.8	754.8	0.20201	0.001125	178.1	9.6
9.8	6.590	4.460	2.130	2775.4	2016.7	758.7	0.19808	0.001126	179.0	9.8
10.0	6.583	4.445	2.138	2776.2	2013.6	762.6	0.19430	0.001127	179.9	0.0
10.5	6.566	4.407	2.159	2778.0	2006.0	772.0	0.18548	0.001130	182.0	0.5
11.0	6.550	4.371	2.179	2779.7	1998.6	781.1	0.17739	0.001133	184.1	1.0
11.5	6.534	4.336	2.198	2781.3	1991.4	789.9	0.17002	0.001136	186.0	1.5
12.0	6.519	4.303	2.216	2782.7	1984.3	798.4	0.16321	0.001139	188.0	2.0
12.5	6.505	4.271	2.234	2784.2	1977.5	806.7	0.15696	0.001141	189.8	2.5

	(p)	(s _g)	(s _{fg})	(s _f)	(h _g)	(h _{fg})	(h _f)	(v _g)	(v _f)	(t)	(p)
)	13.0	6.491	4.240	2.251	2785.4	1970.7	814.7	0.15114	0.001144	191.6	13.0
5	13.5	6.478	4.211	2.267	2786.7	1964.2	822.5	0.14576	0.001146	193.3	13.5
)	14.0	6.465	4.181	2.284	2787.8	1957.7	830.1	0.14073	0.001149	195.0	14.0
5	14.5	6.453	4.154	2.299	2788.9	1951.4	837.5	0.13606	0.001151	196.7	14.5
)	15.0	6.441	4.127	2.314	2789.9	1945.3	844.6	0.13167	0.001154	198.3	15.0
5	15.5	6.429	4.100	2.329	2790.8	1939.2	851.6	0.12756	0.001156	199.8	15.5
)	16.0	6.418	4.074	2.344	2791.7	1933.2	858.5	0.12370	0.001159	201.4	16.0
5	16.5	6.407	4.049	2.358	2792.6	1927.3	865.3	0.12006	0.001161	202.9	16.5
)	17.0	6.396	4.025	2.371	2793.4	1921.6	871.8	0.11664	0.001163	204.3	17.0
5	17.5	6.385	4.001	2.384	2794.1	1915.9	878.2	0.11340	0.001166	205.7	17.5
)	18.0	6.375	3.977	2.398	2794.8	1910.3	884.5	0.11033	0.001168	207.1	18.0
5	18.5	6.365	3.955	2.410	2795.5	1904.8	890.7	0.10742	0.001170	208.5	18.5
)	19.0	6.356	3.933	2.423	2796.1	1899.3	896.8	0.10467	0.001172	209.8	19.0
5	19.5	6.346	3.911	2.435	2796.7	1894.0	902.7	0.10204	0.001174	211.1	19.5
)	20.0	6.337	3.890	2.447	2797.2	1888.7	908.5	0.09955	0.001177	212.4	20.0
)	21.0	6.319	3.849	2.470	2798.2	1878.3	919.9	0.094902	0.001181	214.8	21.0
)	22.0	6.301	3.809	2.492	2799.1	1868.1	930.9	0.090663	0.001185	217.2	22.0
)	23.0	6.285	3.771	2.514	2799.8	1858.2	941.6	0.086780	0.001189	219.6	23.0
)	24.0	6.269	3.735	2.534	2800.4	1848.5	951.9	0.083209	0.001193	221.8	24.0
)	25.0	6.253	3.699	2.554	2801.0	1839.1	961.9	0.079915	0.001197	223.9	25.0
)	26.0	6.239	3.665	2.574	2801.4	1829.7	971.7	0.076865	0.001201	226.0	26.0
)	27.0	6.224	3.632	2.592	2801.7	1820.5	981.2	0.074033	0.001205	228.1	27.0
)	28.0	6.211	3.600	2.611	2802.0	1811.5	990.5	0.071396	0.001209	230.0	28.0
)	29.0	6.197	3.569	2.628	2802.2	1802.7	999.5	0.068935	0.001213	232.0	29.0
)	30.0	6.184	3.538	2.646	2802.3	1794.0	1008.3	0.066632	0.001216	233.8	30.0
)	31.0	6.171	3.509	2.662	2802.3	1785.4	1017.1	0.064473	0.001220	235.7	31.0
)	32.0	6.159	3.480	2.679	2802.3	1776.9	1025.4	0.062443	0.001224	237.4	32.0
)	33.0	6.146	3.452	2.694	2802.3	1768.6	1033.7	0.060533	0.001227	239.2	33.0
)	34.0	6.134	3.424	2.710	2802.1	1760.3	1041.8	0.058731	0.001231	240.9	34.0
)	35.0	6.123	3.398	2.725	2802.0	1752.3	1049.7	0.057028	0.001235	242.5	35.0

							1			
(p)	(t)	(v _f)	(v _g)	(h _f)	(\mathbf{h}_{fg})	(h _g)	(s _f)	(s _{fg})	(s _g)	(p)
36.0	244.2	0.001238	0.055417	1057.5	1744.2	2801.7	2.740	3.371	6.111	36.0
37.0	245.8	0.001242	0.053889	1065.2	1736.2	2801.4	2.755	3.345	6.100	37.0
38.0	247.3	0.001245	0.052439	1072.7	1728.4	2801.1	2.769	3.321	6.090	38.0
39.0	248.8	0.001249	0.051061	1080.1	1720.7	2800.8	2.783	3.296	6.079	39.0
40.0	250.3	0.001252	0.049749	1087.4	1712.9	2800.3	2.797	3.272	6.069	40.0
42.0	253.2	0.001259	0.047306	1101.6	1697.8	2799.4	2.823	3.225	6.048	42.0
44.0	256.1	0.001266	0.045078	1115.4	1682.9	2798.3	2.849	3.180	6.029	44.0
46.0	258.8	0.001273	0.043036	1128.8	1668.2	2797.0	2.874	3.136	6.010	46.0
48.0	261.4	0.001279	0.041158	1141.8	1653.9	2795.7	2.897	3.094	5.991	48.0
50.0	263.9	0.001286	0.039425	1154.5	1639.7	2794.2	2.921	3.053	5.974	50.0
52.0	266.4	0.001293	0.037820	1166.9	1625.7	2792.6	2.943	3.013	5.956	52.0
54.0	268.8	0.001299	0.036330	1179.0	1611.8	2790.8	2.965	2.974	5.939	54.0
56.0	271.1	0.001306	0.034942	1190.8	1598.2	2789.0	2.986	2.937	5.923	56.0
58.0	273.4	0.001312	0.033646	1202.4	1584.6	2787.0	3.007	2.899	5.906	58.0
60.0	275.6	0.001319	0.032433	1213.7	1571.3	2785.0	3.027	2.863	5.890	60.0
62.0	277.7	0.001325	0.031295	1224.9	1558.0	2782.9	3.047	2.828	5.875	62.0
64.0	279.8	0.001332	0.030225	1235.8	1544.8	2780.6	3.066	2.794	5.860	64.0
66.0	281.9	0.001338	0.029218	1246.5	1531.8	2778.3	3.085	2.760	5.845	66.0
68.0	283.9	0.001345	0.028267	1257.1	1518.8	2775.9	3.104	2.727	5.831	68.0
70.0	285.8	0.001 351	0.027368	1267.4	1506.0	2773.4	3.122	2.694	5.816	70.0
72.0	287.7	0.001358	0.026517	1277.7	1493.2	2770.9	3.140	2.662	5.802	72.0
74.0	289.6	0.001365	0.025711	1287.8	1480.4	2768.2	3.157	2.631	5.788	74.0
76.0	291.4	0.001371	0.024944	1297.7	1467.8	2765.5	3.174	2.600	5.774	76.0
78.0	293.2	0.001378	0.024215	1307.5	1455.4	2762.7	3.191	2.569	5.760	78.0
80.0	295.0	0.001384	0.023521	1317.2	1442.7	2759.9	3.208	2.539	5.747	80.0
82.0	296.7	0.001391	0.022860	1326.7	1430.3	2757.0	3.224	2.510	5.734	82.0
84.0	298.4	0.001398	0.022228	1336.2	1417.8	2754.0	3.240	2.481	5.721	84.0
86.0	300.1	0.001404	0.021624	1345.4	1405.5	2750.9	3.256	2.452	5.708	86.0
88.0	301.7	0.001411	0.021046	1354.7	1393.1	2747.8	3.271	2.424	5.695	88.0
90.0	303.3	0.001418	0.020493	1363.8	1380.8	2744.6	3.287	2.395	5.682	90.0

Properties of Pure Substances

(p)	(t)	(V _f)	(v _g)	(h _f)	(h _{fg})	(h _g)	(s _f)	(s _{fg})	(s _g)	(p)
92	304.9	0.001425	0.019962	1372.8	1368.5	2741.3	3.302	2.367	5.669	92
94	305.5	0.001432	0.019453	1381.7	1356.3	2738.0	3.317	2.340	5.657	94
96	308.0	0.001439	0.018964	1390.6	1344.1	2734.7	3.332	2.313	5.644	96
98	309.5	0.001446	0.018493	1399.4	1331.9	2731.2	3.346	2.286	5.632	98
100	301.0	0.001453	0.018041	1408.0	1319.7	2727.7	3.361	2.259	5.620	100
105	314.6	0.001470	0.016981	1429.5	1289.2	2718.7	3.396	2.194	5.590	105
110	318.0	0.001489	0.016007	1450.5	1258.8	2709.3	3.430	2.129	5.560	110
115	321.4	0.001508	0.015114	1471.3	1228.2	2699.5	3.464	2.066	5.530	115
120	324.6	0.001527	0.014285	1491.7	1197.5	2698.2	3.497	2.003	5.500	120
125	327.8	0.001547	0.013518	1511.9	1166.5	2678.4	3.530	1.941	5.471	125
130	330.8	0.001567	0.012800	1531.9	1135.1	2667.0	3.561	1.880	5.441	130
135	333.8	0.001588	0.012130	1551.8	1103.3	2655.1	3.593	1.818	5.411	135
140	336.6	0.001611	0.011498	1571.5	1070.9	2642.4	3.624	1.756	5.380	140
145	339.4	0.001634	0.010905	1591.3	1037.9	2629.2	3.655	1.694	5.349	145
150	342.1	0.001658	0.010343	1610.9	1004.2	2615.1	3.686	1.632	5.318	150
155	344.8	1.001683	0.009813	1630.7	969.7	2600.4	3.716	1.570	5.286	155
160	347.3	0.001710	0.009310	1650.4	934.5	2584.9	3.747	1.506	5.253	160
165	349.7	0.001739	0.008833	1670.4	898.5	2568.9	3.778	1.442	5.220	165
170	352.3	0.001770	0.008372	1691.6	860.0	2551.6	3.811	1.375	5.186	170
175	354.6	0.001803	0.007927	1713.3	820.0	2533.3	3.844	1.306	5.150	175
180	357.0	0.001840	0.007497	1734.8	779.1	2513.9	3.877	1.236	5.113	180
185	359.2	0.001881	0.007082	1756.5	736.5	2493.0	3.910	1.164	5.074	185
190	361.4	0.001926	0.006676	1778.7	691.8	2470.5	3.943	1.090	5.033	190
195	363.6	0.001978	0.006276	1801.9	643.9	2445.8	3.978	1.011	4.989	195
200	365.7	0.002037	0.005875	1826.6	591.6	2418.2	4.015	0.926	4.941	200
205	367.8	0.002110	0.005462	1854.2	532.0	2386.2	4.056	0.830	4.886	205
210	369.8	0.002202	0.005023	1886.3	461.2	2347.5	4.105	0.717	4.822	210
215	371.8	0.002342	0.004509	1929.4	365.2	2294.6	4.170	0.566	4.736	215
220	373.7	0.002668	0.003735	2010.3	186.3	2196.6	4.293	0.288	4.581	220
221.2	374.15	0.003170	0.003170	2107.4	000.0	2107.4	4.443	0.000	4.443	221.2

Mechanical Science-II

			Tabl	le: 7.5. S	pecific V	olume of	Superhe	ated Stea	am			
Absolute Pressure in bar	Saturation Temperature in °C				Specif	fic volume ((v) in m ³ /kg	g at various	s temperatu	res in °C		
(p)	(t _s)	100	150	200	250	300	350	400	500	600	700	800
0.02	17.5	86.08	97.63	109.2	120.7	132.2	143.8	155.3	178.4	201.5	224.6	247.6
0.04	29.0	43.03	48.81	54.58	60.35	66.12	71.89	77.66	89.20	100.7	112.3	123.8
0.06	36.2	28.68	32.53	36.38	40.23	44.08	47.93	51.77	59.47	67.16	74.85	82.54
0.08	41.5	21.50	24.40	27.28	30.17	33.06	35.94	38.83	44.60	50.37	56.14	61.91
0.10	45.8	17.20	19.51	21.83	24.14	26.45	28.75	31.06	35.68	40.30	44.91	49.53
0.15	54.0	11.51	13.06	14.61	16.16	17.71	19.25	20.80	23.89	26.98	30.07	33.16
0.20	60.1	8.585	9.748	10.91	12.07	13.22	14.37	15.53	17.84	20.15	22.45	24.76
0.25	65.0	6.874	7.808	8.737	9.665	10.59	11.52	12.44	14.29	16.14	17.99	19.84
0.30	69.1	5.714	6.493	7.268	8.040	8.811	9.581	10.35	11.89	13.43	14.70	16.51
0.35	72.7	4.898	5.568	6.233	6.896	7.557	8.218	8.879	10.20	11.52	12.84	14.16
0.40	75.9	4.279	4.866	5.448	6.028	6.607	7.185	7.763	8.918	10.07	11.23	12.38
0.45	78.7	3.803	4.325	4.844	5.360	5.875	6.389	6.903	7.930	8.957	9.984	10.99
0.50	81.3	3.418	3.889	4.356	4.821	5.284	5.747	6.209	7.134	8.057	8.981	9.904
0.60	86.0	2.844	3.238	3.628	4.016	4.402	4.788	5.174	5.944	6.714	7.484	8.254
0.70	90.0	2.434	2.773	3.108	3.441	3.772	4.103	4.434	5.095	5.755	6.415	7.074
0.80	93.5	2.126	2.425	2.718	3.010	3.300	3.590	3.879	4.457	5.035	5.613	6.190
0.90	96.7	1.887	2.153	2.415	2.674	2.933	3.190	3.448	3.962	4.475	4.989	5.502
1.00	99.6	1.696	1.936	2.172	2.406	2.639	2.871	3.103	3.565	4.028	4.490	4.952
1.50	111.4		1.285	1.444	1.601	1.757	1.912	2.067	2.376	2.685	2.993	3.301
2.00	120.2		0.9595	1.080	1.199	1.316	1.433	1.549	1.781	2.013	2.244	2.475

eated	Steam	
σ at va	rious temperatures in	0

(p)	(t _s)	100	150	200	250	300	350	400	500	600	700	800
2.5	127.4		0.7641	0.8620	0.9574	1.052	1.145	1.239	1.424	1.610	1.795	1.980
3.0	133.5		0.6337	0.7164	0.7964	0.8753	0.9535	1.031	1.187	1.341	1.496	1.650
3.5	138.9		0.5406	0.6123	0.6814	0.7493	0.8166	0.8835	1.017	1.149	1.282	1.414
4.0	143.6		0.4707	0.5343	0.5952	0.6549	0.713 9	0.7725	0.8892	1.005	1.121	1.237
4.5	147.9		0.4165	0.4738	0.5284	0.5817	0.6343	0.6865	0.7905	0.8939	0.9971	1.100
5.0	151.8			0.4250	0.4744	0.5226	0.5701	0.6172	0.7108	0.8040	0.8969	0.9896
6.0	158.8			0.3520	0.3939	0.4344	0.4742	0.5136	0.5918	0.6696	0.7471	0.8245
7.0	165.0			0.2999	0.3364	0.3714	0.4057	0.4396	0.5069	0.5737	0.6402	0.7066
8.0	170.4			0.2608	0.2932	0.3241	0.3543	0.3842	0.4432	0.5017	0.5600	0.6181
9.0	175.4			0.2303	0.2596	0.2874	0.3144	0.3410	0.3936	0.4458	0.4976	0.5493
10.0	179.9			0.2059	0.2328	0.2580	0.2824	0.3065	0.3540	0.4010	0.4477	0.4943
11.0	184.1			0.1859	0.2108	0.2339	0.2563	0.2782	0.3215	0.3644	0.4069	0.4492
12.0	188.0			0.1692	0.1924	0.2139	0.234 5	0.2547	0.2945	0.3338	0.3729	0.4118
13.0	191.6			0.1551	0.1769	0.1969	0.2161	0.2348	0.2716	0.3080	0.3441	0.3800
14.0	195.0			0.1429	0.1636	0.1823	0.2002	0.2177	0.2520	0.2859	0.3194	0.3528
15.0	198.3			0.1324	0.1520	0.1697	0.1865	0.2029	0.2350	0.2667	0.2980	0.3292
16.0	201.4				0.1419	0.1587	0.1745	0.1900	0.2202	0.2499	0.2793	0.3086
17.0	204.3				0.1329	0.1489	0.1640	0.1786	0.2070	0.2351	0.2628	0.2904
18.0	207.1				0.1250	0.1402	0.1546	0.1684	0.1954	0.2219	0.2481	0.2742
19.0	209.8				0.1179	0.1325	0.1461	0.1593	0.184 9	0.2101	0.2350	0.2597
20.0	212.4				0.1115	0.1255	0.1386	0.1511	0.1756	0.1995	0.2232	0.2467
22.0	217.2				0.1004	0.1134	0.1255	0.1370	0.1593	0.1812	0.2028	0.2242
24.0	221.8				0.09108	0.1034	0.1146	0.1252	0.1458	0.1659	0.1858	0.2054
26.0	226.0				0.08321	0.09483	0.1053	0.1153	0.1344	0.1530	0.1714	0.1895
28.0	230.0				0.07644	0.08751	0.09740	0.1067	0.124 6	0.1419	0.1590	0.1759
30.0	233.8				0.07055	0.08116	0.09053	0.09931	0.1161	0.1323	0.1483	0.1641
32.0	237.4				0.06538	0.07559	0.08451	0.09283	0.1087	0.1239	0.1390	0.1538
34.0	240.9				0.06080	0.07068	0.07920	0.08711	0.1021	0.1165	0.1307	0.1447
36.0	244.2				0.05670	0.06630	0.07448	0.08202	0.09626	0.1100	0.1234	0.1366
38.0	247.3				0.05302	0.06237	0.07025	0.07747	0.09104	0.1041	0.1168	0.1294

Properties of Pure Substances

(p)	(t _s)	100	150	200	250	300	350	400	500	600	700	800
40.0	250.3					0.05883	0.06645	0.07338	0.08634	0.09876	0.1109	0.1229
42.0	253.2					0.05563	0.06300	0.06967	0.08209	0.09397	0.1056	0.1170
44.0	256.0					0.05270	0.05986	0.06630	0.07823	0.08961	0.1007	0.1116
46.0	258.8					0.05003	0.05699	0.06322	0.07470	0.08562	0.09626	0.1067
48.0	261.4					0.04757	0.05436	0.06039	0.07147	0.08197	0.09219	0.1022
50.0	263.9					0.04530	0.05194	0.05779	0.06849	0.07862	0.08845	0.09809
55.0	269.9					0.04343	0.04666	0.05213	0.06202	0.07131	0.08031	0.08912
60.0	275.6					0.03615	0.04222	0.04738	0.05659	0.06518	0.07348	0.08159
65.0	280.8					0.03258	0.00848	0.04338	0.05203	0.06003	0.06774	0.07526
70.0	285.8					0.02946	0.00523	0.03992	0.04809	0.05559	0.06279	0.06980
75.0	290.5					0.02672	0.03243	0.03694	0.04469	0.05176	0.05852	0.06509
80.0	295.0					0.02426	0.02995	0.03431	0.04170	0.04839	0.05477	0.06096
85.0	299.2					0.02191	0.02776	0.03200	0.03908	0.04544	0.05148	0.05732
90.0	303.3						0.02579	0.02993	0.03674	0.04280	0.04853	0.05408
95.0	307.2						0.02403	0.02808	0.03465	0.04045	0.04591	0.05119
100.0	311.0						0.02242	0.02641	0.03276	0.03832	0.04355	0.04858
110.0	318.0						0.01961	0.02351	0.02950	0.03466	0.03947	0.04408
120.0	324.6						0.01721	0.02108	0.02679	0.03160	0.03607	0.04033
130.0	330.8						0.01510	0.01902	0.02449	0.02902	0.03319	0.03716
140.0	336.6						0.01321	0.01723	0.02251	0.02680	0.03072	0.03444
150.0	342.1						0.01146	0.01566	0.02080	0.02488	0.02859	0.03209
160.0	347.3						0.00976	0.01428	0.01929	0.02320	0.02672	0.03003
170.0	352.3							0.01303	0.01797	0.02172	0.02507	0.02821
180.0	357.0							0.01191	0.01679	0.02040	0.02300	0.02659
190.0	361.4							0.01089	0.01573	0.01922	0.02229	0.02515
200.0	365.7							0.00995	0.01477	0.01816	0.02111	0.02385
210.0	369.8							0.00907	0.01391	0.01720	0.02004	0.02267
220.0	373.7							0.00825	0.01312	0.01633	0.01907	0.02160
221.2	374.15							0.00816	0.01303	0.01622	0.01895	0.02135

Absolute	Saturation			Spec	ific enthalp	ohy (h) in k	J/kg at var	ious tempe	ratures in	°C		
Pressure in bar (p)	Temperature in °C (t _s)	100	150	200	250	300	350	400	500	600	700	800
0.02	17.5	2688.5	2783.7	2880.0	2977.7	3076.8	3177.5	3279.7	3489.2	3705.6	3928.8	4158.7
0.04	29.0	2688.3	2783.5	2879.9	2977.6	3076.8	3177.4	3279.7	3489.2	3705.6	3928.8	4158.7
0.06	36.2	2688.0	2783.4	2879.8	2977.6	3076.7	3177.4	3279.6	3489.2	3705.6	3928.8	4158.7
0.08	41.5	2687.8	2783.2	2879.7	2977.5	3076.7	3177.3	3279.6	3489.1	3705.5	3928.8	4158.7
0.10	45.8	2687.5	2783.1	2879.6	2977.4	3076.6	3177.3	3279.6	3489.1	3705.5	3928.8	4158.7
0.15	54.0	2686.9	2782.4	2879.5	2977.3	3076.5	3177.7	3279.5	3489.1	3705.5	3928.7	4158.7
0.20	60.1	2686.3	2782.3	2879.2	2977.1	3076.4	3177.1	3279.4	3489.0	3705.4	3928.7	4158.7
0.25	65.0	2685.7	2782.0	2879.0	2977.0	3076.3	3177.0	3279.3	3489.0	3705.4	3928.7	4158.6
0.30	69.1	2685.1	2781.6	2878.7	2976.8	3076.1	3176.9	3279.3	3488.9	3705.4	3928.7	4158.6
0.35	72.7	2684.5	2781.2	2878.5	2976.7	3076.0	3176.8	3279.2	3488.9	3705.3	3928.7	4158.6
0.40	75.9	2683.8	2780.9	2878.2	2976.5	3075.9	3176.8	3279.1	3488.8	3705.3	3928.6	4158.6
0.45	78.7	2683.2	2780.5	2878.0	2976.3	3075.8	3176.7	3279.1	3488.8	3705.2	3928.6	4158.5
0.50	81.3	2682.6	2780.1	2877.7	2976.1	3075.7	3176.6	3279.0	3488.7	3705.2	3928.6	4158.5
0.60	86.0	2681.3	2779.4	2877.3	2975.8	3075.4	3176.4	3278.8	3488.6	3705.1	3928.5	4158.5
0.70	90.0	2680.0	2778.6	2876.8	2975.5	3075.2	3176.2	3278.7	3488.5	3705.0	3928.4	4158.4
0.80	93.5	2678.8	2777.8	2876.3	2975.2	3075.0	3176.0	3278.5	3488.4	3705.0	3928.4	4158.4
0.90	96.7	2677.5	2777.1	2875.8	2974.8	3074.7	3175.8	3278.4	3488.3	3704.9	3928.3	4158.3
1.00	99.6	2676.2	2176.3	2875.4	2974.5	3074.5	3175.6	3278.2	3488.1	3704.8	3928.2	4158.3
1.50	111.4		2772.5	2872.9	2972.9	3073.3	3174.7	3277.5	3487.6	3704.4	3927.9	4158.0
2.00	120.2		2768.5	2870.5	2971.2	3072.1	3173.8	3276.7	3487.0	3704.0	3927.6	4157.8

Table: 7	7.6.	Specific	Enthalpy	of Su	perheated	Steam
			• • · · ·			

	1	1										
(p)	(t _s)	100	150	200	250	300	350	400	500	600	700	800
2.5	127.4		2764.5	2868.0	2969.6	3070.9	3172.8	3275.9	3486.5	3703.6	3927.3	4157.6
3.0	133.5		2760.4	2865.5	2967.9	3069.7	3 171.9	3275.2	3486.0	3703.2	3927.0	4157.3
3.5	138.9		2756.3	2863.0	2966.2	3068.4	3170.9	3274.4	3485.4	3702.7	3926.7	4157.1
4.0	143.6		2752.0	2860.4	2964.5	3067.2	3170.0	3273.6	3484.9	3702.3	3926.4	4156.9
4.5	147.9		2746.7	2857.8	2962.8	3066.0	3 169.1	3272.9	3484.3	3701.9	3926.1	4156.7
5.0	151.8			2855.1	2961.1	3064.8	3168.1	3272.1	3483.8	3701.5	3925.8	4156.4
6.0	158.8			2849.7	2957.6	3062.3	3166.2	3270.6	3482.7	3700.7	3925.1	4155.9
7.0	165.0			2844.2	2954.0	3059.8	3164.3	3269.0	3481.6	3699.9	3924.5	4155.5
8.0	170.4			2838.6	2950.4	3057.3	3162.4	3267.5	3480.5	3699.1	3923.9	4155.0
9.0	175.4			2832.7	2946.8	3054.7	3160.5	3266.0	3479.4	3698.2	3923.3	4154.5
10.0	179.9			2826.8	2943.0	3052.1	3158.5	3264.4	3478.3	3697.4	3922.7	4154.1
11.0	184.1			2820.7	2939.3	3049.6	3156.6	3262.9	3477.2	3696.6	3922.0	4153.6
12.0	188.0			2814.4	2935.4	3.046.9	3154.6	3261.3	3476.1	3695.8	3921.4	4153.1
13.0	191.6			2808.0	2931.5	3.044.3	3152.7	3259.7	3475.0	3695.0	3920.8	4152.7
14.0	195.0			2801.4	2927.6	3 041.6	3150.7	3258.2	3473.9	3694.1	3920.2	4152.2
15.0	198.3			2794.7	2923.5	3038.9	3148.7	3256.6	3472.8	3693.3	3919.6	4 151.7
16.0	201.4				2919.4	3036.2	3146.7	3255.0	3471.7	3692.5	3918.9	4151.3
17.0	204.3				2915.3	3033.5	3144.7	3253.5	3470.6	3691.7	3918.3	4150.8
18.0	207.1				2911.0	3030.7	3142.7	3251.9	3469.5	3690.9	3917.7	4150.3
19.0	209.8				2906.7	3027.9	3140.7	3250.3	3468.4	3690.0	3917.1	4149.8
20.0	212.4				2902.4	3025.0	3138.6	3248.7	3467.3	3689.2	3916.5	4149.4
22.0	217.2				2893.4	3019.3	3134.5	3245.5	3465.1	3687.6	3915.2	4148.4
24.0	221.8				2884.2	3013.4	3130.4	3242.3	3462.9	3685.9	3914.0	4147.5
26.0	226.0				2874.7	3007.4	3126.1	3239.0	3460.6	3684.3	3912.7	4146.6
28.0	230.0				2864.9	3 001.3	3121.9	3235.8	3458.4	3682.6	3911.5	4145.6
30.0	233.8				2854.8	2995.1	3117.5	3232.5	3456.2	3681.0	3910.3	4144.7
32.0	237.4				2844.4	2988.7	3113.2	3229.2	3454.0	3679.3	3909.0	4143.8
34.0	240.9				2833.6	2982.2	3108.7	3225.9	3451.7	3677.7	3907.8	4142.8
36.0	244.2				2822.5	2975.6	3104.2	3222.5	3449.5	3676.1	3906.5	4141.9
38.0	247.3				2811.0	2968.9	3099.7	3219.1	3447.2	3674.4	3905.3	4 141.0

- 1	N
	<u> </u>
	Ċ
	\sim

(p)	(t _s)	100	150	200	250	300	350	400	500	600	700	800
40.0	250.3					2962.0	3095.1	3215.7	3445.0	3672.8	3904.1	4140.0
42.0	253.2					2955.0	3090.4	3212.3	3442.7	3671.1	3902.8	4139.1
44.0	256.0					2947.8	3085.7	3208.8	3440.5	3669.5	3901.6	4138.2
46.0	258.8					2940.5	3080.9	3205.3	3438.2	3667.8	3900.3	4 137.2
48.0	261.4					2933.1	3076.1	3201.8	3435.9	3666.2	3899.1	4136.3
50.0	263.9					2925.5	3071.2	3198.3	3433.7	3664.5	3897.9	4135.3
55.0	269.9					2905.8	3058.7	3189.3	3427.9	3660.4	3894.8	4 133.0
60.0	275.6					2835.0	3045.8	3180.1	3422.2	3656.2	3891.7	4 130.7
65.0	280.8					2863.0	3032.4	3170.8	3416.4	3652.1	3888.6	4128.8
70.0	285.8				••	2839.0	3018.7	3161.2	3410.6	3647.9	3885.4	4126.0
75.0	290.5					2814.1	3004.5	3151.6	3404.7	3643.7	3882.4	4123.7
80.0	295.0					2786.6	2989.9	3141.6	3398.8	3639.5	3879.2	4 121.3
85.0	299.2					2757.1	2974.7	3131.5	3392.8	3635.4	3876.1	4119.0
90.0	303.3						2959.0	3121.2	3386.8	3631.1	3873.0	4116.7
95.0	307.2						2942.7	3110.7	3380.7	3627.0	3869.9	4114.4
100.0	311.0						2925.8	3099.9	3374.6	3622.7	3866.8	4112.0
110.0	318.0						2889.6	3077.8	3362.2	3614.2	3860.5	4107.3
120.0	324.6						2849.7	3054.8	3349.6	3605.7	3854.3	4102.7
130.0	330.8						2805.0	3030.7	3336.8	3597.1	3848.0	4098.0
140.0	336.6						2754.2	3005.6	3323.8	3588.5	3841.7	4093.3
150.0	342.1						2694.8	2979.1	3310.6	3579.8	3835.4	4088.6
160.0	347.3						2620.8	2951.3	3297.1	3571.0	3829.1	4084.0
170.0	352.3							2921.7	3283.5	3562.2	3822.8	4079.3
180.0	357.0							2890.3	2269.6	3553.4	3816.5	4074.6
190.0	361.4							2856.7	3255.4	3544.5	3810.2	4070.0
200.0	365.7							2820.5	3241.1	3535.5	3803.8	4065.3
210.0	369.8							2781.3	3226.5	3526.5	3797.5	4060.6
220.0	373.7							2738.8	3211.7	3517.4	3791.1	4055.9
221.2	374.15							2734.5	3210.7	3516.4	3789.1	4054.7

Absolute	Saturation	Specific entropy (s) in kJ/kg K at various temperatures in °C											
PressureTemperaturein barin °C(p)(t,)	100	150	200	250	300	350	400	500	600	700	800		
0.02	17.5	9.193	9.433	9.648	9.844	10.025	10.193	10.351	10.641	10.904	11.146	11.371	
0.04	29.0	8.873	9.113	9.328	9.524	9.705	9.874	10.031	10.321	10.585	10.827	11.051	
0.06	36.2	8.685	8.925	9.141	9.337	9.518	9.686	9.844	10.134	10.397	10.639	10.864	
0.08	41.5	8.552	8.792	9.008	9.204	9.385	9.554	9.711	10.001	10.265	10.507	10.731	
0.10	45.8	8.449	8.689	8.905	9.101	9.282	9.450	9.608	9.898	10.162	10.404	10.628	
0.15	54.0	8.261	8.502	8.718	8.915	9.096	9.264	9.422	9.712	9.975	10.217	10.442	
0.20	60.1	8.126	8.368	8.584	8.781	8.962	9.130	9.288	9.578	9.842	10.084	10.309	
0.25	65.0	8.022	8.264	8.481	8.678	8.859	9.028	9.186	9.476	9.739	9.981	10.206	
0.30	69.1	7.936	8.179	8.396	8.593	8.774	8.943	9.101	9.391	9.654	9.897	10.121	
0.35	72.7	7.864	8.107	8.325	8.522	8.703	8.872	9.030	9.320	9.583	9.826	10.050	
0.40	75.9	7.801	8.045	8.263	8.460	8.641	8.810	8.968	9.258	9.522	9.764	9.989	
0.45	78.7	7.745	7.990	8.208	8.405	8.587	8.755	8.914	9.204	9.467	9.709	9.934	
0.50	81.3	7.695	7.941	8.159	8.356	8.538	8.707	8.865	9.155	9.419	9.661	9.886	
0.60	86.0	7.609	7.855	8.074	8.272	8.454	8.622	8.781	9.071	9.334	9.576	9.801	
0.70	90.0	7.535	7.783	8.002	8.200	8.382	8.551	8.709	9.000	9.263	9.505	9.730	
0.80	93.5	7.470	7.720	7.940	8.138	8.320	8.489	8.648	8.938	9.214	9.444	9.669	
0.90	96.7	7.413	7.664	7.884	8.083	8.266	8.435	8.593	8.884	9.147	9.389	9.614	
1.00	99.6	7.362	7.614	7.835	8.034	8.217	8.386	8.544	8.835	9.098	9.341	9.565	
1.50	111.4		7.419	7.644	7.845	8.028	8.198	8.356	8.647	8.911	9.153	9.378	
2.00	120.2		7.279	7.507	7.710	7.894	8.064	8.223	8.514	8.778	9.020	9.245	

 Table: 7.7. Specific Entropy of Superheated Steam

$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	8.119 8.034 7.962	8.410	8.674		
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	8.034 7.962	8 3 2 6		8.917	9.142
3.5 138.5 6.998 7.237 7.444 7.631 7.802 4.0 143.6 6.929 7.171 7.380 7.568 7.740 4.5 147.9 6.866 7.112 7.323 7.512 7.684 5.0 151.8 7.059 7.272 7.461 7.634 6.0 158.8 6.966 7.183 7.374 7.548 7.0 165.0 6.886 7.107 7.300 7.475 8.0 170.4 6.815 7.040 7.235 7.411 9.0 175.4 6.692 6.926 7.177 7.354 10.0 179.9 6.638 6.876 7.078 7.257 12.0 188.0 6.537 6.831 7.034 7.214 13.0 191.6 6.494 6.748 6.956 7.139 14.0 195.0 6.494 6.748 6.956 7.139	7 962	0.520	8.590	8.833	9.058
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	1.902	8.254	8.518	8.761	8.986
4.5 147.9 6.866 7.112 7.323 7.512 7.684 5.0 151.8 7.059 7.272 7.461 7.634 6.0 158.8 6.966 7.183 7.374 7.548 7.0 165.0 6.886 7.107 7.300 7.475 8.0 170.4 6.815 7.040 7.235 7.411 9.0 175.4 6.692 6.926 7.125 7.303 11.0 184.1 6.638 6.876 7.078 7.257 12.0 188.0 6.539 6.788 6.994 7.175 14.0 191.6 6.644 6.748 6.956 7.139	7.899	8.192	8.456	8.699	8.925
5.0 151.8 7.059 7.272 7.461 7.634 6.0 158.8 6.966 7.183 7.374 7.548 7.0 165.0 6.886 7.107 7.300 7.475 8.0 170.4 6.815 7.040 7.235 7.411 9.0 175.4 6.751 6.980 7.177 7.354 10.0 179.9 6.692 6.926 7.125 7.303 11.0 184.1 6.638 6.876 7.078 7.257 12.0 188.0 6.539 6.788 6.994 7.175 14.0 191.6 6.494 6.748 6.956 7.139	7.844	8.137	8.402	8.645	8.870
	7.795	8.088	8.353	8.596	8.821
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	7.709	8.003	8.268	8.511	8.737
8.0 170.4 6.815 7.040 7.235 7.411 9.0 175.4 6.751 6.980 7.177 7.354 10.0 179.9 6.692 6.926 7.125 7.303 11.0 184.1 6.638 6.876 7.078 7.257 12.0 188.0 6.587 6.831 7.034 7.214 13.0 191.6 6.494 6.748 6.956 7.139 14.0 195.0 6.494 6.748 6.956 7.139	7.636	7.931	8.196	8.440	8.665
9.0 175.4 6.751 6.980 7.177 7.354 10.0 179.9 6.692 6.926 7.125 7.303 11.0 184.1 6.638 6.876 7.078 7.257 12.0 188.0 6.587 6.831 7.034 7.214 13.0 191.6 6.494 6.748 6.994 7.175 14.0 195.0 6.494 6.748 6.956 7.139	7.573	7.868	8.134	8.377	8.603
10.0 179.9 6.692 6.926 7.125 7.303 11.0 184.1 6.638 6.876 7.078 7.257 12.0 188.0 6.587 6.831 7.034 7.214 13.0 191.6 6.539 6.788 6.994 7.175 14.0 195.0 6.494 6.748 6.956 7.139	7.517	7.812	8.079	8.323	8.549
11.0 184.1 6.638 6.876 7.078 7.257 12.0 188.0 6.587 6.831 7.034 7.214 13.0 191.6 6.539 6.788 6.994 7.175 14.0 195.0 6.494 6.748 6.956 7.139	7.467	7.763	8.029	8.273	8.500
12.0 188.0 6.587 6.831 7.034 7.214 13.0 191.6 6.539 6.788 6.994 7.175 14.0 195.0 6.494 6.748 6.956 7.139	7.421	7.718	7.985	8.229	8.455
13.0 191.6 6.539 6.788 6.994 7.175 14.0 195.0 6.494 6.748 6.956 7.139 14.0 195.0 6.494 6.748 6.956 7.139	7.379	7.677	7.944	8.188	8.415
14.0 195.0 6.494 6.748 6.956 7.139	7.340	7.639	7.906	8.151	8.378
	7.305	7.603	7.871	8.116	8.343
15.0 198.3 6.451 6.710 6.921 7.104	7.271	7.570	7.839	8.084	8.311
16.0 201.4 6.674 6.887 7.072	7.239	7.540	7.808	8.054	8.281
17.0 204.3 6.640 6.856 7.042	7.210	7.511	7.779	8.025	8.252
18.0 207.1 6.607 6.826 7.013	7.182	7.483	7.752	7.998	8.226
19.0 209.8 6.576 6.797 7.986	7.155	7.457	7.727	7.973	8.200
20.0 212.4 6.545 6.770 6.960	7.130	7.432	7.702	7.949	8.176
22.0 217.2 6.488 6.718 6.911	7.082	7.386	7.657	7.904	8.132
24.0 221.8 6.434 6.670 6.866	7.038	7.344	7.615	7.862	8.091
26.0 226.0 6.382 6.625 6.824	6.998	7.305	7.577	7.825	8.053
28.0 230.0 6.333 6.582 6.784	6.960	7.269	7.541	7.789	8.018
30.0 233.8 6.286 6.542 6.747	6.925	7.235	7.508	7.756	7.986
32.0 237.4 6.240 6.504 6.712	6.891	7.203	7.477	7.726	7.955
34.0 240.9 6.195 6.467 6.679	6.860	7.172	7.447	7.697	7.927
36.0 244.2 6.151 6.432 6.647	(000	7 1 4 4	7 420	7 669	7 900
38.0 247.3 6.109 6.397 6.616	6.829	/.144	7.420	1.000	1.200

7.8.1 Steam Tables and Their Uses

The properties of dry saturated steam like its temperature of formation *i.e.*, saturation temperature, sensible heat, latent heat of vaporisation, enthalpy or total heat, specific volume of water, specific volume of vapour, enthalpy of liquid and enthalpy of vapour, vary with pressure and can be found by experiments only. Then properties have been carefully determined and made available in a tabular form known as steam tables.

There are two important steam tables, one in terms of absolute pressure and other in terms of temperature. The properties of dry saturated steam is shown in the five tables (from 7.3 to 7.7).

7.9 ENTROPY OF STEAM

The entropy of steam consist of

- (a) Increase in entropy of water during heating from freezing point to boiling point, corresponding to the pressure at which the water is being heated.
- (b) Increase in entropy during evaporation.
- (c) Increase in entropy during super heating.

7.10 ENTROPY OF WATER

Consider 1 kg of water being heated at a constant pressure from freezing point (0°C or 273 K) to the boiling temperature. Now consider an instant when the absolute temperature of water is T K.

Let for small rises in temperature of dT, the heat absorbed by 1 kg of water is δq then

 $\delta q = \text{mass} \times \text{specific heat of water} \times \text{rise in temperature}$

$$= 1 \times C_{w} \times dT$$
$$= C_{w} \cdot dT$$

$$C_{w} = \text{specific heat of water} \\ = 4.2 \text{ kJ/kg K}$$

we know that increase in entropy for rise in temperature dT

$$ds = \frac{\delta q}{T} = C_w \cdot \frac{dT}{T}$$

total increase in entropy of water from freezing point to boiling point

$$\int_{0}^{s} ds = \int_{273}^{T} \frac{C_{w} dT}{T}$$

$$\therefore \qquad S_{f} = C_{w} \cdot \ln\left(\frac{T}{273}\right) = 2.3 C_{w} \log\left(\frac{T}{273}\right)$$

$$S_{f} \text{ can be directly steen from the steam tables.}$$

ENTROPY INCREASES DURING EVAPORATION 7.11

When the water is completely evaporated into steam it absorbs full latent heat h_{fg} at constant temperature T, corresponding to the given pressure. We know that

Entropy =
$$\frac{\text{heat absorbed}}{\text{absolute temperature}}$$

Increase of entropy during evaporation

$$S_{fg} = \frac{h_{fg}}{T}$$

If the steam is wet with dryness fraction x, the evaporation will be partial. In such a case heat absorbed = xh_{fg}

 \therefore Increase in entropy $S_{fg} = \frac{xh_{fg}}{T}$

7.12 ENTROPY OF WET AND DRY STEAM

The entropy of wet and dry steam above the freezing point of water is the entropy of water plus the entropy during evaporation.

: Increase in entropy

$$= S_{f} + \frac{xh_{fg}}{T}$$

$$= S_{f} + xs_{fg} \qquad (for wet steam)$$

$$= S_{f} + s_{fg}$$

$$= S_{g} \qquad (for dry steam)$$

 S_{g} can directly be read from steam tables.

7.13 ENTROPY FOR SUPERHEATED STEAM

During super heating heat is supplied at constant pressure and the temperature of dry steam (T) increases to the temperature of superheated steam (T_{sup}) For a small rise in temp dT the heat absorbed.

$$\delta q = 1 \times C_p \cdot dT$$

$$\frac{\delta q}{T} = \frac{C_p \cdot dT}{T}$$

$$dS = C_p \cdot \frac{dT}{T}$$

 C_p = specific heat of super heated steam may be taken between 1.67 to 2.5 kJ/kg K.

 \therefore Total increase in entropy from T K to T_{sup} K.

$$\int_{S_{g}}^{S_{sup}} ds = C_{p} \int_{T}^{T_{sup}} \frac{dT}{T}$$
$$S_{sup} - S_{g} = C_{p} \cdot \ln\left(\frac{T_{sup}}{T}\right)$$

:.

$$\therefore \qquad (S_{\sup} - S_g) = 2.3 \ C_p \cdot \log\left(\frac{T_{\sup}}{T}\right)$$

 $(S_{sup} - S_g)$ = increase in entropy of superheated steam.

: For 1 kg of super heated steam

$$S_{sup} = S_g + 2.3 C_p \log\left(\frac{T_{sup}}{T}\right)$$
$$= \left(S_f + \frac{h_{fg}}{T}\right) + 2.3 C_p \log\left(\frac{T_{sup}}{T}\right)$$

7.14 EXTERNAL WORK OF EVAPORATION

Latent heat of vaporisation supplied during the evaporation not only changes the phase of the substance but also does an external work in moving the piston at constant pressure due to increase in volume from v_t to v_o .

: External work of evaporation per kg of dry saturated steam

$$= p (v_g - v_f) kJ$$
$$= p v_g kJ$$

[::
$$v_f$$
 is very small at low pressure]

: External work of evaporation per kg of wet steam

$$= pxv_a kJ$$

And external work of evaporation per kg of superheated steam

$$= pv_{sup} kJ$$

7.15 INTERNAL LATENT HEAT

It is obtained by subtracting the external work of evaporation from the latent heat of evaporation.

: Internal latent heat of dry saturated steam

$$= (h_{fg} - pv_g) \text{ kJ/kg}$$

7.16 INTERNAL ENERGY OF STEAM

The internal energy of steam is defined as the difference between the enthalpy of the steam and the external work of evaporation.

: Internal energy of dry saturated steam

$$u_g = h_g - pv_g \,\mathrm{kJ/kg}$$

Internal energy of wet steam

$$u = (h_f + xh_{fo} - pxv_o) \text{ kJ/kg}$$

Internal energy of superheated steam

$$u_{\rm sup} = h_{\rm sup} - pv_{\rm sup} \, {\rm kJ/kg}$$

7.17 TEMPERATURE-ENTROPY DIAGRAM OF WATER AND STEAM

It is useful in solving the problem on adiabatic expansion and compression of steam. The abscissa of the diagram represents the entropy of the vertical ordinate showing the values of temperature as shown in figure 7.7. Consider 1 kg of water and steam above the freezing point of water at constant pressure. When heat is added to the water entropy will increase where the logarithmic curve increase in entropy = $S_f = AB$.

On further heating water starts evaporating and receives heat at constant temperature T. Entropy goes on increasing till the entire *latent heat* required to evaporate 1 kg of water has been supplied, increase in entropy $S_{fg} = BC$.



If we draw a family of similar curves at different pressure and plot the points B_1 , B_2 and C_1 , C_2 etc., then the line joining the points A, B_1 , B_2 , B_3 etc. is called water line. Similarly, the line joining the points C, C_1 , C_2 , C_3 etc. is called dry steam line. The point where the water and dry steam line meets (P) is called *critical point*, it is represented by 374.15°C.

7.18 ISOTHERMAL LINES ON T-S DIAGRAM

We know that there is no change of temperature during isothermal process.

1 kg of wet steam at 250°C, x = 0.2 at A expanded isothermally up to x = 0.8 at B

: increase in entropy

$$= S_b - S_a$$

= 5.42 - 3.45
= 1.97 kJ/kg. K.

Isothermal line *AB* is shown in figure 7.9.

7.19 ISENTROPIC LINES

There is no increase or decrease in enthalpy or total heat during an isentropic process. There is no change in entropy during the process.



Fig. 7.9

Consider a sample of 1 kg of wet steam at temperature of 310°C and dryness fraction of 0.8 at *A*. Let the sample expanded isentropically till its dryness fraction 0.6 marks *B*.

$$T_{A} = 310^{\circ}\text{C}, T_{B} = 40^{\circ}\text{C}$$

Fall in temperature = 310 - 40

 $= 270^{\circ}C$

It is shown in figure 7.10 on *t-s* diagram.

7.20 ENTHALPY-ENTROPY (h-s) DIAGRAM FOR WATER AND STEAM OR MOLLIER CHART





Mollier chart is a graphical representation of the steam tables, in which the enthalpy (h) is plotted along the ordinate and entropy (s) along abscissa. First of all, enthalpy and entropy of water and dry saturated steam, for any particular pressure, are obtained from the steam tables. These values of enthalpies and entropies are plotted and then liquid line and dry saturated line is obtained. Both these line meet at *C*, *i.e.*, the critical point as shown in figure 7.11. The critical point corresponds to the enthalpy of liquid and dry saturated steam at 221.2 bar. The enthalpy-entropy chart like temperature-

entropy chart, is also very useful in solving the problems in adiabatic or isentropic expansion and compression of steam. In actual diagram, abscissa of the diagram represents the entropy of 1 kg of water and steam, *i.e.*, specific entropy above the freezing point of water and the vertical ordinate shows the values of specific enthalpy *i.e.*, total heat, as shown in figure 7.12.

The diagram is divided into two parts by a line termed as saturation line. Upper region of the saturation line is called superheated region where temperature of steam increases at given pressure and lower region of saturation line is called wet region where temperature of steam remains constant at a given pressure.

The Mollier diagram has the following lines

- 1. Dryness fraction lines
- 2. Constant volume (i.e., specific volume) line
- 3. Constant pressure lines
- 4. Isothermal lines
- 5. Isentropic lines and
- 6. Throttling line



Fig. 7.12

7.21 DRYNESS FRACTION LINES ON H-S DIAGRAM

Below the saturation line is known as wet region, so the dryness fraction lines are drawn only below the saturation (which represents dryness fraction equal to 1.0). Ξ These lines represent the condition of wet $\frac{\Delta}{B}$ H steam between various values of enthalpy and entropy as shown in figure 7.13. For example, a sample of 1 kg of wet steam of dryness of 0.9 is represented by the line *AB* which is shown in figure 7.13. If *A* and *B* be the initial and final stage, then enthalpy at point *A* is $h_A = 2500$ kJ/kg and enthalpy at point *B* is $h_B = 2400$ kJ/kg is shown in figure 7.13.

7.22 CONSTANT VOLUME LINE

The constant volume lines are drawn in both the wet region and superheated steam region. This lines are straight in the wet steam region, but curved upwards above the saturation curve, *i.e.*, superheated region as shown in figure 7.14 by parts *AB* and *CD* of line of constant volume 1.0.

7.23 CONSTANT PRESSURE LINE

Constant pressure lines are drawn in both upper and lower region of saturated steam line. These lines are straight in the wet region because during vaporisation the increase of enthalpy is directly proportional to the increase in quality, and hence to the increase in entropy. It is curve and it is closer to wet region than superheated region is shown in figure 7.15.

7.24 ISOTHERMAL LINE

The isothermal lines or constant pressure lines are drawn only above the saturation line. These lines represent the condition of

superheated steam between various values of enthalpy and entropy, as shown in figure 7.16 by the line AB.





7.25 ISENTROPIC LINE ON (H-S) DIAGRAM

Isentropic process is a reversible adiabatic process. In this process entropy is constant. So isentropic line is parallel to vertical axis by the line *AB* which is shown in figure 7.17.



7.26 THROTTLING LINES ON H-S DIAGRAM

In throttling process enthalpy before throttling is equal to enthalpy after throttling *i.e.*, there is no change of enthalpy. So enthalpy line *AB* will be parallel to horizontal axis as shown in figure 7.18.



Multiple Choice Questions

Choose the correct answer

- (a) Specific volume of water decreases on freezing; (b) Boiling point of water decreases with increasing pressure; (c) Specific volume of CO₂ increases on freezing; (d) Freezing temperature of water decreases with increasing pressure.
- (a) The slope of vaporisation curve is always negative; (b) The slope of vaporisation curve is always positive; (c) The slope of sublimation curve in negative for all pure substances; (d) The slope of fusion curve is positive for all pure substances; (e) The slope of fusion curve is positive for all pure substances.
- 3. (a) The process of passing from liquid to vapour is condensation; (b) An isothermal line is also a constant pressure line during wet region; (c) Pressure and temperature are independent during phase change; (d) The term dryness fraction is used to describe the fraction by mass of liquid in the mixture of liquid water and water vapour.
- 4. The latent heat of vaporisation at critical point is

(a) less than zero; (b) greater than zero; (c) equal to zero; (d) none of above.

- (a) Critical point involves equilibrium of solid and vapour phases; (b) Critical point involves equilibrium of solid and liquid phases; (c) Critical point involves equilibrium of solid, liquid and vapour phases;
 (d) Triple point involves equilibrium of solid, liquid and vapour phases.
- 6. With the increase in pressure

(a) boiling point of water increases and enthalpy of evaporation increases.

7. With increase in pressure

(a) enthalpy of dry saturated steam increases;(b) enthalpy of dry saturated steam decreases.(c) enthalpy of dry saturated steam remains same;(d) enthalpy of dry saturated steam first increases and then decreases.

8. Dryness fraction of steam is defined as

(a) mass of water vapour in suspension/(mass of water vapour in suspension + mass of dry steam);
(b) mass of dry steam/mass of water vapour in suspension;
(c) mass of dry steam/(mass of dry steam + mass of water vapour in suspension);
(d) mass of water vapour in suspension/mass of dry steam.

9. The specific volume of water when heated at 0°C

(a) first increases and then decreases;(b) first decreases and then increases;(c) increases steadily

10. Only throttling calorimeter is used for measuring:

(a) Very low dryness fraction up to 0.7; (b) Very high dryness fraction up to 0.98; (c) dryness fraction of only low pressure steam; (d) dryness fraction of only high pressure steam

11. Heat of superheated steam is given by

(a)
$$h_{sup} = h_f + h_{fg} + C_{ps} \log_e \frac{T_{sup}}{T_s}$$
; (b) $h_{sup} = h_f + xh_{fg}$; (c) $h_{sup} = h_f + h_{fg}$; (d) $h_{sup} = h_f + xh_{fg} + C_{ps} \log_e \frac{T_s}{273}$

- 12. Volume of wet steam (per kg) with dryness fraction x is given by (a) x^3v_{g} ; (b) xv_{f} ; (c) $x^2(v_g - v_f)$; (d) x^2v_{g} ; (e) none of these
- 13. Internal latent heat is given by

(a)
$$h_{fg} - \frac{pv_g}{J}$$
; (b) $h_g - \frac{pv_g}{J}$; (c) $h_{sup} - \frac{pv_f}{J}$; (d) $h_{fg} + \frac{pv_g}{J}$; (e) none of the above

14. Entropy of 1 kg of water at TK is given by

(a) $C_{pw} \log_e \frac{T}{273}$; (b) $C_{pw} \log_e \frac{T_2}{T_1}$; (c) $C_{pw} \log_e \frac{T}{273}$; (d) $C_{pw} \log_e \frac{T_2}{273}$; (e) none of the above

15. Entropy of wet steam (1 kg) is given by

(a)
$$s_f + \frac{xh_{fg}}{T_s}$$
; (b) $s_g + \frac{xh_{fg}}{T_s}$; (c) $s_f + \frac{h_{fg}}{T_s}$; (d) $s_f + C_{ps} \log_e \frac{T_{sup}}{T_s}$; (e) none of the above

16. In throttling process

(a)
$$h_1^2 = h_2$$
; (b) $h_1 = h_2$; (c) $h_1 = h_2 + \frac{h_{fg}}{T_s}$; (d) $h_2 = h_1 + \frac{h_{fg}}{T_s}$; (e) none of the above

17. In isentropic process

(a)
$$w = 2(u_2 - u_1)$$
; (b) $w = (u_2 - u_1)^2$; (c) $w = u_2 - u_1$; (d) $w = (u_2 - u_1)^{1/2}$; (e) none of the above

Answers

1. (d) 2. (a) 3. (b) 4. (c) 5. (d) 6. (b) 7. (b) 8. (c) 9. (b) 10. (b) 11. (a) 12. (e) 13. (a) 14. (a) 15. (a) 16. (b) 17 (c)

NUMERICAL EXAMPLES

EXAMPLE 1

Determine the quantity of heat required to produce 1 kg of steam at a pressure of 6 bar at a temperature of 25°C under the following conditions.

(a) When the steam is wet having a dryness fraction 0.9, (b) When the steam is dry satuarted, (c) When it is super heated at a constant pressure at 250°C. Assume the mean specific heat of superheated steam to be 2.3 kJ/kg.

SOLUTION

Here,

$$p = 6 \text{ bar}, t_w = 25^{\circ}\text{C}, x = 0.9$$

 $t_{sup} = 250^{\circ}\mathrm{C}$

Following values are obtained from steam table for p = 6 bar

$$h_f = 670.4 \text{ kJ/kg}, h_{fg} = 2085 \text{ kJ/kg}$$

$$t = 158.8$$
°C

(a) When the steam is wet, entropy of total heat of 1 kg of wet steam

$$h = h_f + x h_{fg}$$

= 670.4 + 0.9 × 2085
= 2546.9 kJ

Since the water is at a temperature of 25°C therefore heat already in water

$=4.2 \times 25 = 105 \text{ kJ}$

Heat actually required

= 2546.9 - 105 = 2441.9 kJ

(b) When the steam is dry, saturated entropy or total heat of 1 kg of dry saturated steam

$$h_g = h_f + h_{fg}$$

= 670.4 + 2085 = 2755.4 kJ

Heat actually required

$$=2755.4 - 105 = 2650.4 \, \text{kJ}$$

(c) When the steam is superheated, enthalpy or total heat of 1 kg of superheated steam

$$h_{sup} = h_g + C_p (t_{sup} - t)$$

= 2755.4 + 2.3 [(250 + 273) - (158.8 + 273)]
= 2965.16 kJ

Heat actually required

 $= 2965.16 - 105 = 2860.16 \, \text{kJ}$

EXAMPLE 2

Steam enters an engine at a pressure of 12 bar with a 67°C of superheat. It is exhausted at a pressure of 0.15 bar and 0.95 dryness fraction. Find the drop in entropy of the steam. $C_p = 2.04 \text{ kJ/kg}$

SOLUTION

Here,

$$p_1 = 12 \text{ bar}, (t_{sup} - t) = 67^{\circ}\text{C}$$

$$p_2 = 0.15 \text{ bar}, x = 0.95$$
From steam table corresponding to a pressure of 12 bar, we have

$$h_f = 798.4 \text{ kJ/kg}, h_{fg} = 1984.3 \text{ kJ/kg}$$
We know that enthalpy or total heat of 1 kg of superheated stream

$$h_{sup} = h_f + h_{fg} + C_p (t_{sup} - t) = 798.4 + 1984.3 + 2.04 \times 67 = 2919.38 \text{ kJ/kg}$$
Similarly for pressure 0.15 bar from steam tables

$$h_f = 226 \text{ kJ/kg}, h_{fg} = 2373.2 \text{ kJ/kg}$$
We know enthalpy or total heat of 1 kg of wet steam

$$h = h_f + xh_{fg} = 226 + 0.95 \times 2373.2 = 2480.54 \text{ kJ/kg}$$
Drop in enthalpy of the steam

$$=h_{sup}-h=(2919.38-2480.54)=438.84$$
 kJ/kg

EXAMPLE 3

A steam engine obtains steam from a boiler at a pressure of 15 bar and 0.98 dry. It was observed that the steam loses 21 kJ of heat per kg as it flows through the pipeline, pressure remaining constant. Calculate dryness fraction of steam at the engine end of the pipeline.

SOLUTION

Given p = 15 bar, heat loss = 21 kJ/kg.

From steam tables corresponding to the pressure of 15 bar, we have $h_{fg} = 1945.3 \text{ kJ/kg}$ $h_{f} = 844.6 \, \text{kJ/kg}$ and Enthalpy of wet steam at the boiler end *.*.. $h_1 = h_f + x \cdot h_{fg}$ $=844.6 \pm 0.98 \times 1945.3$ $= 2751 \, \text{kJ/kg}$ As steam loses 21 kJ/kg of steam, enthalpy of wet steam at the engine end $h_2 = h_f + x_2 h_{fg}$ and $h_2 = 2751 - 21 = 2730 \text{ kJ}$ $2730 = 844.6 + x_2 \times 1945.3$ *:*.

... $x_2 = 0.97$

Therefore, at the end of pipeline dryness fraction is 0.97.

EXAMPLE 4

Calculate the enthalpy of 1 kg of steam at a pressure of 8 bar and dryness fraction of 0.8. How must heat would be rejected to raise 2 kg of its steam from water at 20°C?

and

SOLUTION

p = 8 barHere, and x = 0.8Enthalpy of 1 kg of steam, from steam table $h_f = 720.9 \, \text{kJ/kg}$ $h_{fg} = 2046.5 \text{ kJ/kg}$

:..

Heat

$$\dot{h} = h_f + x h_{fg}$$

= 720.9 + 2046.5 × 0.8
= 2358 1 kJ

Heat required to rise 2 kg of this steam from water at 20°C.

already in water =
$$4.2 \times 20 = 84$$
 kJ

$$= 2358.1 - 84$$

 $= 2274.1 \, \text{kJ}$

: Heat required for 2 kg of steam

 $= 2 \times 2274.1 = 4548.2 \text{ kJ}$

EXAMPLE 5

Find the entropy of 1 kg of dry saturated steam at a pressure of 5.2 bar. The boiling point of water at this pressure is given by 152.6°C and the total heat at this temperature is utilised.

SOLUTION

 $T = 152.6^{\circ}\text{C} = 152.6 + 273 = 425.6 \text{ K}$ Here, p = 5.2 barand From steam table, $h_{fg} = 2102.7 \text{ kJ/kg}$ We know entropy of 1 kg of dry saturated steam

$$S_g = S_f + \frac{h_{fg}}{T}$$

$$= 2.3 C_{w} \log\left(\frac{T}{273}\right) + \frac{h_{fg}}{T}$$
$$= 2.3 \times 4.2 \log\left(\frac{425.6}{273}\right) + \frac{2102.7}{425.6}$$
$$= 1.863 + 4.94$$
$$= 6.803 \text{ kJ/kg K}$$

EXAMPLE 6

Calculate the entropy of 1 kg of wet steam with dryness fraction of 0.9 at a pressure of 8.4 bar.

SOLUTION

Here, x = 0.9 and p = 8.4 bar From steam table for pressure 8.4 bar, boiling point is 172.4°C \therefore T = 172.4 + 273 = 445.4 K

$$h_{fr} = 2039.6 \text{ kJ/kg}$$
 $S_f = 2.066 \text{ kJ/kgK}$

For 1 kg of wet steam, entropy

$$S = S_f + \frac{xh_{fg}}{T}$$

$$= 2.066 + 4.12 = 6.186 \text{ kJ/kg K}$$

Also $S = S_f + xS_{fg}$ directly can be read from steam table.

EXAMPLE 7

In a Carnot cycle, heat is supplied at 350°C and is rejected at 25°C. The working fluid is water, which while receiving heat, evaporates from liquid at 350°C to steam at 350°C. From the steam tables the entropy change for this process 1.438 kJ/kg. If the cycle operates on a stationary mass of 1 kg of water, find the heat supplied, work done and heat rejected per cycle, what is the pressure of water during heat reception?

SOLUTION

Here,

$$(S_2 - S_1) = 1.438 \, \text{kJ/kg K}$$

Heat supplied per cycle

$$=(S_2-S_1)T_1=1.438\times 623=895.87$$
 kJ/kg

Work done per cycle

$$=(S_2 - S_1)(T_1 - T_2) = 1.438(623 - 298) = 467.35 \text{ kJ/kg}$$

 $T_1 = 350^{\circ}\text{C} = 350 + 273 = 623 \text{ K}, T_2 = 25^{\circ}\text{C} = 25 + 273 = 298 \text{ K}$

Heat rejected per cycle

$$=(S_2 - S_1) T_{200} = 1.438 \times 298 = 428.52 \text{ kJ/kg}$$

From the steam tables corresponding to 350°C the pressure is 165.35 bar.

EXERCISE

- 1. What is pure substance?
- 2. Draw and explain a p-T (pressure-temperature) diagram for a pure substance.
- 3. What is a triple point?
- 4. Explain with a neat diagram $p \forall -T$ surface.
- 5. Does wet steam obey laws of perfect gases?
- 6. Describe the process of formation of steam and give its graphical representation also.
- 7. Explain the following terms relating to steam formation:

(i) Sensible heat of water (ii) Latent heat of steam (iii) Dryness fraction of steam (iv) Enthalpy of wet steam (v) Superheated steam.

- 8. What advantages are obtained if superheated steam is used in steam prime movers?
- 9. What do you ment by following:

(i) Internal latent heat (ii) Internal energy of steam (iii) External work of evaporation (iv) Entropy of evaporation (v) Entropy of wet steam (vi) Entropy of superheated steam.

- 10. Write a short note on Mollier Chart.
- 11. 300 tonnes per second of steam is expanded in a turbine from an initial pressure of 90 bar to 0.1 bar. The specific enthalpy of steam at inlet and exit of the turbine are 3300 kJ/kg and 2200 kJ/kg, respectively. Neglecting potential energy and kinetic energy terms and heat transfer loss, find the output of the turbine in MW.
- 12. A 0.025 m³ vessel contains 0.3 kg of steam at 2 MPa. Determine the quality and enthalpy of the steam.
- 13. In a throttling calorimeter, the steam produced by a boiler is admitted at a pressure of 17 bar absolute. After throttling, the steam is discharged into the atmosphere at 120°C. Find the quality of boiler steam if the barometer reads 760 mm of Hg. Take for superheated steam, C_p as 2.1 kJ/kgK. Show the process diagrammatically.
- 14. Steam at 6 bar absolute from a fire-tube boiler is passed through a throttling calorimeter. After throttling, the steam is let out into the atmosphere at 120°C. If the barometer reading is 750 mm of Hg, find the change of entropy per kg of steam while C_n is 2.1 kJ/kgK for superheated steam.
- 15. 0.5 kg of steam at 00 kPa and 0.98 dry is superheated at constant pressure or 300°C. Using steam table find the amount of heat added to the steam in the superheater. Take C_p as 2.1 kJ/kgK for superheated steam.
- Determine the quantity of heat required to produce 1 kg of steam at a pressure of 6 bar at 25°C when steam is (i) wet with 0.90 dry (ii) dry saturated (iii) superheated at constant 250°C. Assume mean specific heat for superheated steam 2.3 kJ/kgK.
- 17. A steam engine obtains steam from a boiler at 15 bar and 0.98 dry. While flowing through pipeline at constant pressure, rate of heat loss is 2.1 kJ/kg. What is dryness fraction at the end of pipeline?
- 18. At 300°C and 20 bar pressure, what is the volume of 1 kg superheated steam?
- 19. Feed water is supplied to the boiler at 45°C and is converted into steam at 5.5 bar pressure at 188°C. Assuming relevant data, compute rate of heat supply.

- 20. At constant 18 bar pressure 0.9 dry steam is heated until dry and saturated. What is increase in volume, heat supplied, and work done per kg of steam? Now volume remaining constant, what amount of heat is to be extracted for reducing pressure to 14 bar?
- 21. What is the mass of 0.8 dry steam of 0.50 m^3 at 4 bar? What is enthalpy of 1 m³ of steam?
- 22. 0.8 dry steam is produced at 8 bar. Calculate (i) external work done during evaporation (ii) internal latent heat of steam.
- 23. Determine internal energy per kg of superheated steam at 14 bar and 573 K. If the steam is expanded to 1.4 bar and 0.8 dry, what is the change in internal energy?

This page intentionally left blank



THERMODYNAMIC AIR STANDARD CYCLES

8.1 INTRODUCTION

Thermodynamic cyclic process consists of a series of thermodynamic processes, which takes place in a certain order, and the initial conditions are restored at the end of the processes. The study of various thermodynamic cycles is very much essential for the power developing system.

A cycle, which requires four piston strokes and two complete revolutions of the crank, is known as *four stroke cycle*.

When a cycle requires only two piston strokes and one revolution of the crank, is known as *two stroke cycle*. When air is assumed to be the working fluid (substance) inside the engine cylinder, the cycle is called as an *air cycle*.

8.2 ASSUMPTIONS IN THERMODYNAMIC CYCLES

The analysis of all thermodynamic cycles or air cycles is based on the following assumptions:

- 1. The gas in the engine cylinder is a perfect gas, *i.e.*, it obeys the gas laws and constant specific heats.
- 2. All the expansion and compression are isentropic or reversible adiabatic and takes place without any internal friction.
- 3. The properties of the working substance can be calculated by the application of the characteristic equation for a perfect gas *i.e.*, $p \forall = mRT$

or

$$= \frac{p\forall}{mT}$$
$$= \frac{R_m}{M} = \frac{\text{universal gas constant}}{\text{molecular weight of gas}}$$

$$= \frac{8314}{29} \approx 287 \text{ J/kg K}$$

= 0.287 kJ/kg K for air

This value of R are same at moderate temperatures.

R

- 4. The working substance is a fixed mass of air either contained in a closed system or flowing at constant rate round a closed circuit.
- 5. The kinetic and potential energies of working substances are neglected.
- 6. The gas does not undergo any chemical changes, the addition of heat from outside source and removal of heat to outside sink is by simple heat transfer.

8.3 CLASSIFICATION OF THERMODYNAMIC CYCLES

The thermodynamic cycles, in general, may be classified into the following two types;

(1) Reversible or ideal and (2) Irreversible or natural or real cycle. These have already been discussed, in detail, in chapter 5.

8.4 IMPORTANT PARAMETERS IN AIR STANDARD CYCLE ANALYSIS

(a) Air Standard Efficiency (Thermal Efficiency)

It is defined as the ratio of net work transfer during cycle to the net heat transfer to the cycle. It is denoted by η .

(b) Specific Work Transfer

It is work transfer per unit mass of the working substance (*w*)

(c) Specific Air Consumption

It is the quantity of working substance required for doing unit work transfer or the flow rate of working substance for unit power. It is inverse of specific work transfer.

(d) Work Ratio

It is the ratio of net work transfer in a cycle to the positive work transfer during the cycle. It is denoted by r_w and is given by

positive work transfer

net work transfer in a cycle

positive work transfer in a cycle

8.5 IMPORTANT TERMS USED IN THERMODYNAMIC CYCLES

- (a) Cylinder bore: The inner diameter of the cylinder, in which the piston moves, is known as cylinder bore.
- (b) Stroke length: The piston moves through the cylinder due to rotation of the crank. Its extreme positions are known as top dead centre (TDC) and bottom dead centre (BDC) respectively as shown in figure 8.1. The distance between these two extreme positions is known as stroke length or simply stroke.
- (c) Clearance volume: The volume occupied by the working fluid, when piston reaches the top dead centre, is known as clearance volume. It is generally denoted by \forall_c .

(d) Swept volume: The volume swept by the piston, when it moves between the two extreme positions is known as swept volume or displacement volume or stroke volume. It is generally denoted by \forall_{s} .

$$\therefore \qquad \forall_s = \text{piston area} \times \text{stroke length} = \frac{\pi d^2}{4} \times l$$



- (e) Total cylinder volume: The volume occupied by the working fluid, when the piston is at the bottom dead centre, is known as total cylinder volume. Total cylinder volume (\forall) is equal to the sum of clearance volume \forall_c and swept volume (\forall_s). $\therefore \qquad \forall = \forall_c + \forall_s \qquad (8.1)$
- (f) Compression ratio: The ratio of total cylinder volume to clearance volume is known as compression ratio. It is denoted generally by (r)

$$\therefore \qquad r = \frac{\forall_c + \forall_s}{\forall_s} = 1 + \frac{\forall_s}{\forall_c}$$

- (g) Expansion ratio: It is the ratio of cylinder to the volume of cylinder when gas starts to expand. Sometimes it is equal to compression ratio.
- (h) Cut-off ratio: It is the ratio of the volume at the point of cut off of fuel to the clearance volume. It is generally denoted by ρ

$$\rho = \frac{\text{cut-off volume}}{\text{clearance volume}}$$
(8.2)

(i) Mean effective pressure: As a matter of fact, pressure in the cylinder keeps on changing with the position of the piston. For all sorts of calculations, we need the mean effective pressure, which may be defined as the constant pressure acting on the piston during the working stroke. It will be able to do the same amount of work, as done by the actual varying pressure, produced during the cycle. Mean effective pressure is the ratio of work done to the stroke volume or piston displacement volume. It is denoted by p_m

$$p_m = \frac{\text{work done}}{\text{wtroke volume}}$$
(8.3)

(j) Efficiency of cycle: It may be defined as the ratio of work done to the heat supplied during a cycle. Mathematically

Efficiency of cycle
$$(\eta) = \frac{\text{work done}}{\text{heat supplied}}$$

Again work done is equal to heat supplied minus the heat rejected, the efficiency of cycle therefore may also be expressed

$$\eta = \frac{\text{heat supplied} - \text{heat rejected}}{\text{heat supplied}}$$
(8.4)

This efficiency is the theoretical efficiency of the cycle. Therefore it is known as *theoretical thermal efficiency*. This efficiency does not take into account the practical losses, which occur in running of the engine. When air is assumed to be the working substance inside the engine cylinder, the efficiency obtained is known as *air standard efficiency*. It is also called *ideal efficiency*.

(k) Indicated thermal efficiency: It is the ratio of the heat equivalent to one kW hour to the heat in fuel per I.P. hour. Mathematically it is

$$\eta_{t} = \frac{\text{heat equivalent to 1 kWh}}{\text{heat in fuel per I.P. hour}}$$

$$= \frac{3600}{\frac{(m_{f} \cdot C)}{I.P.}} = \frac{3600 I.P.}{m_{f} \cdot C}$$
(8.5)

Here m_f is mass of fuel in kg/h, C is calorific value of fuel in kJ/kg and I.P. is the indicated power. (I) **Relative efficiency:** It is also known as efficiency ratio. The relative efficiency of an I.C.

engine is the ratio of indicated thermal efficiency to air standard efficiency. Mathematically

$$\eta_R = \frac{\text{indicated thermal efficiency}}{\text{air standard efficiency}}$$
(8.6)

8.6 TYPES OF THERMODYNAMIC CYCLES

Though there are many types of thermodynamic cycles, yet the following are important from the subject point of view:

(1) Carnot cycle (2) Stirling cycle (3) Ericsson cycle (4) Joule cycle (5) Otto cycle (6) Diesel cycle and (7) Dual combustion cycle.

8.7 CARNOT CYCLE

We have already discussed the Carnot cycle in chapter 5.

8.8 OTTO CYCLE

The first successful engine working on this cycle was built by A. Otto. These days many gas, petrol and many of the oil engines run on this cycle. It is also known as *constant volume cycle*. Here heat is received and rejected at constant volume. In comparison to the other cycles the air is assumed to be the working substance.



Fig. 8.2

The ideal otto cycle consists of two constant volume and two reversible adiabatic or isentropic processes as shown in p- \forall and T-S diagram in figure 8.2. Let engine cylinder contains m kg of air at point 1 and at this point pressure = p_1 , volume = \forall_1 , temperature = T_1 .

1. First stage (reversible adiabatic or isentropic expansion)

The air is expanded reversibly and adiabatically from temperature T_1 to T_2 . In this process no heat is absorbed or rejected by the air as shown by (1–2).

2. Second stage (constant volume cooling)

The air is cooled at constant volume from temp T_2 to T_3 , shown by (2–3). Heat rejected by air during this process $Q_{2-3} = mC_v (T_2 - T_3)$

3. Third stage (reversible adiabatic or isentropic compression)

The air is compressed reversibly and adiabatically from temperature T_3 to T_4 as shown by the curve (3–4). Here no heat is absorbed or rejected by the air.

4. Fourth stage (constant volume heating)

The air is now heated at constant volume from temp T_4 to T_1 , given by curve (4–1). Heat absorbed by the air during this process $Q_{4-1} = mC_{\nu} (T_1 - T_4)$

We see that the air has been brought back to its original condition of pressure volume and temperature, thus completing the cycle.

:. Work done = heat absorbed – heat rejected

$$= mC_{v} (T_{1} - T_{4}) - mC_{v} (T_{2} - T_{3})$$

$$\therefore \qquad \eta_{\text{otto}} = \frac{\text{work done}}{\text{heat absorbed}}$$

$$= \frac{mC_{v} (T_{1} - T_{4}) - mC_{v} (T_{2} - T_{3})}{mC_{v} (T_{1} - T_{4})}$$

$$= 1 - \frac{T_{2} - T_{3}}{T_{1} - T_{4}}$$

$$= 1 - \frac{T_{3} (T_{2} / T_{3} - 1)}{T_{4} (T_{1} / T_{4} - 1)} \qquad (8.7)$$

Also for reversible adiabatic compression process (3–4), we know

$$\left(\frac{T_3}{T_4}\right) = \left(\frac{\forall_4}{\forall_3}\right)^{\gamma-1} = \left(\frac{1}{r}\right)^{\gamma-1} \qquad \left[\because r = \frac{\forall_3}{\forall_4} = \frac{\forall_2}{\forall_1}\right]$$

$$\cdot \qquad \frac{T_2}{T_1} = \frac{T_3}{T_4} = \left(\frac{1}{r}\right)^{\gamma-1}$$

$$\cdot \qquad \frac{T_1}{T_4} = \frac{T_2}{T_3}$$

$$\therefore \qquad \eta_{\text{otto}} = 1 - \frac{T_3}{T_4} = 1 - \frac{T_2}{T_1} = 1 - \frac{1}{(r)^{\gamma - 1}}$$
(8.8)

8.9 JOULE'S CYCLE

It consists of two constant pressure and two reversible adiabatic or isentropic processes. It is shown in $p - \forall$ and T - S diagram in figure 8.3. Let the engine cylinder contains m kg air and its pressure = p_1 , volume = \forall_1 , temperature = T_1



1. First stage (constant pressure heating)

Air is heated at constant pressure from initial temperature T_1 to final temperature T_2 . So heat $Q_{1-2} = mC_p (T_2 - T_1)$ supplied to the air

2. Second stage (isentropic expansion) Air is allowed to expand reversibly and adiabatically from \forall_2 to \forall_3 . Temperature falls from T_1 to

 T_3 and no heat is absorbed or rejected.

3. Third stage (constant pressure cooling)

The air is now cooled at constant pressure from temperature T_3 to temperature T_4 . Heat rejected

by the air $Q_{3-4} = mC_p (T_3 - T_4)$ 4. Fourth stage (isentropic compression)

 η_{ioule}

Temperature increases from T_4 to T_1 . No heat is absorbed or rejected by the air. No interchange of heat takes place during two adiabatic processes. So work done by constant pressure process

= Heat supplied – Heat rejected
=
$$mC_p (T_2 - T_1) - mC_p (T_3 - T_4)$$

= $mC_p [(T_2 - T_1) - (T_3 - T_4)]$
= $\frac{\text{Work done}}{\text{Heat supplied}}$

:..

$$= \frac{mC_p \left\{ (T_2 - T_1) - (T_3 - T_4) \right\}}{mC_p (T_2 - T_1)}$$
$$= 1 - \frac{T_3 - T_4}{T_2 - T_1}$$

$$=1-\frac{T_{3}\left(1-T_{4}/T_{3}\right)}{T_{2}\left(1-T_{1}/T_{2}\right)}$$
(8.9)

Again for isentropic expansion process

$$\frac{T_3}{T_2} = \left(\frac{\forall_2}{\forall_3}\right)^{\gamma-1} = \left(\frac{p_3}{p_2}\right)^{\frac{\gamma-1}{\gamma}}$$

For isentropic compression process

$$\frac{T_4}{T_1} = \left(\frac{\forall_1}{\forall_4}\right)^{\gamma-1} = \left(\frac{p_4}{p_1}\right)^{\frac{\gamma-1}{\gamma}} \\
\frac{T_3}{T_2} = \frac{T_4}{T_1} \\
\frac{T_3}{T_2} = \frac{T_4}{T_3} \qquad [\because p_1 = p_2 \text{ and } p_3 = p_4] \\
\eta_{\text{joule}} = 1 - \frac{T_3}{T_2} = 1 - \frac{T_4}{T_1} = 1 - \frac{1}{r^{\gamma-1}}$$
(8.10)

8.10 DIESEL CYCLE

As

:.

:.

This cycle is devised by Dr Rudloph Diesel with an idea to attain a higher thermal efficiency with a high compression ratio. It is also known as *constant pressure cycle* as heat is received at constant pressure.

An ideal diesel cycle consists of two reversible adiabatic or isentropic, a constant pressure and a constant volume processes as shown in $p-\forall$ and T-S diagram in figure 8.4. Let the engine contains m kg of air at point where, pressure = p_1 , volume = \forall_1 , temperature = T_1



1. First stage (constant pressure heating)

Air is heated at constant pressure from initial temperature T_1 to final temperature T_2 . So heat supplied to the air $Q_{1-2} = mC_p (T_2 - T_1)$. Since supply of heat is cut off at point 2 therefore it is known as *cut off point*.
2. Second stage (reversible adiabatic or isentropic expansion) The air is expanded reversibly and adiabatically from temperature T_2 to T_3 . No heat is absorbed

or rejected by the air.

3. Third stage (constant volume cooling)

Air is now cooled at constant volume from temperature T_3 to T_4 . So heat rejected by the air $Q_{3-4} = mC_v (T_3 - T_4)$

4. Fourth stage (reversible adiabatic or isentropic compression)

The air compressed reversibly and adiabatically from temperature T_4 to T_1 and in this process, no heat is absorbed or rejected thus the air has been brought to its original conditions completing the cycle.

Therefore,
Work done = heat absorbed – heat rejected

$$= mC_{p} (T_{2} - T_{1}) - mC_{v} (T_{3} - T_{4})$$

$$\eta_{\text{diesel}} = \frac{\text{Work done}}{\text{Heat supplied}}$$

$$= \frac{mC_{p} (T_{2} - T_{1}) - mC_{v} (T_{3} - T_{4})}{mC_{p} (T_{2} - T_{1})}$$

$$= 1 - \frac{C_{v} (T_{3} - T_{4})}{C_{p} (T_{2} - T_{1})} = 1 - \frac{1}{\gamma} \frac{(T_{3} - T_{4})}{(T_{2} - T_{1})}$$
(8.11)
Compression ratio $= r = \left(\frac{\forall_{4}}{\forall_{1}}\right)$
Cut off ratio $\rho = \frac{\forall_{2}}{\forall_{1}}$
Expansion ratio $r_{1} = \left(\frac{\forall_{3}}{\forall_{2}}\right) = \left(\frac{\forall_{4}}{\forall_{2}}\right)$

$$= \frac{\forall_{4}}{\forall_{1}} \times \frac{\forall_{1}}{\forall_{2}} = r \times \frac{1}{\rho} = \frac{r}{\rho}$$
At constant pressure heating process $1 - 2$

$$\frac{\forall_{1}}{T_{1}} = \frac{\forall_{2}}{T_{2}}$$

Similarly for isentropic expansion process (2-3)

$$\left(\frac{T_3}{T_2}\right) = \left(\frac{\forall_2}{\forall_3}\right)^{\gamma-1} = \left(\frac{1}{r_1}\right)^{\gamma-1} = \left(\frac{\rho}{r}\right)^{\gamma-1}$$

 $T_2 = T_1 \frac{\forall_2}{\forall_1} = T_1 \times \rho$

At

:.

:.

$$T_{3} = T_{2} \left(\frac{\rho}{r}\right)^{\gamma-1} = T_{1} \rho \left(\frac{\rho}{r}\right)^{\gamma-1}$$

For isentropic compression process (4–1)

$$\frac{T_1}{T_4} = \left(\frac{\forall_4}{\forall_1}\right)^{\gamma-1} = (r)^{\gamma-1}$$

 $\therefore \qquad \qquad T_1 = T_4 \cdot (r)^{\gamma - 1}$

$$\therefore \qquad T_2 = T_1 \rho = T_4 \cdot (r)^{\gamma - 1} \rho$$

$$T_{3} = T_{4} \left(r \right)^{\gamma - 1} \rho \left(\frac{\rho}{r} \right)^{\gamma - 1} = T_{4} \cdot \rho^{\gamma}$$

$$\therefore \qquad \qquad \eta_{\text{diesel}} = 1 - \frac{1}{\gamma} \left[\frac{T_3 - T_4}{T_2 - T_1} \right]$$

$$= 1 - \frac{1}{\gamma} \left[\frac{T_4 \cdot \rho^{\gamma} - T_4}{T_4 \cdot r^{\gamma - 1} \cdot \rho - T_4 \cdot r^{\gamma - 1}} \right]$$
$$= 1 - \frac{1}{(r)^{\gamma - 1}} \left[\frac{\left(\rho^{\gamma} - 1\right)}{\gamma \left(\rho - 1\right)} \right]$$
(8.12)

8.11 COMPARISON BETWEEN THE EFFICIENCY OF OTTO AND DIESEL CYCLE FOR SAME COMPRESSION RATIO

Efficiency of Otto cycles

$$\eta_{otto} = 1 - \frac{1}{r^{\gamma - 1}}$$

And efficiency of diesel cycle

$$\eta_{diesel} = 1 - \frac{1}{r^{\gamma - 1}} \left[\frac{\left(\rho^{\gamma} - 1 \right)}{\gamma(\rho - 1)} \right]$$

The efficiency of the ideal Diesel cycle is lower than that of Otto cycle, for the same compression ratio. This is due to the fact that the cut-off ratio ρ is always greater than unity and hence the term with in the large bracket of η_{diesel} increases with increase of cut-off ratio. Thus the negative term increase and the efficiency is reduced.

The Diesel cycle efficiency increases with decrease in cut-off ratio and approaches maximum (equal to Otto cycle efficiency) when the term within the large bracket is unity.

Multiple Choice Questions

- 1. The air standard otto cycle comprises:
 - (a) two constant pressure processes and two constant volume processes
 - (b) two constant pressure processes and two constant entropy processes
 - (c) two constant volume processes and two constant entropy processes
 - (d) none of the above
- 2. The air standard efficiency of otto cycle is given by:

(a)
$$\eta = 1 + \frac{1}{r^{\gamma + 1}}$$
 (b) $\eta = 1 - \frac{1}{r^{\gamma - 1}}$
(c) $\eta = 1 + \frac{1}{r^{\gamma - 1}}$ (d) $\eta = 2 - \frac{1}{r^{\gamma - 1}}$

- 3. The thermal efficiency of theoretical otto cycle:
 - (a) increase with increase in compression ratio
 - (b) increases with increase in isentropic index γ
 - (c) does not depend upon the pressure ratio
 - (d) follows all the above
- 4. The work output of theoretical otto cycle:
 - (a) increases with increase in compression ratio
 - (b) increases with increase in pressure ratio
 - (c) increases with increase in adiabatic index γ
 - (d) follows all the above
- 5. For same compression ratio:
 - (a) thermal efficiency of otto cycle is greater than that of diesel cycle
 - (b) thermal efficiency of otto cycle is less than that of diesel cycle
 - (c) thermal efficiency of otto cycle is same as that for diesel cycle
 - (d) thermal efficiency of otto cycle cannot be predicted
- 6. In air standard diesel cycle, at fixed compression ratio and fixed value of adiabatic index (γ):
 - (a) thermal efficiency increases with increase in heat addition cut-off ratio
 - (b) thermal efficiency decreases with increase in heat addition cut-off ratio
 - (c) thermal efficiency remains same with increase in heat addition cut-off ratio
 - (d) none of the above

Answers

1. (b) 2. (b) 3. (d) 4. (d) 5. (a) 6. (b)

NUMERICAL EXAMPLES

EXAMPLE 1

In an Otto cycle the temperature at the beginning and end of the isentropic compression are 316 K and 596 K respectively. Determine the air standard efficiency and the compression ratio, take $\gamma = 1.4$ SOLUTION

He

Here
$$T_3 = 316 \text{ K}, T_4 = 596 \text{ K}, \gamma = 1.4, r = \frac{\forall_3}{\forall_4}$$

We know for isentropic compression

$$\frac{T_3}{T_4} = \left(\frac{\forall_4}{\forall_3}\right)^{\gamma-1} = \frac{1}{(r)^{\gamma-1}} = \frac{1}{(r)^{1.4-1}} = \frac{1}{r^{0.4}}$$
$$(r)^{0.4} = \left(\frac{T_4}{T_3}\right) = \left(\frac{596}{316}\right)$$

...

:.

$$r = \left(\frac{596}{316}\right)^{1/0.4} = 4.885$$

Air standard efficiency

$$\eta = 1 - \frac{1}{(r)^{\gamma - 1}}$$

$$\eta = 1 - \frac{T_3}{T_4} = 1 - \frac{316}{596} = \frac{596 - 316}{596} = \frac{280}{596}$$

= 0.47

EXAMPLE 2

An air motor works on Joule's cycle between 5 bar and 1 bar. The temperature at the beginning of isentropic expansion is 773 K and at the beginning of isentropic compression is 293 K. Determine the work done per kg of air and ideal efficiency. Assume $C_p = 1 \text{ kJ/kgK}$ and $\gamma = 1.4$.

SOLUTION

Given

$$p_1 = p_2 = 5$$
 bar $p_3 = p_4 = 1$ bar
 $T_2 = 773$ K, $T_4 = 293$ K, $C_p = 1$ kJ/kg K, $\gamma = 1.4$.

We know that for isentropic expansion process

p

$$\frac{T_3}{T_2} = \left(\frac{p_3}{p_2}\right)^{\frac{\gamma-1}{\gamma}} = \left(\frac{1}{5}\right)^{\frac{1.4-1}{1.4}} = (0.2)^{0.286} = 0.6311$$
$$T_3 = T_2 \times 0.6311 = 773 \times 0.6311 = 488 \text{ K}$$

...

Similarly for isentropic compression process

$$\frac{T_4}{T_1} = \left(\frac{p_4}{p_1}\right)^{\frac{\gamma-1}{\gamma}} = \left(\frac{1}{5}\right)^{\frac{14-1}{1.4}} = (0.2)^{0.286} = 0.6311$$

$$\therefore \qquad T_1 = \frac{T_4}{0.6311} = \frac{293}{0.6311} = 464 \text{ K}$$

$$\therefore \text{ Heat supplied} = mC_p (T_2 - T_1) = 1 \times 1 (773 - 464) = 309 \text{ kJ}$$

$$\therefore \text{ Heat rejected} = mC_p (T_3 - T_4) = 1 \times 1 (488 - 293) = 195 \text{ kJ}$$

$$\therefore \text{ Work done} = \text{Heat supplied} - \text{Heat rejected} = (309 - 195) = 114 \text{ kJ}$$

$$\text{Ideal efficiency} = \frac{\text{Work done}}{\text{Heat supplied}} = \frac{114}{309} = 0.369$$

EXAMPLE 3

In a diesel engine the compression ratio is 13:1 and the fuel is cut off at 8% of the stroke. Find the air standard efficiency of the engine, take γ for air = 1.4.

SOLUTION

Given,
$$r = \text{compression ratio} = \frac{\forall_4}{\forall_1} = 13$$

Since the cut off takes place at 8% of the stroke, therefore volume at cut off,

$$\forall_2 = \forall_1 + 8\%$$
 of the stroke volume

$$= \forall_1 + 0.08 (\forall_4 - \forall_1)$$

Let us assume the clearance volume $\forall_1 = 1 \text{ m}^3$, So $\forall_4 = 13 \text{ m}^3$ \therefore $\forall_2 = \forall_1 + 0.08 (13 - 1) = 1 + 0.08 \times 12 = 1.96 \text{ m}^3$

=

Cut off ratio

$$=\frac{\forall_2}{\forall_1}=\frac{1.96}{1}=1.96$$

Efficiency

$$\eta = 1 - \frac{1}{(r)^{\gamma - 1}} \left[\frac{\rho^{\gamma} - 1}{\gamma(\rho - 1)} \right]$$
$$= 1 - \frac{1}{(r)^{1/4 - 1}} \left[\frac{1.96^{1.4} - 1}{1.4(r)^{1/4 - 1}} \right]$$

$$= 1 - \frac{1}{(13)^{1.4-1}} \left[\frac{1.96^{1.4} - 1}{1.4(1.96 - 1)} \right]$$
$$= 1 - 0.417 = 0.583$$

EXAMPLE 4

In an ideal diesel cycle, the temperature at the beginning and end of compression are 57°C and 603°C respectively. The temperatures at the beginning and end of expansion are 1950°C and 870°C respectively. Determine the ideal efficiency of the cycle. Take $\gamma = 1.4$, if the compression ratio is 14 and the pressure at the beginning of the compression is 1 bar, calculate the maximum pressure in the cycle.

SOLUTION

Therefore,

Here,

$$T_{4} = 57 + 273 = 330 \text{ K}, T_{1} = 603 + 273 = 876 \text{ K}$$

$$T_{2} = 1950 + 273 = 2223 \text{ K}, T_{3} = 870 + 273 = 1143 \text{ K}.$$

$$r = \frac{\forall_{4}}{\forall_{1}} = 14, \ p_{4} = 1 \text{ bar}$$

$$\eta = 1 - \frac{1}{\gamma} \left[\frac{T_{3} - T_{4}}{T_{2} - T_{1}} \right]$$

$$= 1 - \frac{1}{\gamma} \left[\frac{1143 - 330}{2223 - 876} \right]$$

$$= 1 - 0.431 = 0.569$$

Maximum pressure in the cycle = p_1

For reversible adiabatic compression

$$p_1 \forall_1^{\gamma} = p_4 \forall_4^{\gamma}$$
$$p_1 = \left(\frac{\forall_4}{\forall_1}\right)^{\gamma} \cdot p_4 = (14)^{1.4} \times 1 = 40.23 \text{ bar}$$

So,

EXAMPLE 5

An air standard diesel cycle has a compression ratio of 18 and the heat transfer to the working fluid per cycle is 1800 kJ/kg. At the beginning of the compression process the pressure is 0.1 MPa and the temperature is 15°C. Determine (a) the pressure and temperature at each point in the cycle (b) thermal efficiency (c) effective pressure.





SOLUTION Here,

Now,

...

:..

:..

:..

Again,

$$p_1 = p_4 \times 57.2 = \frac{100 \times 57.2}{1000} = 5.72 \text{ MPa}$$

EXAMPLE 6

Gas is entering in a turbine @ 15 kg/s, while 1 kg gas occupies 0.45 m³. Power developed by the pump is 12,000 kW. Velocities and enthalpies at inlet and outlet section of the turbine are 50 m/s, 110 m/s, 1260 kJ/kg and 400 kJ/kg respectively. Find out the amount of heat rejected and also the cross-section at the inlet.

SOLUTION

Rate of flow of gases m = 15 kg/s,

 $v_{s} = 0.45 \text{ m}^{3}/\text{kg}$

Power developed by the turbine P = 12000 kW

:. Work done
$$w = \frac{12000}{15} = 800 \, \text{kW/kg}$$

Enthalpy of the gases at the inlet $h_1 = 1260 \text{ kJ/kg}$ Enthalpy of the gases at the outlet $h_2 = 400 \text{ kJ/kg}$ Velocity of the inlet $V_1 = 50$ m/sec Velocity at the outlet $V_2 = 110$ m/sec

(i) Heat rejected

$$h_1 + \frac{V_1^2}{2} + q = h_2 + \frac{V_2^2}{2} + w$$



or,	$1260 + \frac{50^2}{2 \times 10^3} + q = 400 + \frac{110^2}{2 \times 10^3} + 800$
or,	$1260 + \frac{2500}{2 \times 1000} + q = 400 + \frac{110 \times 110}{2 \times 1000} + 800$
or,	1260 + 1.25 + q = 400 + 6.05 + 800
<i>.</i>	q = -55.2 kJ/kg
So, heat reject	ted (<i>i</i>) 55.2 kJ/kg = 55.2×15 kJ/s. = 828 kW
Therefore,	Inlet area = $\frac{\text{volume flow rate}}{\text{velocity}} = \frac{0.45 \times 15}{50} = 0.135 \text{ m}^2$

EXAMPLE 7

Some amount of fluid is heated from 50°C to 450°C at constant 2 bar pressure. Initial and final specific volumes of the fluid are 1 m³/kg and 1.8 m³/kg. Determine (i) heat added to the fluid (ii) work done of the fluid (iii) change in internal energy (iv) change in enthalpy for a steady flow process. Assume $C_p = 2.5 + \frac{40}{T+20}$ SOLUTION

$$v_{s1} = 1 \text{ m}^3/\text{kg}, T_1 = 50^{\circ}\text{C}, v_{s2} = 1.8 \text{ m}^3/\text{kg}, T_2 = 450^{\circ}\text{C}$$

Specific heat at constant 2 bar pressure

$$C_p = 2.5 + \frac{40}{T + 20}$$

(i) Heat added per kg of the fluid is given by

$$q = \int_{T_1}^{T_2} C_p \cdot dT = \int_{50}^{450} \left(2.5 + \frac{40}{T + 20} \right) dT$$
$$= \left[2.5T + 40 \ln \left(T + 20 \right) \right]_{50}^{450}$$
$$= 2.5 \left(450 - 50 \right) + 40 \ln \left(\frac{450 + 20}{50 + 20} \right)$$
$$= 2.5 \times 400 + 40 \ln \left(\frac{470}{70} \right)$$
$$= 1076 \text{ kJ/kg}$$

(ii) Work done/kg of the fluid

$$w = \int_{v_{s1}}^{v_{s2}} p dv_s$$

= $p (v_{s2} - v_{s1})$
= $2 \times 10^5 (1.8 - 1) = 2 \times 10^5 \times 0.8 \text{ N.m/kg}$
= $160 \times 10^3 = 160 \text{ kJ/kg}.$

(iii) Change in internal energy

$$\Delta u = q - w = 1076 - 160 = 996 \text{ kJ/kg}$$

(iv) Change in enthalpy (for steady flow process)

 $\Delta h = q = 1076 \text{ kJ/kg}$

EXAMPLE 8

Calculate the entropy change of 10 gm of water at 20°C when it is converted to ice at -10°C. Assume the specific heat of water to remain constant at 4.2 J/gm K and that of ice to be half of the value and taking the latent heat of fusion of ice at 0°C to be 335 J/gm

SOLUTION

Entropy change of water from 20°C to 0°C

$$\Delta S_1 = \int_{273+20}^{273+0} \frac{mC_p \cdot dT}{T} = 10 \times 4.2 \left[\ln T \right]_{293}^{273} = -2.97 \text{ J/K}$$

Entropy change of water from 0°C to ice at 0°C

$$\Delta S_2 = \frac{mL}{T} = \frac{10 \times 335}{273 + 0} = -12.27 \text{ J/K}$$

Entropy change of ice at 0°C to ice at -10°C

$$\Delta S_3 = \int_{273+0}^{273-10} \frac{mC_p \cdot dT}{T} = 10 \times \frac{4.2}{2} \left[\ln T \right]_{273}^{263} = -0.783 \text{ J/K}$$

:. Net entropy change

$$\Delta S = \Delta S_1 + \Delta S_2 + \Delta S_3$$

= -2.97 - 12.27 - 0.783 = -16.023 J/K

EXAMPLE 9

A lump of steel of mass 10 kg at 627°C is dropped in 100 kg oil at 30°C. The specific heats of steel and oil are 0.5 kJ/kg K and 3.5 kJ/kg K respectively. Calculate the entropy change of the steel the oil and universe.

SOLUTION

Let the final temperature at equilibrium = TK

Here for thermal equilibrium heat lost by steel = Heat gained by oil

$$\therefore \qquad [mC_p \cdot (\Delta T)]_{\text{steel}} = [mC_p \cdot (\Delta T)]_{\text{oil}}$$

$$\therefore \qquad 10 \times 0.5 (627 + 273 - T) = 100 \times 3.5 \times [T - (273 + 30)]$$

 $T = 311.41 \,\mathrm{K}$

: Change of entropy of steel

$$(\Delta S)_{\text{steel}} = \int_{627+273}^{311.41} \frac{mC_p \cdot dT}{T}$$

$$= 10 \times 0.5 \ln\left(\frac{311.41}{900}\right) = -5.306 \text{ J/K}$$

: Change of entropy of oil

$$(\Delta S)_{\text{oil}} = \int_{30+273}^{311.41} \frac{mC_p \cdot dT}{T}$$
$$= 100 \times 3.5 \ln\left(\frac{311.41}{303}\right) = 9.582 \text{ J/K}$$

:. Change of entropy of the universe

$$(\Delta S)_{\text{universe}} = (\Delta S)_{\text{steel}} + (\Delta S)_{\text{oil}}$$
$$= -5.306 + 9.582 = 4.276 \text{ kJ/K}.$$

EXAMPLE 10

The compression ratio in an air standard Otto cycle is 8. At the beginning of the compression stroke the pressure is 0.1 MPa and the temperature is 15° C. The heat transfer to the air per cycle is 1800 kJ/kg. Determine (a) the pressure and temperature at the end of each process of the cycle (b) thermal efficiency (c) mean effective pressure

SOLUTION



Heat rejected in constant volume process (2-3)

 $Q_{2-3} = mC_v (T_2 - T_3)$

Heat absorbed in constant volume process (4–1)

$$Q_{4-1} = mC_{v}(T_{1} - T_{4})$$

$$p_{3} = 0.1 \text{ MPa} = 100 \text{ kPa}, T_{3} = (273 + 15) = 288\text{K}$$

$$\forall_{3} \quad \forall_{2} = 0 \text{ heat transfer} = 1800 \text{ LMa}$$

 $r = \frac{\nabla_3}{\nabla_4} = \frac{\nabla_2}{\nabla_1} = 8$, heat transfer = 1800 kJ/kg

Given

Now,

$$\frac{T_4}{T_5} = \frac{T_1}{T_2} = (r)^{\gamma-1} = (8)^{0.4} = 2.297$$

$$\therefore \qquad T_4 = T_3 \times 2.297 = 661.54 \text{ K} \qquad T_1 = 2.297 T_2 \qquad (i)$$
Again, $mC_v[(T_1 - T_a) - (T_2 - T_3)] = 1800$
or
 $C_v(2.297 T_2 - 661.54 - T_2 + 288) = 1800$ [For unit mass, is, $m = 1 \text{ kg}$]
or
 $C_v(1.197 T_2 - 373.54) = 1800$

$$\therefore \qquad T_2 = \left(\frac{1800}{C_v} + 373.54\right) \times \frac{1}{1.197} \qquad (ii)$$
Assuming, $C_\rho = 1 \text{ kJ/kgK}$, $C_v = \frac{C_\rho}{\gamma} = \frac{1}{1.4} = 0.714$
From equation (ii), $T_1 = 5554.54 \text{ K}$
Again, $\frac{p_4}{P_3} = (r)^{r} - (8)^{1.4} = 18.379 - \frac{p_1}{P_2}$ (iii)
 $\therefore \qquad p_4 = 100 \times 18.379 = 1837.9 \text{ kPa}$
For constant volume process,
 $\frac{P_2}{T_2} = \frac{P_3}{T_3}$

$$\therefore \qquad p_2 = 100 \times \frac{2418.17}{1.288} = 839.64 \text{ kPa}$$
From equation (ii), $p_1 = 839.64 \times 18.379 = 15431.74 \text{ kPa}$
Thermal efficiency $\eta = 1 - \frac{1}{(r)^{\gamma-1}} = 1 - \frac{1}{(8)^{0.4}} = 0.5647$
Again we can write $\forall_5 = 8\forall_4 = 8\forall_1 = \forall_2$
 $\therefore \qquad Work done = \frac{1}{\gamma - 1} [(p_1\forall_1 - p_2\forall_2) - (p_4\forall_4 - p_3\forall_3)]$
 $\therefore \qquad 1800 = \frac{1}{0.4} [15431.74\forall_1 - 839.64 \times 8\forall_1 - 1837.9\forall_1 + 100 \times 8\forall_1]$
 $\therefore \qquad \forall_1 = 0.0938 \text{ m}^3\text{kg}$
 $\therefore \qquad \text{Mean effective pressure} = \frac{Work done}{Stroke volume}$
 $= \frac{1800}{0.7504} = 2398.72 \text{ kPa}$

EXAMPLE 11

An air standard otto cycle operates with the compression and expansion strokes being polytropic with n = 1.35 for each stroke. Calculate the thermal efficiency of the cycle for a compression ratio of 9.

SOLUTION

Compression ratio (r) = 9, Polytropic constant (n) = 1.35 Efficiency of the otto cycle is

$$\eta_{\text{otto}} = 1 - \frac{1}{(r)^{\gamma - 1}}$$
$$= 1 - \frac{1}{q^{1.35 - 1}} = 0.5365$$

EXAMPLE 12

An air standard diesel cycle has a compression ratio of 22 and cut off ratio of 2.2. Determine the thermal efficiency of the cycle.

SOLUTION

Compression ratio (r) = 22, Cut-off ratio (ρ) = 2.2

So,

$$\eta_{\text{diesel}} = 1 - \frac{1}{(r)^{\gamma - 1}} \left[\frac{\rho^{\gamma} - 1}{\gamma(\rho - 1)} \right]$$
$$= 1 - \frac{1}{(22)^{1.4 - 1}} \left[\frac{(2.2)^{1.4} - 1}{1.4(2.2 - 1)} \right]$$
$$= 0.6515$$

- 1. What is a cycle? what is the difference between an ideal and actual cycle?
- 2. What is an air-standard efficiency?
- 3. What is relative efficiency?
- 4. Derive expressions of efficiency in the following cases:

(i) Carnot cycle; (ii) Diesel cycle; (iii) Dual combustion cycle.

- 5. Explain "Air standard analysis" which cycle has been adopted for I.C. engine cycle. State the assumption made for air standard cycle.
- 6. Efficiency of Otto cycle and Diesel cycle—which one is more?
- 7. A 6-cylinder petrol engine has volume compression ratio 5:1 and clearance volume in each cylinder is 110 cc. Engine consumes fuel 10 kg/hr. having calorific value 10,000 k cal/kg. Engine runs at 2400 rpm and efficiency ratio 0.66. Estimate mean effective pressure developed.

Ans. [7.038 kgf/cm²]

- 8. Carnot engine produces 150 kJ work between 650 K and 37°C. what is the thermal efficiency and heat added during the process?
- 9. Being operated between two reservoirs at T_1 and T_3 , output of Carnot engine is 0.6 times the heat rejected. Determine thermal efficiency, temperatures of source and sink, if source is 200°C higher than sink.

Ans. [37.5%, 260.3°C, 60.3°C]

10. Initial and final temperatures of isentropic compression in an Otto cycle are 316 K and 323°C, respectively. What is the compression ratio and air-standard efficiency? Take $\gamma = 1.4$.

Ans. [4.885, 47%]

11. Engine, working on Otto cycle has stroke 22.5 cm and 0.15 m diameter. Calculate air-standard efficiency of the engine, while $\gamma = 1.4$ and clearance volume = 1.25×10^{-3} m³.

Ans. [43.6%]

12. In a diesel engine, fuel cut-off is at 8% of the stroke. What is its air-standard efficiency, being compression ratio 13:1 and $\gamma(air) = 1.4$?

Ans. [58.3%]

 In a ideal diesel engine clearance volume is 10% of swept volume while stroke is 0.2 m and diameter 150 mm. If cut-off occurs at 6% of the stroke, calculate its air-standard efficiency and compression ratio.

Ans. [57.53%, 11]

14. A 4-cylinder 4-stroke petrol engine has bore 6.5 cm, stroke 95 mm, speed 3000 rpm, clearance volume 65 cc, relative efficiency 50%, calorific value of fuel used 46 MJ/kg. When tested it develops 70 Nm torque. Determine, specific fuel consumption, brake mean effective pressure of the engine.

Ans. [0.308925 kg/kWh, 6.976 bar]

15. A 4-stroke petrol engine having 6 cylinders is to operate with a compression ratio of 8, and delivers power of 300 kW at 2400 rpm. Find out engine bore and stroke length and fuel consumption rate while stroke is 1.6 times the bore, mechanical efficiency 88% indicated mean effective pressure 9.5 kgf/cm², relative efficiency 50% and calorific value of fuel 11000 kcal/kg.

Ans. [134.36 mm, 214.98 mm, 83.067 kg/hr.]

- 16. Following are the particulars of a 2-stroke diesel engine; bore = 10 cm stroke = 150 mm, piston speed = 300 m/min, torque developed = 58 Nm, mechanical efficiency = 0.80, indicated thermal efficiency = 0.40, calorific value of fuel = 44 MJ/kg. Determine (i) indicated power (ii) indicated mean effective pressure (iii) fuel consumption rate on brake power basis.
- 17. A 4-cylinder 2-stroke petrol engine develops 30 kW at 2500 rpm. Mean effective pressure of each is 800 kN/m² and mechanical efficiency 80%. Calculate diameter and stroke of each cylinder, of stroke to bore ratio is 1.5. Also calculate the brake specific fuel consumption of the engine, if brake thermal efficiency is 28% and calorific value of petrol is 44×10^6 J/kg.

Ans. [5.75 cm, 8.625 cm, 0.2922 kg/kwh]



STEAM POWER CYCLE

9.1 INTRODUCTION

We have already discussed the air standard cycle in the previous chapter. In this chapter, we shall discuss the steam power cycle, *i.e.*, thermodynamic cycle of Q vapour. In a vapour cycle, all the theory remains the same, except the working substance, which is steam. Though there are many vapour cycles, from the subject point of view, important vapour power cycles are (i) *Carnot vapour power cycle* which is already dicussed in chapter 5 (ii) *Rankine vapour cycle*.

9.2 RANKINE CYCLE

Rankine cycle is a thermodynamic cycle. It is *modified form of Carnot cycle*. It has practical application on thermal power plant (steam power plant). It consists of four units, *i.e.*, boiler, turbine, condenser and fuel pump. Schematic diagram of steam power plant or turbine plant is shown in figure 9.1.

Consider 1 kg of saturated water at pressure p_1 and absolute temperature T_1 as represented by point 1 in figure 9.2 (a), (b) and (c). The cycle is completed by the following four processes: *isothermal* expansion, *isentropic expansion*, *isothermal* compression and *isentropic compression*. Process: (1–2)

Reversible adiabatic expansion in the turbine. *Process:* (2–3)

Constant pressure transfer of heat in the condenser.







Fig. 9.2 (a)

Process: (3-4)

Reversible adiabatic pumping process in the fuel pump.



Process: (4-1)

Constant pressure transfer of heat in the boiler.

Figure 9.1 shows the Rankine cycle and figure 9.2, $p-\forall$, T-S and h-s diagram (when the saturated steam enters the turbine, the steam can be wet or superheated also). Consider 1 kg of fluid, applying steady flow energy equation (S.F.E.E) to boiler, turbine, condenser and pump.

(i) For boiler (as control volume), we get

$$h_{f4} + q_1 = h_1 \qquad \Rightarrow q_1 = h_1 - h_{f4}$$
(9.1)
(ii) For turbine (as control volume), we get $[w_r = \text{Turbine work}]$

$$h_1 = w_T + h_2 \qquad \Rightarrow w_T = h_1 - h_2 \tag{9.2}$$

(iii) For condenser, we get

$$h_2 = q_2 + h_{f3} \implies q_2 = h_2 - h_{f3}$$
 (9.3)

(iv) For the feed pump, we get $h_{f3} + w_p = h_{f4}$

$$\Rightarrow w_p = h_{f4} - h_{f3} \tag{9.4}$$

[w = Pump work]

Efficiency of the Rankine cycle *:*.

$$\eta = \frac{w_{net}}{q_1} = \frac{w_T - w_p}{q_1}$$
$$= \frac{(h_1 - h_2) - (h_{f4} - h_{f3})}{q_1}$$
(9.5)

Using general property relation for reversible adiabatic compression,

$$T \cdot ds = dh - v dp$$
 and $ds = 0$

dh = vdp*.*..

$$\Rightarrow \qquad \Delta h = v_3 \cdot \Delta p$$

 $\Delta h = v_3 \cdot \Delta p$ $h_{f4} - h_{f3} = v_3(p_1 - p_2)$ \Rightarrow

As $(h_{f4} - h_{f3})$ is very small in comparison to w_T and can be neglected, especially when the boiler pressure is low.

$$\eta_{\text{Rankine}} = \frac{h_1 - h_2}{h_1 - h_{f4}}$$

9.3 VAPOUR COMPRESSION REFRIGERATION CYCLE

The refrigeration cycles are classified as vapour compression refrigeration cycle and gas refrigeration cycle based on the working medium employed in the refrigerator. The basic components involved in a vapour compression refrigeration plant are shown in schematic diagram of vapour compression refrigeration plant in figure 9.3. There are four basic units, compressor, condenser, throttle valve and evaporator. The $p-\forall$, T-S, h-s and schematic diagram is shown in figures 9.4 and 9.5.



Fig. 9.3

Fig. 9.4

The vapor compression refrigeration cycle consists of following processes.



251

Process 1–2: The working fluid at state 1 (a mixture of liquid and vapour) is isentropically compressed to state 2 (saturated vapour) inside the compressor.

Process 2–3: Isothermal energy rejection as heat (Q_{μ}) to the surroundings at T_{μ} in a condenser. At the end of the process, the working fluid is available at state 3(saturated liquid).

Process 3–4: In throttle valve, the working fluid is expanded isentropically to a low pressure. The temperature of the working fluid is T_{i} at the end of the expansion process.

Process 4–1: Isothermal energy transfer as heat (Q_L) from the low temperature body at T_L to the working fluid. If the flow rate of the refrigerant or working through the device is *m* then, energy extracted from the cold body

$$Q_{L} = m (h_{1} - h_{4}) \tag{9.6}$$

Work done on the compressor

$$W = m (h_2 - h_1)$$
(9.7)

$$COP = \frac{Q_L}{W} = \frac{m(h_1 - h_4)}{m(h_2 - h_1)} = \frac{h_1 - h_4}{h_2 - h_1}$$
(9.8)

Since process 3–4 is a throttling process, we have $h_4 = h_3$

Multiple Choice Questions

- The Rankine Cycle, as compared to Carnot Cycle, has work ratio

 (a) high
 (b) low
- 2. The ideal cycle on which a steam engine works is
 - (a) Carnot Cycle (b) Rankine Cycle
 - (c) Otto Cycle (d) Joule Cycle
- 3. In a Rankine Cycle with superheated steam
 - (a) the work done increases
 - (b) the specific steam consumption decreases
 - (c) the dryness fraction of steam after isentropic expansion increases
 - (d) all of the above
- 4. Rankine Cycle comprises
 - (a) two isothermal and two isentropics
 - (c) two isobarics and two isentropics
- 5. Which one is a thermodynamic cycle?
 - (a) Diesel Cycle
 - (c) Otto Cycle

- (b) two isobarics and two isothermals
- (d) two isothermals and two isochorics
- (b) Rankine Cycle
- (d) Joule Cycle
- 6. The highest temperature during the cycle, in a vapour compression system, occurs after
 - (a) compression (b)
 - (c) expansion
- (b) condensation(d) evaporation

...

7.	7. In a vapour compression system, the lowest temperature during the cycle occurs after												
	(a)	con	pression	1			(b)	condensatio	n				
	(c)	exp	ansion				(d)	evaporation					
8.	Sub)-coo	ling in a	refri	gerat	or cycle							
	(a)	doe	s not alte	er CC	OP	(b)	increases (СОР	(c)	deo	creases C	OP	
							Answers						
1. (a	I)	2.	(b)	3.	(d)	4. (c	:) 5.	(b) 6.	(a)	7.	(d)	8.	(c)
	NUMERICAL EXAMPLES												

EXAMPLE 1

In a steam power cycle, the steam supply is at 15 bar and dry and saturated. The condenser pressure is 0.4 bar. Calculate the Carnot and Rankine efficiency of the cycle. Neglect pump work.

SOLUTION

Steam supply pressure $(p_1) = 15$ bar, $x_1 = 1$, and condensor pressure $(p_2) = 0.4$ bar

From steam tables, at 15 bar

$$t_s = 198.3^{\circ}\text{C}, h_g = 2789.9 \text{ kJ/kg}, s_g = 6.4406 \text{ kJ/kg K}$$

So,

$$T_1 = 198.3 + 273 = 471.3 \,\mathrm{K}$$

From steam tables, at 0.4 bar

$$t_s = 75.9^{\circ}\text{C}, h_f = 317.7 \text{ kJ/kg}, h_{fg} = 2319.2 \text{ kJ/kg}$$

 $s_f = 1.0261 \text{ kJ/kg K}, s_{fg} = 6.6448 \text{ kJ/kg K}$
 $T_2 = 75.9 + 273 = 348.9 \text{ K}$

So,

$$\eta_{\text{Carnot}} = \frac{T_1 - T_2}{T_1} = \frac{471.3 - 348.9}{471.3} = 0.259$$

$$\eta_{\text{Rankine}} = \frac{\text{Adiabatic or isentropic heat drop}}{\text{Heat supplied}} = \frac{h_1 - h_2}{h_1 - h_f}$$

But,

But,
$$h_2 = h_{f2} + x_2 h_{fg2} = 317.7 + x_2 \times 2319.2$$

As steam expands isentropically, $s_1 = s_2$

$$\therefore \qquad 6.6448 = 1.0261 + x_2 \times 6.6448$$

$$\therefore x_2 = 0.845$$

:.
$$h_2 = 317.7 + 0.845 \times 2319.2 = 2277.4 \text{ kJ/kg}$$

$$\therefore \qquad \eta_{\text{Rankine}} = \frac{2789.9 - 2277.4}{2789.9 - 317.7} = 0.2073.$$

EXAMPLE 2

In a Rankine cycle, the steam at inlet to turbine is saturated at a pressure of 35 bar and the exhaust pressure is 0.2 bar. Determine (i) pump work (ii) turbine work (iii) Rankine efficiency (iv) condensor heat flow (v) dryness at the end of expansion. Assume flow rate of 9.5 kg/s.

SOLUTION

Pressure and condition of steam at inlet to the turbine $p_1 = 35$ bar, x = 1, exhaust pressure $p_2 = 0.2$ bar, flow rate m = 9.5 kg/sec From the steam tables, at 35 bar $h_1 = 2802 \text{ kJ/kg}, s_{g1} = 6.1228 \text{ kJ/kg K}$ at 0.20 bar, $h_{f3} = h_f = 251.5 \text{ kJ/kg}, h_{fg} = 2358.4 \text{ kJ/kg}$ $v_f = 0.001017 \text{ m}^3/\text{kg}, s_f = 0.8321 \text{ kJ/kgK}$ $s_{fg} = 7.0773 \text{ kJ/kg K}$

(i) The pump work

$$= (p_5 - p_3) v_f$$

= [(35 - 0.2) 10⁵ × 0.001017] × 10⁻³ kJ/kg
= 3.54 kJ/kg

also

 $h_{f4} - h_{f3} = 3.54 \text{ kJ/kg}$ $h_{f4} = 3.54 + 251.5 = 255.04 \text{ kJ/kg}$ *:*.. Power required to drive the pump = $9.5 \times 3.54 = 33.63$ kW

(ii) Turbine work

$$s_1 = s_2 = s_{f2} + x_2 s_{fg2}$$

$$\therefore \qquad 6.1228 = 0.8321 + x_2 \times 7.0773$$

:..

:..

$$x_{2} = \frac{6.1228 - 0.8321}{7.0773} = 0.747$$
$$h_{2} = h_{f2} + x_{2}hf_{g2}$$

Turbine work =
$$m(h_1 - h_2) = 9.5 (2802 - 2013) = 7495.5 \text{ kW}$$

[Pump work is very small in comparison to turbine work] (iii) Rankine efficiency

$$\eta_{\text{Ran}} = \frac{h_1 - h_2}{h_1 - h_{f2}} = \frac{2802 - 2013}{2802 - 251.5} = 0.3093$$

(iv) Condensor heat flow

$$= m (h_2 - h_{f_2})$$

= 9.5 (2013 - 251.5) = 16734.25 kW

(v) The dryness at the end of expansion

$$x_2 = 0.747$$



Steam Power Cycle

EXAMPLE 3

Determine the amount of heat which should be supplied to 2 kg of water at 25° C to convert it into steam at 5 bar and 0.9 dry. Consider specific heat of water = 4.18 kJ/kgK.

SOLUTION

Mass of water to be converted into steam $m_w = 2$ kg, temperature of water $t_w = 25^{\circ}$ C, pressure of steam 5 bar, dryness fraction = 0.9

From steam table, at 5 bar,

 $h_f = 640.1 \text{ kJ/kg}, h_{fg} = 2107.4 \text{ kJ/kg}$

:. Enthalpy of 1 kg of steam (above 0° C)

$$h = h_f + x h_{fg}$$

= 640.1 + 0.9 × 2107.4
= 2536.76 kJ/kg

:. Sensible heat associated with 1 kg of water

$$= m_{w} \times C_{pw} (t_{w} - 0)$$

= 1 × 4.18 (25 - 0) = 104.5 kJ

: Net quantity of heat to be supplied per kg of water

 $= 2536.76 - 104.5 = 2432.26 \, \text{kJ}$

 \therefore Total heat to be supplied =2×2432.26=4864.52 kJ

EXAMPLE 4

What amount of heat would be required to produce 4.4 kg of steam at pressure of 6 bar and temperature of 250°C from water at 30°C. Take specific heat for superheated steam as 2.2 kJ/kg K.

SOLUTION

Here,

 $m = 4.4 \text{ kg}, p = 6 \text{ bar}, t_{sup} = 250^{\circ}\text{C}$ = 30°C, $C_{ps} = 2.2 \text{ kJ/kgK}.$

From steam table at 6 bar,

temperature of water

$$t_s = 158.8^{\circ}\text{C}, h_f = 670.4 \text{ kJ/kg}, h_{f\sigma} = 2085 \text{ kJ/kg}$$

Enthalpy of 1 kg superheated steam recorded from 0°C

$$h_{sup} = h_f + h_{fg} + C_{ps} [t_{sup} - t_s]$$

= 670.4 + 2085 + 2.2 (250 - 158.8)
= 2956 kJ

Heat already with 1 kg of water

 $= 1 \times 4.18 (30 - 0) = 125.4 \text{ kJ}$

Net amount of heat required to be supplied/kg

= 2956 - 125.4 = 2830.6 kJ/kgTotal Heat required $= 4.4 \times 2830.6 = 12454.6 \text{ kJ}$

EXAMPLE 5

If a certain amount of steam is produced at a pressure of 8 bar and dry = 0.8. Calculate (i) external work done during evaporation (ii) internal latent heat of steam.

 $h_{fg} = 2046.5 \text{ kJ/kg}$

x = 0.8

SOLUTION

Here, p = 8 bar and From steam tables at 8 bar $v_g = 0.240 \text{ m}^3/\text{kg}$

(i) External work during evaporation

=

_

$$= p \cdot x \cdot v_g = 8 \times 10^5 \times 0.8 \times 0.24$$
 N-m

$$\frac{8 \times 10^5 \times 0.8 \times 0.24}{1000} \text{ kJ} = 153.6 \text{ kJ}$$

(ii) Internal latent heat

$$xh_{fg}$$
 - external work
= 0.8 × 2046.5 - 153.6 = 1483.6 kJ

EXAMPLE 6

Find the specific volume enthalpy and internal energy of wet steam at 18 bar. Dryness fraction = 0.85

SOLUTION

p = 18 barx = 0.85Here, and From steam tables at p = 18 bar $h_f = 884.6 \, \text{kJ/kg}$ $h_{fg} = 1910.3 \text{ kJ/kg}$ $h_f = 884.5 \text{ kJ/kg}$ $v_g = 0.110 \,\mathrm{m^{3}/kg}$ $h_{\sigma} = 2794.8 \, \text{kJ/kg}$ Specific volume of wet steam (i) $v = x v_g = 0.85 \times 0.110 = 0.0935 \text{ m}^3/\text{kg}$ (ii) Specific enthalpy of wet steam $h = h_f + x h_{fg}$ $= 884.5 + 0.85 \times 1910.3 = 2508.25 \text{ kJ/kg}$ (iii) Specific internal energy of wet steam $w = (1-x)h_f + xh_g = (1-0.85)884.5 + 0.85 \times 2794.8$ $= 2508.25 \, \text{kJ/kg}$

EXAMPLE 7

10 kg of water at 45°C is heated at constant pressure of 10 bar until it becomes superheated vapour at 300°C. Find the change in volume, enthalpy, internal energy and entropy. Consider specific heats for superheated steam and water are 2.25 kJ/kgK and 4.187 kJ/kgK respectively.

Solution
At
$$p = 10$$
 bar = 1000 kPa
 $t_s = 273 + 179.9 = 452.9$ K, $t_{sup} = 300 + 273 = 573$ K
 $v_g = 0.194$, $C_{ps} = 4.187$, $C_{ps} = 2.25$, $h_{gs} = 2013.6$ kJ/kg, $h_g = 2776.1$ kJ/kg
Now, $v_{sup} = \frac{v_g \times t_{sup}}{t_s} = \frac{0.194 \times 573}{452.9}$
 $= 0.245$ m³/kg
At 45°C, $v_f = 0.00101$ m³/kg
Change in volume
 $= m(v_{sup} - v_f)$
 $= 10(0.245 - 0.00101) = 2.4399$ m³ = 2.44 m³
Change in enthalpy
 $\therefore \qquad \Delta H = m[C_{ps}(t_s - t_f) + h_{gs} + C_{ps}(t_{sup} - t_f)]$
 $= 10[4.187 (179.9 - 45) + 2013.6 + 2.25 (300 - 179.9)]$
 $= 28486.5$ kJ
 $h_{sup} = h_g + C_{ps} (300 - t_f)$
 $= 2776.1 + 2.25 (300 - 179.9) = 3046.3$ kJ/kg
So, $h_1 = h_{sup} - \frac{\Delta H}{m}$
 $= 3046.3 - \frac{28486.5}{10} = 3046.3 - 2848.65$
 $= 197.65$ kJ/kg
Change in internal energy
 $u_1 = h_1 - p_1v_1 = 197.65 - 1000 (0.00101) = 196.64$ kJ/kg
 $u_{sup} = h_{sup} - p_{sup} \times v_{sup}$
 $= 3046.3 - 1000 \times 0.245 = 2801.3$ kJ/kg
 $\therefore \qquad \Delta U = m(u_{sup} - u_f) = 10(2801.3 - 196.64) = 26046.6$ kJ
Change in entropy

$$\Delta S = m \left[C_{pw} \int_{273+45}^{273+176.9} \frac{dT}{T} + S_{fg} + C_{ps} \int_{273+179.9}^{273+300} \frac{dT}{T} \right]$$
$$= 10 \left[4.187 \ln \left(\frac{452.9}{318} \right) + 4.45 + 2.25 \ln \left(\frac{573}{452.9} \right) \right]$$
$$= 64.59 \text{ kJ/K}$$

EXERCISE

- 1. Draw the Carnot vapour power cycle on $p-\forall$, T-S and h-S diagrams. Establish that its efficiency depends on temperature limits only.
- 2. Explain the working of a steam power plant with the help of Rankine cycle.
- 3. What are the limitations for which the Carnot cycle is not used in steam power plants?
- 4. Show Rankine cycle on $p \forall$ and T S diagrams and explain the processes involved.
- 5. Draw the Rankine cycle on *T*–*S* diagram using dry saturated steam and obtain an expression for the Rankine cycle efficiency.
- 6. Draw the schematic diagram of steam power plant.
- 7. Draw layout of a vapour compression refrigerating system and also show corresponding pr.-vol. and temp.-enthalpy diagram. State functions of each components.

BIBLIOGRAPHY

THERMODYNAMICS

- 1. Arora, C.P., Thermodynamics, Tata Mc Graw-Hill Publishing Company Limited, New Delhi, 2001.
- 2. Ballaney, P.L., Thermal Engineering, Khanna Publishers, Delhi, 2003.
- 3. Khurmi, R.S. and Gupta, J.K., A Text book of Thermal Engineering, S. Chand & Company Ltd., 1996.
- 4. Rao, Y.V.C., Engineering Thermodynamics, University Press (India) Private Limited, Hyderabad, 2003.
- 5. Rajput, R.K., Thermal Engineering (6th ed.), Laxmi Publications Pvt. Ltd., New Delhi, 2006.

This page intentionally left blank

Group–B Fluid Mechanics

This page intentionally left blank



INTRODUCTION AND FUNDAMENTAL CONCEPTS

1.1 DEFINITION OF FLUID

To speak very loosely, fluid is a state of matter which flows. From the view point of Fluid Mechanics, all matters are segregated into two states, one is solid and the other is fluid. Hence, both liquid and gas are termed as fluid. Fluid is so defined as *one which under continuously applied tangential stress, starts deforming continuously as long as shear stress is applied*. By nature, any fluid may be Real or Ideal, Newtonian or Non-Newtonian, Compressible or Incompressible. It is of special mention that no fluid is truly ideal, but is nearly ideal.

1.2 FLUID MECHANICS AND ITS PERVIEW

Fluid mechanics is that branch of mechanics, which studies behavioural aspects of fluids, either in motion or at rest, with subsequent effects of boundaries, which may be either solid or fluid. Fluid statics deals with fluid at rest. But fluid dynamics concerns about the fluid at motion.

From our day-to-day activities to highly specialised phenomenon may be included within the perview of fluid mechanics. From breathing to blood flow, floating of ships to swimming against current, motion of submarines to flight of Concord aircrafts, water distribution network to turbine action in hydel plant, design of missile to rotation of windmill are a few glimpses where fluid mechanics has direct involvement.

1.3 FLUID AS A CONTINUUM

Continuum is an idealisation of continuous description of matter and properties of matter considered as continuous function of space variables. In molecular concept there exists intermolecular space, but in continuum approach, actual conglomeration of separate molecules occur, leaving no empty space within the matter or system. Schematically it is shown in figure 1.1.



In liquids, the cohesive force amongst all molecules are strong enough, so that the entire mass of liquid can behave as a continuous mass of substance. So, liquids can be analysed through continuum approach. This is true for solid also.

But in case of gases, the story is somewhat different. Mean free path of a gas is an average distance of molecules between two successive collisions. If this mean free path of a gas is less than 1/100 th part of a characteristic length, that gas can be considered continuous and can be analysed through continuum approach. Deviation of any gas from this minimum value will cease it to form continuum.

1.4 PROPERTIES OF FLUID

There exist a few characteristics of a fluid in continuum, independent of motion of fluid. These characteristics are termed as basic properties of a fluid. Such a few properties are discussed hereinafter.

1.4.1 Density

The density of fluid is defined as mass per unit volume. It is symbolised by ' ρ ' (*rho*). The density at a point within fluid is expressed as

$$\rho = \frac{Lt}{\delta \forall \to 0} \quad \frac{\delta m}{\delta \forall} = \frac{dm}{d \forall} \tag{1.1}$$

where δm and $\delta \forall$ represent elemental mass and elemental volume, respectively. Otherwise it can be said as specific mass. Sometimes the term 'mass density' is also used. The dimensional expression of density is ML^{-3} .

1.4.2 Specific Weight

The specific weight is the weight of fluid per unit volume, and is symbolised by ' γ ' (gamma). It can be expressed as,

$$\gamma = \frac{dw}{d\forall} = \frac{dm \times g}{d\forall} = \left(\frac{dm}{d\forall}\right) \times g = \rho \times g \tag{1.2}$$

where 'g' is the acceleration due to gravity. Dimensionally, specific weight is $ML^{-2}T^{-2}$. Weight density is another name of specific weight.

1.4.3 Specific Volume

The specific volume is the volume of the fluid per unit mass. It is the reciprocal of mass density. So,

$$v_s = \frac{d\forall}{dm} = \frac{1}{\rho} \tag{1.3}$$

Dimensionally specific volume is $M^{-1}L^3$.

264

1.4.4 Specific Gravity

The specific gravity of a liquid (S) is the ratio of its density or specific weight to that of water under standard temperature and pressure, (standard pressure = 1atm. and standard temperature = 4° C). So,

$$S = \frac{\gamma_{\text{liq}}}{\gamma_{\text{water/4°C}}} = \frac{\rho_{\text{liq}}}{\rho_{\text{water/4°C}}}$$
(1.4)

Being a ratio, specific gravity has no unit and is a pure number.

1.4.5 Viscosity and Newton's Law of Viscosity

It is a property of the fluid, the effect of which is better understood when the fluid is in motion. In a fluid flow, while different fluid elements move with different velocities, each element will feel some resistance due to friction. Now it begins to move at a strain rate proportional to shearing stress. Referring to figure 1.2(a), showing shearing effect of two adjacent layers, displacement occurred in elemental time ' δt ' is ' $\delta u \cdot \delta t$ ', subtending shearing strain ' $\delta \theta$ '. So,



$$\propto \frac{\delta \theta}{\delta t} \tag{1.5a}$$

From geometry, we can say, $\tan \delta \theta = \frac{\delta u \cdot \delta t}{\delta y}$

Considering ' $\delta\theta$ ' be very small and within limits, tan $d\theta \approx d\theta$

 $d\Theta = \frac{du \cdot dt}{dy}$ $\frac{d\Theta}{dt} = \frac{du}{dy}$

τ

Following equation (1.5a), within limits and introducing constant of proportionality, we obtain

$$\tau = \mu \, \frac{d\theta}{dt} = \mu \frac{du}{dy} \tag{1.5c}$$

In fluid mechanics our concentration focuses on velocity distribution [figure 1.2 (b)] and hence,

$$\tau = \mu \, \frac{du}{dy} \tag{1.5d}$$

(1.5b)

This relationship is famous as *Newton's law of viscosity*, after *Sir Isaac Newton*, who proposed this law in 1687. This constant of proportionality ' μ ' (*mu*) is termed as *coefficient of viscosity* or *coefficient of dynamic* or *kinetic viscosity*. Dimensionally, it is expressed as $ML^{-1}T^{-1}$.

Another type of viscosity exists which is kinematic viscosity. Coefficient of kinematic viscosity is the ratio of kinetic viscosity and mass density. It is expressed as 'v'(nu).

$$v = \frac{\mu}{\rho} \tag{1.5e}$$

Dimensionally, kinematic viscosity is L^2T^{-1} .

Two factors are generally responsible for generation of viscosity:

- 1. Intermolecular force of cohesion.
- 2. Exchange of molecular momentum.

1.4.5.1 Newtonian and Non-Newtonian Fluid

In this article, a term has an obvious appearance that is *Rheology*.

As per the strict definition, rheology is concerned with the description of the flow behaviour of all types of matter. By convention, however, the main interests of rheology are restricted to industrially relevant materials with properties intermediate between those of ideal solids and liquids. So, an useful engineering definition of rheology is the description of materials using '*constitutive equation*' between the stress history and strain history.

The fluid which obeys Newton's law of viscosity is called Newtonian Fluid, *e.g.*, air and water.

There are certain fluids, where the linear relationship between shear stress (γ) and deformation

rate $\left(\frac{d\theta}{dt}\right)$ or velocity gradient $\left(\frac{du}{dy}\right)$ does not exist. And these are specified as Non-Newtonian

Fluid. It is also of various types, which is shown in this tree diagram.



A generalised Power-law model can describe both Newtonian and Non-Newtonian fluid as following:

$$\tau = A \left(\frac{du}{dy}\right)^n + B \tag{1.6}$$

- (i) When n = 1, B = 0, it is Newtonian fluid.
- (ii) When n > 1, the fluid is dilatant, i.e., quicksand, butter, printing ink.
- (iii) When n < 1, the fluid is pseudoplastic, i.e., gelatine, blood, milk, paper pulp, etc.
- (iv) When $n = 1, B \neq 0$, the fluid is Bingham plastic, i.e., sewage sludge, drilling muds, etc.

Thixotropic fluid is a special type of Non-Newtonian fluid, which shows an increase in apparent viscosity with time. Lipstick and certain paints show thixotropic behaviour.

In reverse, the Non-Newtonian fluids showing decrease in apparent viscosity with time are Rheopectic fluid, e.g, bentonite solution, gypsum suspensions in water etc. Illustration of various type of fluids is shown in figure 1.3.

1.4.5.2 Real Fluid and Ideal Fluid

All the fluid which we encounter in our day to day activities may have some viscosity. Some of them have the property of compressibility, *i.e.*, under pressure the volume gets reduced. These are the real fluids. In order to achieve simplifications



in various theories of fluid, sometimes it is assumed to be inviscid or non-viscous and incompressible. And any inviscid, incompressible fluid is an ideal fluid. But in reality, there exists no ideal fluid.

1.4.6 Bulk Modulus of Elasticity and Compressibility

Bulk modulus of elasticity (E) is defined as the ratio of volumetric stress to volumetric strain. So,

$$E = \frac{Lt}{\delta \forall \to 0} \frac{-\delta p}{\left(\frac{\delta \forall}{\forall}\right)}$$
$$= -\frac{dp}{\left(\frac{d \forall}{\forall}\right)}$$

We can express mass of substance as $m = \rho \forall$ Taking differential of above equation, we obtain $dm = \rho d \forall + \forall d \rho$ Mass being an invariant parameter, dm = 0 $\rho d \forall + \forall d \rho = 0$ So, $-\frac{d\forall}{\forall} = \frac{dp}{\rho}$ Therefore,

Substituting in the expression of bulk modulus of elasticity, we obtain

Ì

$$E = \frac{dp}{d\rho/\rho} = \rho \frac{dp}{d\rho}$$
(1.7)

The bulk modulus of elasticity is dimensionally expressed as $ML^{-1}T^{-2}$ and that values for water and air at atmospheric pressure are approximately 2000 MPa and 0.101 MPa.

Compressibility (K) or coefficient of compressibility is defined as reciprocal of bulk modulus of elasticity. So,

$$K = \frac{1}{E} = \frac{1}{\rho} \frac{d\rho}{dp}$$
(1.8)

The property of compressibility is usually referred in the context of gas. Dimensionally, it is $M^{-1}LT^{-2}$.

For an isothermal process,
$$\frac{p}{\rho} = \text{Constant}$$

 $\therefore \qquad \frac{dp}{d\rho} = \text{Constant} = \frac{p}{\rho}$
 $\therefore \qquad \text{Bulk modulus}(E) = \rho \cdot \frac{dp}{d\rho} = p$ (1.8a)
For an isentropic process, $\frac{p}{\rho^{\gamma}} = \text{Constant} = C$
 $\therefore \qquad dp = C \cdot \gamma \cdot \rho^{\gamma - 1} \cdot d\rho$
 $\Rightarrow \qquad \frac{dp}{d\rho} = C \cdot \gamma \cdot \rho^{\gamma - 1}$

 \Rightarrow

$$= \left(\frac{p}{\rho^{\gamma}}\right) \cdot \gamma \cdot \rho^{\gamma - 1} = \gamma \left(\frac{p}{\rho}\right)$$

$$\therefore \qquad \text{Bluk modulus } (E) = \rho \cdot \frac{dp}{d\rho} = \gamma \cdot p \tag{1.8b}$$

Mach Number (Ma) is a very useful dimensionless parameter for the analysis of flow, especially compressible. It is the ratio of velocity of fluid flow to the velocity of sound in the flowing medium at that condition. So,

$$Ma = \frac{V}{a}$$
(1.9)

The Mach Number is named after Ernst Mach (1838–1916), an Austrian physicist. The value of Mach number 0.30 is a very significant one, below which a flow normally is considered to be incompressible. Variety of fluid flow according to Mach Number are:

0.3< Ma < 0.8	: Subsonic Flow
0.8< Ma < 1.0	: Subsonic-Transonic Flow
1.0< Ma < 1.2	: Supersonic-Transonic Flow
1.2< Ma < 3.0	: Supersonic Flow
3.0< Ma	: Hypersonic Flow

1.4.7 Surface Tension of Liquid

Referring to figure 1.4 (a), the molecule 'p' with diameter '2a' experiences equal attraction from surrounding molecules at all direction. But the molecule 'q' on the surface experiences a resultant inward pull due to unbalanced cohesive force of the molecules. The horizontal components of cohesive force of the molecules keep a fluid particle on the surface under tension and this tensile force acting normal to a unit length on the surface is called surface tension σ (sigma).

Figures 1.4 (*b*) and (*c*) show a small differential membrane area of the interface with radii of curvature of R_1 and R_2 , If the pressure difference be Δp , then from equilibrium in vertical direction gives,







 $\Delta p \cdot ds_1 \cdot ds_2 = 2(\sigma \cdot ds_2) \sin \frac{d\alpha_1}{2} + 2(\sigma \cdot ds_1) \sin \frac{d\alpha_2}{2}$

$$= \sigma \cdot \left(ds_2 \cdot d\alpha_1 + ds_1 \cdot d\alpha_2 \right)$$

$$= \sigma \cdot \left(ds_2 \cdot \frac{ds_1}{R_1} + ds_1 \cdot \frac{ds_2}{R_2} \right)$$

$$= \left(ds_1 \cdot ds_2 \right) \cdot \sigma \cdot \left(\frac{1}{R_1} + \frac{1}{R_2} \right)$$

$$\Rightarrow \qquad \Delta p = \sigma \cdot \left(\frac{1}{R_1} + \frac{1}{R_2} \right)$$

$$\sigma = \frac{\Delta p}{\left(\frac{1}{R_1} + \frac{1}{R_2} \right)} \qquad (1.10)$$

The dimensional formula for surface tension is MT^{-2} , as ' σ ' is considered as force per unit length.

1.4.8 Capillarity

It is a phenomenon by virtue of which it is understood that if a liquid either rises or falls in a tube dipped into the liquid. This rise or fall depends upon whether force of adhesion is more than that of cohesion or not. In the adjoining figure, considering vertical equilibrium, we have capillary rise or depression 'h' as

$$\frac{\pi d^2}{4} h\rho g = \sigma \pi D \cos \theta$$

$$h = \frac{4\sigma \cos \theta}{\rho g D}$$
(1.11)





For pure water in air in clean glass tube, $\theta = 0^{\circ}$ and in case of mercury it is capillary depression and $\theta = \text{nearly } 130^{\circ}$ (Fig. 1.5).

Daramotor	Dimonsion	Units				
1 al ameter	Dimension	C.G.S	S.I.	Commonly Used		
Density	ML ⁻³	gm/cc	kg/m ³	kg/m ³		
Specific Weight	$ML^{-2}T^{-2}$	dyne/cc	N/m ³	N/m ³		
Specific Volume	$M^{-1}L^{3}$	M ⁻¹ L ³ cc/gm m ³ /kg		m³/kg		
Specific Gravity	PURE NUMBER					
Dynamic Viscosity or Kinetic Viscosity	$ML^{-1}T^{-1}$	dyne.s /cm ² = poise	Ns/m ² = Pas	poise or centipoise		
Kinematic Viscosity	L^2T^{-1}	cm ² /s = stoke	m²/s	stoke or centistoke		
Bulk Modulus	ML ⁻¹ T ⁻² dyne/cm ²		N/m ² = Pascal (Pa)	$ \frac{1 \text{ MPa}}{= 10^6 \text{ N/m}^2} = 1 \text{ N/mm}^2 $		
Compressibility	$M^{-1}LT^2$	cm ² /dyne	m²/N			
Surface Tension	M T ⁻²	dyne/cm	N/m	N/m		

Parameters Dimensions and Unit

Multiple Choice Questions

- 1. An ideal fluid is
 - (a) very viscous
 - (c) frictionless and incompressible
- (b) one which obeys Newton's law of viscosity
- (d) none of these.
- 2. A fluid is a substance that
 - (a) always expands until it fills any containers
 - (b) is practically incompressible
 - (c) cannot withstand any shear force
 - (d) cannot remain at rest under the action of any shear force
 - (e) obeys Newton's law of viscosity
 - (f) none of the above.
- 3. Shear stress in a fluid flowing between two parallel plates
 - (a) is constant over the c/s
 - (b) is zero at the plate and increase linearly to the midpoint
 - (c) varies parabolically across the section
 - (d) is zero at the midpoint and varies linearly with the distance from the midpoint.
| 4. | The | he property of fluid by virtue of which it offers resistance to shear is called | | | | | | | | |
|-----|------------|---|--------|---|--|--|--|--|--|--|
| | (a) | surface tension | | | | | | | | |
| | (b) | adhesion | | | | | | | | |
| | (c) | cohesion | | | | | | | | |
| | (d) | viscosity | | | | | | | | |
| | (e) | all of the above. | | | | | | | | |
| 5. | Nev | wton's law of viscosity states that shear sta | ress i | s directly proportional to | | | | | | |
| | (a) | velocity | (b) | velocity gradient | | | | | | |
| | (c) | shear strain | (d) | viscosity | | | | | | |
| 6. | Nev | wton's law of viscosity relates | | | | | | | | |
| | (a) | pressure, velocity and viscosity | | | | | | | | |
| | (b) | shear stress and rate of angular deformati | on in | a fluid | | | | | | |
| | (c) | shear stress, temperature, viscosity and v | eloci | ity | | | | | | |
| | (d) | pressure, viscosity and rate of angular de | form | ation | | | | | | |
| | (e) | none of the above. | | | | | | | | |
| 7. | Sto | ke is the unit of | | | | | | | | |
| | (a) | surface tension | (b) | viscosity | | | | | | |
| | (c) | kinematic viscosity | (d) | none of the above | | | | | | |
| 8. | Poi | se is the unit of | | | | | | | | |
| | (a) | mass density | (b) | kinematic viscosity | | | | | | |
| | (c) | velocity gradient | (d) | dynamic viscosity | | | | | | |
| 9. | Spe | ecific volume has the dimension of | | | | | | | | |
| | (a) | $M^{-1}L^3$ | (b) | ML ⁻³ | | | | | | |
| | (c) | MLT ⁻¹ | (d) | $M^{-1}L^{-1}T.$ | | | | | | |
| 10 | The | e S.I. unit of specific gravity is | | | | | | | | |
| | (a) | N/m ³ | (b) | Pascal | | | | | | |
| | (c) | dyne | (d) | none of the above. | | | | | | |
| 11. | The | e ratio of velocity of fluid and that of air is | defin | ed by | | | | | | |
| | (a) | Froude's Number | (b) | Nusselt's Number | | | | | | |
| | (c) | Mach Number | (d) | Chezy's coefficient. | | | | | | |
| 12. | Coi | mpressibility of a fluid is the reciprocal of | | | | | | | | |
| | (a) | acceleration due to gravity | (b) 1 | nodulus of rigidity | | | | | | |
| | (c) | bulk modulus of elasticity | (d) | Poisson's Ratio. | | | | | | |
| 13. | A l
obt | iquid has a viscosity of 0.005 Pas and a d ained correctly as | ensit | y of 850 kg/m ³ . The kinematic viscosity may be | | | | | | |
| | (a) | $5.882\mu m^{2}/s$ | (b) | $0.0652 \times 10^{-4} m^2/s$ | | | | | | |

(c) 5.895 stoke (d) 0.00582 kg/m³

14.	 The atmospheric pressure measured as 760 (a) 1.033 m of water column (b) 0.1033 N/mm² (c) 10.33 N/mm² (d) 10.33 m of water column 	nm of mercury co	olumn is equal to									
	(e) none of the above											
15.	Surface tension has the dimension											
	(a) F (c) EI^{-2}	(b) FL^{-1} (d) FL^{-3}										
16	(c) FL	(u) FL^{+}										
10.	(a) adhesion	(b) cohesion										
	(c) adhesion and cohesion	(d) neither ac	dhesion nor coh	esion.								
17.	A perfect gas											
	(a) has zero viscosity	(b) has const	tant viscosity									
	(c) is incompressible	(d) satisfies p	$\rho = \rho RT$									
18.	The bulk modulus of elasticity											
	(a) is independent of pressure and visco	ty										
	(b) increases with the pressure											
	(c) has the dimensions of $\left(\frac{1}{p}\right)$											
	(d) is larger when the fluid is more compr	ssible.										
19.	An ideal fluid is the one which is											
	(a) non-viscous and incompressible	(b) compress	ible and has low	v density								
	(c) elastic and viscous	(d) steady an	steady and incompressible									
20.	Which fluid does not experience shear stre	s during flow?										
	(a) pseudoplastic	(b) dilatant										
	(c) inviscid	(d) Newtonia	in									
	Δ	swers										
1. (C	c) 2. (d) 3. (b) 4. (d) 5. (b) c) 12. (c) 13. (a) 14. (d) 15. (b)	6. (b) 7. (c) 16. (c) 17. (c)	;) 8. (d) I) 18. (b)	9. (a) 19. (a)	10. (d) 20. (c)							
			-, ··· (-,		()							
	Fill in the blanks											
1.	1. The rise of sap in a tree is a phenomenon of											

- 2. The spherical shape of rainwater drop is a phenomenon of _____.
- 3. Cavitation is a phenomenon due to _____.

- 4. The flow of an oil jet without break up is due to _____.
- 5. Cohesion acts between molecules of _____.
- 6. If Mach Number is less than 0.3, the fluid is _____.
- 7. Mass ______ and density are the same parameter.
- 8. Pascal is the S.I. unit of _____.
- 9. An ideal fluid is _____ and _____.
- 10. Printers ink is a _____ Newtonian fluid.
- 11. In Newton's law of viscosity the flow must be _____.
- 12. The velocity distribution of fluid following Newton's law of viscosity may be linear or ______.
- 13. Capillary angle in pure water in air in a clear glass tube is _____
- 14. Capillary rise of mercury in glass tube is _____ upward.

Answers

capillarity 2. surface tension 3. vapour pressure 4. viscosity 5. same fluid 6. incompressible 7. density
 stress 9. inviscid incompressible 10. non 11. laminar 12. non-linear 13. zero degree 14. convex

NUMERICAL EXAMPLES

EXAMPLE 1

A circular disc of diameter 'd' is slowly rotated in a liquid of large viscosity ' μ 'at a small distance 'h' from a fixed surface. Derive an expression for torque 'T' necessary to maintain an angular velocity ' ω '. Solution

Let us consider an annular element of width 'dr'at a distance 'r' from centre line.

From Newton's law, shear stress
$$\tau = \mu \frac{du}{dy}$$
 (a)

$$\therefore \qquad \text{Shear force } dF_s = \tau \times \text{Area} = \mu \frac{du}{dy} (2\pi r \cdot dr) \qquad (b)$$

$$\therefore \qquad \text{Torque produced } dT = dF_s \times r = \mu \frac{du}{dy} \cdot (2\pi r^2 \cdot dr)$$

As 'h' is very small, velocity distribution may be considered linear.

$$\therefore \qquad \frac{du}{dy} = \frac{r\omega}{h}$$



So, $dT = \mu \left(\frac{r\omega}{h}\right) \cdot 2\pi r^2 dr$

$$= \frac{2\pi\mu\omega}{h} \cdot r^{3}dr$$

Total torque $T = \int_{0}^{d/2} dT = \frac{2\pi\mu\omega}{h} \int_{0}^{d/2} r^{3}dr = \frac{\pi\mu\omega d^{4}}{32h}$ (d)

:.

A fluid has an absolute viscosity of 0.048 Pas and a specific gravity of 0.913. Such fluid flow over flat solid surface, the velocity at a point 75 mm away from the surface is 1.125 m/s. Calculate shear stresses at solid boundary, at points 25 mm, 50 mm and 75 mm away from boundary surface. Assume (i) linear and (ii) parabolic velocity distribution with the vertex at the point 75 mm away from the surface where the velocity is 1.125 m/s.

SOLUTION

(i) For linear distribution, v = 1.125 m/s at y = 75 mm = 0.075 m

So,
$$\frac{dv}{dy} = \frac{1.125}{0.075} = 15$$

$$\therefore \text{ Shear stress} \qquad \tau = \mu \frac{dv}{dy} = 0.048 \times 15 = 0.75 \text{ N/m}^2$$

,

As, the velocity distribution is linear, $\frac{dv}{dy}$ = constant and so shear stress has the uniform value.

(ii) Let the equation of parabolic distribution:
$$v = A + By + Cy^2$$

at $y = 0, v = 0$ $\Rightarrow A = 0$
at $y = 0.075$ m, $v = 1.125$ m/s $\Rightarrow 0.075 B + (0.075)^2 C = 1.125$ (a)

and
$$\frac{dv}{dy} = 0$$
 $\Rightarrow B + 0.15 C = 0$ (b)

Solving (a) and (b), Therefore, B = 30, C = -200 $v = 30y - 200y^2$

÷



Hence at	y=0,	$\frac{dv}{dy} = 30,$	v = 0	$\Rightarrow \tau = 1.44 \text{ N/m}^2$
	<i>y</i> = 0.025,	$\frac{dv}{dy} = 20,$	v=0.625	$\Rightarrow \tau {=} 0.96 \ \text{N/m^2}$
	<i>y</i> = 0.050,	$\frac{dv}{dy} = 10,$	v = 0.880	$\Rightarrow \tau = 0.48 \text{ N/m}^2$
	<i>y</i> = 0.050,	$\frac{dv}{dy} = 0,$	v=1.125	$\Rightarrow \tau = 0.$
EXAMPLE 3				

A cylinder of 0.12 radius rotates concentrically inside a fixed cylinder of 0.13 m radius. Both the cylinders are 0.3 m long. Determine the viscosity of the liquid which fills the space between the cylinders if a torque of 0.88 Nm is required to maintain angular velocity of 2π rad/s.

SOLUTION

Torque = Shear stress \times Surface area \times Lever arm $0.88 = \tau \times (2\pi r \times 0.3) \times r$ \Rightarrow (a) $\tau = \frac{0.467}{r^2}$ \Rightarrow $\frac{dv}{dr} = \frac{\tau}{\mu} = \frac{0.467}{\mu r^2}$ Now $dv = \frac{0.467 dr}{\mu r^2}$ \Rightarrow (b) Integrating (b), $\int_{v_{out}}^{v_{in}} dv = \frac{0.467}{\mu} \int_{0.13}^{0.12} \frac{dr}{r^2}$ 0.01 m 0.24 m $V_{\rm in} - V_{\rm out} = \frac{0.467}{\mu} \left(\frac{1}{0.12} - \frac{1}{0.13} \right)$ 0.3 m \Rightarrow $\omega \times 0.12 - 0 = \frac{0.2993}{\mu}$ \Rightarrow $2\pi \times 0.12 \quad = \quad \frac{0.2993}{\mu}$ \Rightarrow = 0.3970 Pas \Rightarrow μ

Hence, absolute value is 0.3970 Pas

The velocity profile is a laminar flow through a round pipe is expressed as $u = 2U\left(1 - \frac{r^2}{r_0^2}\right)$, where U = average velocity, r_0 = radius of pipe. Draw dimensionless shear stress profile $\left(\frac{\tau}{\tau_0}\right)$ against $\left(\frac{r}{r_0}\right)$ where τ_0 is wall shear stress. Find τ_0 , when oil flows with absolute viscosity 4 × 10⁻² Ns/m² and velocity of 4 m/s in a pipe of diameter 150 mm.

SOLUTION

Given:

÷

and

So,

$$\frac{\tau}{\tau_0} = \frac{r}{r_0}$$
 and the plot is shown.

 $u = 2U\left(1 - \frac{r^2}{r_0^2}\right)$

 $\tau = \mu \frac{du}{dr} = -\frac{4\mu Ur}{r_0^2}$

 $\tau_0 = \mu \frac{du}{dr} = -\frac{4\mu U}{r}$

 $\frac{du}{dr} = -\frac{4Ur}{r_0^2}$



From (c), $\tau_0 = -\frac{4 \times 4 \times 10^{-2} \times 4}{0.075} = 8.534 \text{ N/m}^2$ (in absolute value)

EXAMPLE 5

Neglecting temperature effect, the empirical pressure-density relation for water is $\left(\frac{p}{p_a}\right) = 3001 \times \left(\frac{\rho}{\rho_a}\right)^7 - 3000$. Determine isothermal bulk modulus of elasticity and compressibility of water at 1,10 and 100 atmospheric pressure.

SOLUTION

The relation given, $\left(\frac{p}{p_a}\right) = 3001 \times \left(\frac{\rho}{\rho_a}\right)^{\prime} - 3000.$

Differentiating with respect to p,

$$\frac{1}{p_a} \cdot \frac{dp}{d\rho} = 7 \times 3001 \times \frac{\rho^6}{\rho_a^7}$$
$$\frac{dp}{d\rho} = 7 \times 3001 \times p_a \times \frac{\rho^6}{\rho_a^7}$$

 \Rightarrow

(a)

$$\Rightarrow \qquad \rho \frac{dp}{d\rho} = 7 \times 3001 \times p_a \times \left(\frac{\rho}{\rho_a}\right)^7 \qquad (b)$$

From equation (1.7),
$$E = \rho \frac{dp}{d\rho}$$

$$= 7 \times 3001 \times p_a \times \left(\frac{\rho}{\rho_a}\right)^7$$

$$= 7 \times 3001 \times p_a \times \left(\frac{p}{\rho_a} + 3000\right) \times \frac{1}{3001}$$

$$= 7 \times p_a \times \left(\frac{p}{\rho_a} + 3000\right) \qquad (c)$$

From equation (c)
So, at
$$p = 1p_{a}, \qquad E_1 = 21007p_a$$

$$p = 10p_{a}, \qquad E_{10} = 21700p_a$$

$$p = 100p_a, \qquad E_{100} = 21700p_a$$
If $p_a = 1.0132 \times 10^5$ N/m²,

$$E_1 = 0.4698 \times 10^{-9}$$
 m²/N

$$E_{10} = 2.1348 \times 10^9$$
 N/m², $K_1 = 0.4698 \times 10^{-9}$ m²/N

$$E_{100} = 2.1986 \times 10^9 \text{ N/m}^2, K_{100} = 0.4548 \times 10^{-9} \text{ m}^2/\text{N}$$

A cylinder 0.35 m³ of air at 50°C and 276 kN/m² absolute. The air is compressed to 0.071 m³. (i) Assuming isothermal condition, what is the pressure at new volume and bulk modulus of elasticity at new state? (ii) Assuming isentropic condition, what is the pressure and bulk modulus of elasticity for $\gamma = 1.4$.

SOLUTION

(i) For isothermal condition

$$p_1 \forall_1 = p_2 \forall_2$$

$$\Rightarrow \qquad 0.35 \times (276 \times 10^3) = 0.071 \times p_2$$

$$\Rightarrow \qquad p_2 = 1360 \times 10^3 \text{ N/m}^2 = 1.36 \text{ MPa}$$

⇒ $p_2 = 1360 \times 10^{\circ}$ N/m ∴ Isothermal bulk modulus = $p_2 = 1.36$ MPa

(ii) For isentropic condition,

 $p_1 \forall_1^{1.4} = p_2 \forall_2^{1.4} \\ \Rightarrow \qquad (0.35)^{1.4} \times (276 \times 10^3) = (0.071)^{1.4} \times p_2 \\ \Rightarrow \qquad \qquad p_2 = 2.575 \text{ MPa}$

:. Isentropic bulk modulus $=\gamma p_2 = 1.4 \times 2.575 = 3.605$ MPa

The mean free path of a gas is $\lambda = 1.26 \frac{\mu}{\rho} (RT)^{-1/2}$. Find out air density when mean free path is

10 mm, temperature 20°C, $\mu = 1.8 \times 10^{-5}$ kg/ms, R = 287 J/kg K.

 $\lambda = 1.26 \, \frac{\mu}{\rho} (RT)^{-1/2}$

SOLUTION

Here, $\lambda = 10 \text{ mm} = 0.01 \text{ m}, \mu = 1.8 \times 10^{-5} \text{ kg/ms}$ $T = 20^{\circ}\text{C} = 293 \text{ K}, R = 287 \text{ J/kg K}$

From the given relation:

 \Rightarrow

$$\rho = 1.26 \frac{\mu}{\lambda} (RT)^{-1/2}$$

= 1.26 \times $\frac{1.8 \times 10^{-5}}{0.01} (287 \times 293)^{-1/2}$
= 7.82 \times 10^{-6} kg/m³.

EXAMPLE 8

A shaft 80 mm in diameter is being pushed through a bearing sleeve of 80.2mm diameter and 0.3 m long. Assumed uniform clearance is flooded with lubricating oil of viscosity 0.1 kg/ms and specific gravity 0.9. (i) If the shaft moves axially at 0.8 m/s, calculate resistance force exerted by oil on the shaft. (ii) If shaft is axially fixed and rotated at 1800 r.p.m., estimate resisting torque exerted by oil and power required.

SOLUTION

(i) Given,

...

$$\mu = 0.1 \text{ kg/ms}, l = 0.3 \text{ m}, y = \frac{0.2 \text{ mm}}{2} = 0.0001 \text{ m},$$

$$v = 0.8 \text{ m/s}, d = 80 \text{ mm} = 0.08 \text{ m}$$
Resisting force = $\tau \times \pi dl$

$$= \left(\mu \times \frac{v}{y} \times \pi dl\right)$$

$$= \left(0.1 \times \frac{0.8}{0.0001} \times \pi \times 0.08 \times 0.3\right) \text{N}$$

$$= 60.32 \text{ N}$$

$$0.1 \text{ mm}$$

$$(////////) \downarrow$$

(ii) Given,

 $\mu = 0.1$ kg/ms, l = 0.3 m, y = 0.0001 m, d = 0.08 m, N = 1800 rpm

$$v = \frac{\pi dN}{60} = 7.539 \text{ m/s}$$

$$\therefore \text{ Resisting force} = \left(\mu \times \frac{\nu}{\nu}\right) \times \pi dl$$
$$= \left(0.1 \times \frac{7.539}{0.0001} \times \pi \times 0.08 \times 0.3\right) \text{N}$$
$$= 568.4272 \text{ N}$$
$$\therefore \text{ Resisting torque} = (568.4272 \times 0.04) \text{ Nm}$$
$$= 22.737 \text{ Nm}$$
$$\therefore \text{ Power transmitted} = \left(22.737 \times \frac{2\pi \times 1800}{60} \times \frac{1}{1000}\right) \text{kW}$$
$$= 4.285 \text{ kW}$$

A body weighing 1000 N slides down at a uniform speed of 1 m/s along a lubricated inclined plane making 30° with the horizontal. The viscosity of lubricant is 0.1 kg/ms and contact area of the body is 0.25 m^2 . Determine the lubricant thickness assuming linear velocity distribution.

SOLUTION

Let the thickness of oil = h m

$$\therefore$$
 Velocity gradient $\left(\frac{du}{dy}\right) = \frac{1}{h}$ per sec

:. Shearing viscous force acting(F)

$$= \mu \cdot \frac{du}{dy} \cdot A$$



$$= 0.1 \times \frac{1}{h} \times 0.25 = \frac{0.025}{h} \text{ kgm/s}^2 = \frac{0.025}{h} \text{ N}$$

Horizontal component of the body = $1000 \sin 30^\circ = 500 \text{ N}$

... From condition of equilibrium,
$$\frac{0.025}{h} = 500 \text{ mm}$$

 $\Rightarrow \qquad h = 5 \times 10^{-5} \text{ m} = 0.05 \text{ mm}$

EXAMPLE 10

A uniform film of oil 0.13 mm thick separates two discs, each of 200 mm diameter, mounted co-axially. Ignoring the edge effects, calculate the torque necessary to rotate one disc relative to other at a speed of 7 r.p.m., if the oil has a viscosity of 0.14 Pas.



SOLUTION

Let us consider an annular element of width 'dr', at a distance r from centreline.

Here,
$$\mu \frac{du}{dy} = \frac{2\pi rN}{dy} \ \mu = \frac{2\pi r \times 7}{0.13 \times 10^{-3}} \times 0.14 = 47365.55 \ r$$

:. Force acting on this element $(dF) = 47365.55 \times 2\pi r dr$

:. Torque acting on this element $(dT) = dF \times r = 297606.53 r^3 dr$

:. Total torque
$$(T) = \int dT = \int_{0}^{0.1} 297606.53r^3 dr = 7.44$$
 Nm.

EXAMPLE 11

A piston 79.6 mm diameter and 210 mm long works in a cylinder of 80 mm diameter. If the annular space is filled with a lubricating oil of viscosity 0.065 kg/ms, calculate the speed at which the piston will move through the cylinder, when an axial load of 85.6 N is applied. Neglect inertia of the piston.

SOLUTION

Effective area subjected to shear $(A) = (\pi \times 0.0796 \times 0.21) \text{ m}^2$ Let the speed of rotation = v m/s

 \therefore Viscous force acting on the piston (*F*)

$$= 0.065 \times \frac{\nu}{\left(\frac{80 - 79.6}{2}\right) \times 10^{-3}} \times (\pi \times 0.0796 \times 0.21) \,\mathrm{N}$$

$$=(17.0673 \times v) N$$

As per given condition:

$$17.0673 \times v = 85.6$$

 $v = 5.01$ m/s.

 \Rightarrow



Two large plane surfaces are 3 cm apart. A very thin plate of negligible thickness with surface area 0.5 m² is placed between the surfaces while above the plate is filled up with glycerine and below the plate is filled up with lubricating oil. Find out the force required to drag the plate at a velocity of 0.45 m/s when (i) it lies exactly at middle of two surfaces (ii) it is 2 cm away from bottom surface, $\mu_g = 0.81$ Pas, $\mu_0 = 0.245$ Pas.

SOLUTION

Effective area subjected to shear $(A) = 0.5 \text{ m}^2$ Velocity of drag (V) = 0.45 m/s

Total Drag force = Drag force due to glycerine + Drag force due to oil \Rightarrow $F = F_c + F_0$

(i) Here
$$y_G = y_0^G = 3 \text{ cm}/2 = 1.5 \text{ cm}$$

:.
$$F_G = \mu_G \times \frac{V}{y_G} \times A = \left(0.81 \times \frac{0.45}{1.5 \times 10^{-2}} \times 0.5 \right) \text{N}$$

= 12.15 N

and

$$F_0 = \mu_0 \times \frac{V}{y_0} \times A = \left(0.245 \times \frac{0.45}{1.5 \times 10^{-2}} \times 0.5\right) N$$
$$= 3.675 N$$

:. Total Force (F)
$$= (12.15 + 3.675) \text{ N} = 15.825 \text{ N}$$

(i) Here
$$y_0 = 2 \text{ cm}, y_G = 3 - 2 = 1 \text{ cm}$$

:.
$$F_G = \mu_G \times \frac{V}{y_G} \times A = \left(0.81 \times \frac{0.45}{1 \times 10^{-2}} \times 0.5\right) \text{N} = 18.225 \text{ N}$$

$$F_0 = \mu_0 \times \frac{V}{y_0} \times A = \left(0.245 \times \frac{0.45}{2 \times 10^{-2}} \times 0.5\right)$$
N = 2.756 N

:. Total Force
$$(F) = (18.225 + 2.756)$$
 N = 20.981 N.

EXAMPLE 13

A vertical gap 2.25cm wide of infinite extent contains a fluid of viscosity 1.98 Pas and specific gravity 0.92. A metal plate of $1.25 \text{ m} \times 1.25 \text{ m} \times 0.25 \text{ cm}$ needs to be lifted with a constant velocity of 0.18 m/s, through the gap. Assuming the plate lying exactly at the middle of the gap, find the required force of lift. Mass density of plate is 1415 kg/m³.

SOLUTION

Volume of the plate $= (1.25 \times 1.25 \times 0.25 \times 10^{-2}) \text{ m}^3$ $= 3.90625 \times 10^{-3} \text{ m}^3$





Area of plate subjected to shear

$$=(1.25 \times 1.25) \text{ m}^2 = 1.5625 \text{ m}^2$$

 $= 1 \times 10^{-2} \text{ m}$

:. Shear force acting on each side of plate

$$= \left(1.98 \times \frac{0.18}{1 \times 10^{-2}} \times 1.5625\right)$$
 N = 55.687 N

 \therefore Total shear force on both side of plate (F_2)

$$= (2 \times 55.687) \text{ N} = 111.374 \text{ N}$$

To lift up the plate, the total shear force (F_2) and effective downward force (F_1) of the plate will have to overcome.

So the net drag force $F = F_1 + F_2$

$$= (18.963 + 111.374) N$$
$$= 130.337 N.$$

EXAMPLE 14

The surface tension of water in contact with air at 20° C is 0.0705 N/m. The pressure inside a water droplet is 0.018 N/cm² higher than outside pressure. What is the diameter of the water droplet?

SOLUTION

From equation (1.10), $\Delta p = 0.018 \text{ N/cm}^2 = 180 \text{ N/m}^2$

$$\sigma = 0.0705 \text{ N/m}$$

and $R_1 = R_2 = R$

...

$$\sigma = \frac{\Delta p}{\left(\frac{2}{R}\right)}$$
$$R = \frac{2\sigma}{\Delta p} = \frac{2 \times 0.0705}{180} = 7.834 \times 10^{-4} \text{ m}$$

 \Rightarrow

So, diameter of the droplet = $(2 \times 7.834 \times 10^{-4})$ m

 $= 1.5668 \times 10^{-3}$ m = 1.57 mm.

EXAMPLE 15

Find surface tension of a soap bubble of 48 mm diameter while pressure inside is 3.12 Pa higher than atmospheric one.

SOLUTION

As it is a hollow bubble it will have two sets of interface surface, one is inner surface and inside air, other is outer surface and outside air. And surface tension will act both outside surface and inside surface of bubble. Hence

$$R = \frac{4\sigma}{\Delta p}$$
$$\sigma = \frac{\Delta p \cdot R}{4} = \frac{3.12 \times 48 \times 10^{-3}}{4 \times 2} \text{ N/m} = 0.01872 \text{ Nm}.$$

 \Rightarrow

EXAMPLE 16

Calculate capillary rise in a glass table of 2.35 mm diameter, when vertically immersed in (i) mercury (ii) water. Surface tension in contact with air is 0.518 N/m for mercury and 0.07238 N/m for water. Specific gravity and angle of contact for mercury is 13.596 and 130° respectively.

SOLUTION

Here,
$$\rho_w = 1000 \text{ kg/m}^3$$
, $\sigma_w = 0.07238 \text{ N/m}$
 $\rho_m = (13.596 \times 1000) \text{ kg/m}^3$, $\sigma_m = 0.518 \text{ N/m}$
 $\theta_w = 0^\circ, \theta_m = 130^\circ$
 \therefore Capillary rise for water

$$h_w = \frac{4\sigma_w \cos \theta_w}{\rho_w \times g \times d} = \frac{4 \times 0.07238 \times \cos 0^{\circ}}{1000 \times 9.807 \times 2.35 \times 10^{-3}} \text{ m}$$
$$= 0.01256 \text{ m} = 12.56 \text{ mm}$$

Capillary depression for mercury

$$h_m = \frac{4\sigma_m \cos\theta_m}{\rho_m \times g \times d} = \frac{4 \times 0.518 \times \cos 130^\circ}{(13.596 \times 1000) \times 9.807 \times 2.35 \times 10^{-3}} \text{ m}$$
$$= -4.2505 \times 10^{-3} \text{ m} = -4.25 \text{ mm}$$

Here, (-) ve sign indicates capillary fall instead of capillary rise.

EXAMPLE 17

The space between a shaft and a concentric sleeve is filled up with Newtonian fluid. If a force of 550 N is applied parallel to the shaft, the sleeve attains a speed of 1.3 m/s. How much speed will attain the sleeve, if above force is increased to 1550 N, the temperature being the same?

SOLUTION

Let the clearance between shaft and sleeve be 'dy', radius of the shaft 'r' and length of sleeve 'l'.

Viscosity remaining same, we obtain two sets of equations

$$550 = \mu \times \frac{1.3}{dy} \times (2\pi rl)$$

$$1550 = \mu \times \frac{V}{dy} \times (2\pi rl)$$

V = 3.66 m/s.

Doing (b) ÷ (a),
$$\frac{1550}{550} = \frac{V}{1.3}$$



⇒

EXAMPLE 18

A hydraulic ram 200 mm diameter and 1.2 m long moves within a concentric cylinder 200.2 mm diameter. The annular clearance is filled with oil of relative density 0.85 and kinematic viscosity 400 mm²/s. What is the viscous force resisting the motion when the ram moves at a speed of 120 mm/s.

(a)

(b)

SOLUTION

Dynamic viscosity (µ)	= vp = (400 × 10 ⁻⁶ × 0.85 × 10 ³) Ns/m ² = 0.34 Ns/m ²
So, Shear stress τ	$=\mu \times \frac{du}{dy}$
	$= 0.34 \times \frac{120}{\left(\frac{200.2 - 200}{2}\right)}$
	$=408 \text{ N/m}^2$
\therefore Shear force (<i>F</i>)	$= \tau \times (\pi dl)$
	$=408 \times (\pi \times 200 \times 10^{-3} \times 1.2) \text{ N}$
	= 307.62 N.
EXAMPLE 19	

The space between two large flat and parallel walls 25 mm apart is filled with liquid of viscosity 0.7 Ns/m^2 . Within this space a thin flat plate 250×250 is towed at a velocity of 150 mm/s at a distance of 19 mm from one wall, the plate and its movement being parallel to the walls. Assuming linear variations of velocity between the plate and the walls, determine the force exerted by the liquid on the plate.



A large plate moves with speed v_0 over a stationary plate on a layer of oil. If the velocity profile is that of a parabola, with the oil at the plates having the same velocity as the plates, what is the shear stress on the moving plate from the oil? If a linear profile is assumed, what is the shear stress on the upper plate?

SOLUTION

Let us assume equation of parabolic velocity profile



At

:..

$$y = d, \tau = \mu \frac{dv}{dy}\Big|_{y=d} = \frac{\mu}{2} \cdot \frac{v_0}{\sqrt{d}} \cdot \frac{1}{\sqrt{d}} = \frac{\mu v_0}{2d}$$

dering linear profile $v = k v$

Considering linear profile, $v = k_2 y$

at
$$y = d$$
, $v = v_0$ $\Rightarrow k_2 = \frac{v_0}{d}$
 $v = \frac{v_0}{d} \cdot y$

$$\Rightarrow \qquad \frac{dv}{dy} = \frac{v_0}{d}$$

$$\therefore \qquad \tau = \mu \frac{dv}{dy} = \mu \frac{v_0}{d}$$

At
$$y = d, \tau = \mu \frac{v_0}{d}.$$

1. (a) State and explain Newton's law of viscosity.

(b) Define Newtonian and Non-Newtonian fluids with examples.

- 2. Define the following terms
 - (a) Surface Tension (b) Rheology

EXERCISE

- (c) Ideal fluid (d) Bulk modulus of elasticity of a compressible fluid
- 3. Explain dynamic viscosity and kinematic viscosity. Give their dimensions.
- 4. Explain the phenomenon of capillarity. Obtain an expression for capillary rise of a fluid.
- 5. A plunger is moving through a cylinder at a speed of 6.1 m/s. The diameter and length of the plunger are 126.7 mm and 25 mm respectively. The diameter of the cylinder is 127 mm. The film of oil separating plunger from the cylinder has a viscosity of 0.95765 N-s/m². What is the force required to maintain this motion?

Ans. [12.275 N]

 Calculate the density, specific weight and volume of chloride gas at 25°C and pressure of 6,00,000 N/m² absolute.

Ans. [17.1 kg/m³, 168 N/m³, 0.0585 m³/kg]

7. A high pressure steel container is partially full of a liquid at a pressure of 10 atm. The volume of the liquid is 1.232 L. At a pressure of 25 atm, the volume of the liquid equals 1.231 L. What is the average bulk modulus of elasticity of the liquid over the given range of pressure of the temperature after compression is allowed to return to the original temperature. What is the coefficient of compressibility?

Ans. $[1.872 \times 10^6 \text{ k Pa}, 0.000534 \text{ m}^2/\text{MN}]$

8. Water is moving through a pipe. The velocity profile is given by $v = \frac{\beta}{4\mu} \left(\frac{d^2}{4} - r^2\right)$ What is the shear

stress at the wall of the pipe and at r = d/4. If the given profile persists a distance L along the pipe, what drag is induced by the water on the pipe in the direction of flow over this distance?

Ans.
$$\left[-\frac{2\beta r}{4}, -\frac{\beta d}{8}, \frac{\beta \pi d^2 L}{4}\right]$$

9. A cylinder of weight 9.1 kg, diameter 152 mm and height 127 mm slides through a lubricated pipe. The clearance between cylinder and pipe is 0.025 mm. If the cylinder is observed to decelerate @ 0.61 m/s² when the speed is 6.1 m/s, what is the viscosity of the oil?

Ans. [6.411 × 10⁻³ Pas]

10. In a stream of glycerine in motion, at a certain point, the velocity gradient is 0.25 m/s/m. Mass density of fluid is 1268.4 kg/m^3 and kinematic viscosity is $6.30 \times 10^{-4} \text{ m}^2/\text{s}$. Calculate shear stress at that point.

Ans. $[0.2 \text{ N/m}^2]$

11. Surface tension of water in contact with air at 20°C is given as 0.0716 N/m. Pressure inside a droplet of water is to be 0.0147 N/cm² greater than the outside pressure, calculate diameter of droplet of water.

Ans. [1.94 mm]

12. Find surface tension in a soap bubble of 30 mm diameter when the inside pressure is 1.962 N/m² above atmosphere.

Ans. [0.00735 N/mm]

13. Capillary rise in the glass tube used for measuring water level is not to exceed 0.5 mm. Determine its minimum size, given that surface tension for water in contract with air = 0.07112 N/m.

Ans. [5.8 cm]

14. A very large thin plate is centered in a gap of width 6 cm with different oil of unknown viscosities above and below, the viscosity of one being twice that of the other. When pulled at a velocity of 30 cm/s, the resulting force on 1 m² of plate due to viscous shear on both sides is 29.4 N. Assuming viscous flow and neglecting all end effects, calculate viscosities of the oil.

Ans. [0.98 Pas, 1.96 Pas]



FLUID STATICS

2.1 FORCES ON FLUID ELEMENT

If an infinitesimally small region in the continuum of fluid can be separated out, then that particular region is defined as a *fluid element*. A fluid element, on isolation from all its surroundings is subjected to variety of forces, whether it is in rest or in motion. All the forces acting on any fluid element can be classified broadly into two groups:

- (a) **Body Forces:** are forces acting throughout the body of the fluid element and distributed over entire volume of the element. Forces due to gravitational acceleration, electromagnetic field are such body forces. These are expressed as force per unit mass.
- (b) Surface Forces: are forces acting directly on the surface of fluid element in contact. Viscous force is such a surface force. All surface forces are normally resolved along normal direction and perpendicular to the normal (tangential) direction. Hence surface forces have two components, one is *normal force* and the other is *tangential force*.

2.1.1 Pascal's Law of Hydrostatics

Let us consider a small tetrahedral fluid element OABC as shown in Figure 2.1. Triangular face ABC has an infinitesimally small area dA, which makes angle α with *x*-plane(bounded by *y*-axis and *z*-axis),angle β with *y*-plane(bounded by *x*-axis and *z*-axis) and angle γ with *z*-plane(bounded by *x*-axis and *y*-axis). Now projection of 'dA' on *x*-plane, *y*-plane and *z*-plane are

$$dA \cdot \cos \alpha = \frac{1}{2} dy \cdot dz$$
$$dA \cdot \cos \beta = \frac{1}{2} dx \cdot dz$$
$$dA \cdot \cos \gamma = \frac{1}{2} dx \cdot dy$$



Now, σ is pressure of fluid in any arbitrary direction and σ_x , σ_y , σ_z indicate pressure along *x*-axis, *y*-axis and *z*-axis respectively.

To attain conditions of equilibrium of the tetrahedral element, the normal surface forces due to fluid pressure along the co-ordinate axes will have to be resolved. Therefore

$$\sigma_{x} \frac{dy \, dz}{2} - \sigma \cdot dA \cdot \cos \alpha = 0$$

$$\sigma_{y} \frac{dx \, dz}{2} - \sigma \cdot dA \cdot \cos \beta = 0$$

$$\sigma_{z} \frac{dx \, dy}{2} - \sigma \cdot dA \cdot \cos \gamma = 0$$
(2.2)

Clubbing (2.1) & (2.2), we obtain,

$$\begin{array}{l}
\sigma_x = 0 \\
\sigma_y = 0 \\
\sigma_z = 0
\end{array}$$
(2.3)

or

$$=\sigma_{x}=\sigma_{y}=\sigma_{z}$$
(2.4)

As the direction of σ acting normal to the surface *dA* is arbitrary, *pressure at a point in a static fluid is the same in all direction*. This is known as *Pascal's Law* of hydrostatics, as established by a French Mathematician, *B.Pascal* in 1653. Very commonly the stress σ in equation (2.4) is denoted as *p*, a scalar quantity.

2.2 EQUILIBRIUM OF STATIC FLUID ELEMENT

Let us consider a differential fluid element in the shape of a parallelepiped, shown in figure 2.2. It is subjected to both body force and surface force. Body force arises due to earth's field of gravitation and surface force due to normal pressure acting on element.

σ

The component of body force along co-ordinate axes are X, Y and Z, per unit mass. Mass density of the element is ρ . Normal pressure p is acting on rear surface of the element in x-plane, as shown.

Rate of change of p along dx length = $\frac{dp}{\partial x}$



Fig. 2.2

So, the normal pressure on front surface of the element in x-plane = $\left(p + \frac{\partial p}{\partial x} dx\right)$

Considering equilibrium of forces along x-axis

$$p \, dy \, dz - \left(p + \frac{\partial p}{\partial x} \, dx\right) dy \, dz + X \, \rho \, dx \, dy \, dz = 0$$

$$\Rightarrow \qquad p - p - \frac{\partial p}{\partial x} dx + X \, \rho dx = 0$$

$$\Rightarrow \qquad \frac{1}{\rho} \frac{\partial p}{\partial x} = X \qquad (2.5)$$

In the similar fashion, we can consider the equilibrium along *y*-axis and *z*-axis, finally yielding *Euler's Equations of Equilibrium* as

$$\frac{1}{\rho}\frac{\partial p}{\partial x} = X \tag{2.6a}$$

$$\frac{1}{\rho}\frac{\partial p}{\partial y} = Y \tag{2.6b}$$

$$\frac{1}{\rho}\frac{\partial p}{\partial z} = Z \tag{2.6c}$$

2.3 SOLUTION OF EULER'S EQUATIONS

Let us consider a specific case, where the fluid is static and *z*-axis coincides with the direction of gravitational acceleration. Here component of body forces along *x*-axis and *y*-axis are zero, *i.e.*, X = 0 = Y.

Due to special choice of coordinate axes, the partial differential form in equation (2.6c) can be changed to total differential form. Replacing Z by -g [(–)ve sign due to downward direction of g], above equation can be re-written as

$$\frac{1}{\rho}\frac{dp}{dz} = -g \tag{2.7}$$

Solution of this equation may be found for the following cases.

2.3.1 Incompressible Fluid (uniform density)

Adjoining figure 2.3 illustrates this specific condition. To find expression of pressure p at any z height, we are to integrate equation (2.7).

$$\int \frac{1}{\rho} dp = \int -g \, dz$$
$$\frac{p}{\rho} = -gz + C_0$$



 \Rightarrow

291

$$p = -\rho gz + C, \text{ where } C = \rho \cdot C_0$$
Applying boundary condition as, at $z = H, p = p_a$

$$p_a = -\rho gH + C$$

$$C = p_a + \rho gH$$
Substituting in equation (2.8), $p = -\rho gz + p_a + \rho gH$

$$(2.8)$$

 \Rightarrow

 \Rightarrow

Substituting in equation (2.8), $p = -\rho gz + p_a + \rho gH$ $p = p_a + \rho g (H - z) = p_a + \rho gh$ (2.10) Here this expression is formed in terms of absolute pressure. In terms of gauge pressure, $p_a = 0$ and it will be, $p = \rho g h$ (2.11)

2.3.2 Isothermal State

At initial reference level, say, on earth's surface, the pressure and mass density of compressible fluid are p_0 and ρ_0 respectively. Applying Boyle's Law,

$$\frac{p}{\rho} = \frac{p_0}{\rho_0} \tag{2.12}$$

Substituting in equation(2.7), $\frac{1}{\rho}dp = -gdz$

$$\frac{dp}{p} \cdot \frac{p_0}{\rho_0} = -gdz \qquad [substituting from eq. 2.12] \quad (2.13)$$

Integrating

 \Rightarrow

$$\frac{p_0}{\rho_0} \int_{p_0}^{p} \frac{dp}{p} = -g \int_{z_0}^{z} dz$$

$$\frac{p_0}{\rho_0} \ln\left(\frac{p}{p_0}\right) = -g(z - z_0)$$

$$\ln\left(\frac{p}{p_0}\right) = -\frac{\rho_0 g}{p_0}(z - z_0)$$
(2.14)

$$\Rightarrow$$

 \Rightarrow

 \Rightarrow

 $\frac{p}{p_0} = \exp\left[-\frac{\rho_0 g}{p_0} \left(z - z_0\right)\right]$ (2.15)

For a perfect gas,

$$\frac{p_0}{\rho_0} = R T_0 \tag{2.16}$$

Substituting in equation (2.15),

$$\frac{p}{p_0} = \exp\left[-\frac{g}{RT_0}(z-z_0)\right]$$

 \Rightarrow

 \Rightarrow

 \Rightarrow

$$p = p_0 \cdot \exp\left[\frac{-g\left(z - z_0\right)}{RT_0}\right]$$
(2.17)

2.3.3 Adiabatic State

For an adiabatic process,

$$\frac{p}{\rho^{\kappa}} = \frac{p_0}{\rho_0^{\kappa}} \tag{2.18}$$

$$\rho = \rho_0 \left(\frac{p}{p_0}\right)^{1/k} \tag{2.19}$$

Substituting in equation (2.7),

$$\frac{1}{\rho_0} \left(\frac{p_0}{p}\right)^{1/k} \cdot dp = -gdz \tag{2.20}$$

Integrating

 \Rightarrow

 \Rightarrow

$$\frac{1}{\rho_0} \cdot p_0^{1/k} \int_{p_0}^p p^{-1/k} \cdot dp = -g \int_{z_0}^Z dz$$
(2.21)

$$p_0^{1/k} \cdot \left[\frac{p^{1-1/k}}{1-1/k} \right]_{p_0}^p = -g\rho_0 \left(z - z_0 \right)$$

$$\Rightarrow \qquad p_0^{1/k} \cdot \left(\frac{k}{k-1}\right) \cdot \left[p^{1-1/k} - p_0^{1-1/k}\right] = -g\rho_0(z-z_0) \tag{2.22}$$

$$g\rho_0(z-z_0) = p_0 \cdot p_0^{1/k-1} \cdot \left(\frac{k}{k-1}\right) \cdot \left[p_0^{1-1/k} - p^{1-1/k}\right]$$
(2.23)

So, from equation (2.23),

$$(z - z_0) = \left(\frac{k}{k-1}\right) \frac{p_0}{\rho_0 g} \left[1 - \left(\frac{p}{p_0}\right)^{\frac{k-1}{k}}\right]$$
$$\frac{p}{p_0} = \left[1 - (z - z_0) \frac{\rho_0 g}{p_0} \left(\frac{k-1}{k}\right)\right]^{\frac{k}{k-1}}$$
$$= \left[1 - (z - z_0) \frac{g}{RT_0} \left(\frac{k-1}{k}\right)\right]^{\frac{k}{k-1}}$$
(2.24)

 \Rightarrow

2.3.4 **Polytropic State**

In a polytropic process,

$$\frac{p}{\rho^n} = \frac{p_0}{\rho_0^n} \tag{2.25}$$

Following equation (2.24) and (2.25), we can write,

$$\frac{p}{p_0} = \left[1 - \left(z - z_0\right) \frac{g}{RT_0} \left(\frac{n-1}{n}\right)\right]^{\frac{n}{n-1}}$$
(2.26)

100

80

60

40

20

0

-100

-80

-60

Altitude km

lonosphere

Mesosphere

-40

Temperature °C

Fig. 2.4

Stratosphere

-20

Troposphere

0

20

2.3.5 Temperature Lapse Rate

The variation of temperature in atmosphere with altitude from M.S.L. is shown in figure 2.4. Troposphere is in contact with the surface of the earth up to an altitude of 7 km in polar and 14 km in equatorial region, averaging 11 km in general. And lapse rate (α) of temperature is constant, 6.5 K/km, and $T_0 = 288$ K, while

$$T = T_0 - \alpha z \qquad (2.27)$$
 As we know, $p = \rho RT$ and from equation (2.7)

dzdn 28)

$$\frac{dp}{p} = -\frac{g}{R} \frac{dz}{(T_0 - \alpha z)}$$
(2.1)

Integrating within limits,

$$\ln\left(\frac{p}{p_0}\right) = \frac{g}{R\alpha} \ln\left(\frac{T_0 - \alpha z}{T_0 - \alpha z_0}\right)$$
$$\Rightarrow \qquad \frac{p}{p_0} = \exp\left[\frac{g}{R\alpha} \ln\left(\frac{T_0 - \alpha z}{T_0 - \alpha z_0}\right)\right] \tag{2.29}$$

2.4 GAUGE PRESSURE AND ABSOLUTE PRESSURE

The air in the atmosphere exerts a natural pressure on all the surfaces in contact, which is known as Atmospheric Pressure. It varies with change of altitude and can be measured with the help of a barometer. Hence atmospheric pressure is sometimes denoted as Barometric Pressure. In normal condition this pressure at mean sea level is

76 cm of mercury column

=10.33 m of water column

 $=10,330 \text{ kgf/m}^2 = 101.32 \text{ kN/m}^2 = 1.013 \text{ bar}$

The pressure exerted by any fluid can be measured either (i) absolute zero as datum, or (ii) atmospheric pressure as datum. When pressure is measured considering vacuum or absolute zero as datum, it is called an *Absolute Pressure*. If the datum be atmospheric pressure, the measured pressure is termed as *Gauge Pressure*. So,

Gauge pressure = Absolute pressure – Atmospheric pressure.

If gauge pressure is below atmospheric pressure, then it is called *Negative gauge pressure* or *Vacuum pressure* or *Suction pressure*.

(2.30)

2.4.1 Barometer

The instrument which measures atmospheric pressure is known as barometer and is shown in figure 2.5. The column of mercury rises in the inverted tube by capillarity and the space above this column apparently empty but actually is filled up with mercury vapour, is known as *Torricellian Vacuum*. Hence the atmospheric pressure can be expressed as,

$$p_{\rm atm} = p_{\rm vap} + \rho g h$$

Generally, the effect of vapour pressure is neglected in computation of atmospheric pressure but an error of
$$0.14\%$$
 to 0.16% is invoked.

2.5 MEASUREMENT OF PRESSURE

There happens to exist number of devices, applying which the fluid pressure can be measured.





Vapour Pressure

p,

2.6 MANOMETER

A manometer is a device for measuring the difference between two pressures by balancing a column of liquid.

2.6.1 Simple Manometer

It measures pressure at point in a fluid contained in a conduit or vessel.

(a) **Piezometer:** It is an ordinary vertical tube, connected to a conduit carrying liquid is used to measure gauge pressure, and is called a piezometer tube or simply a piezometer, as 'piezo' means 'pressure'.

The liquid in piezometer rises freely until equilibrium. In figure 2.6, 'h' is the pressure head at A. Piezometer cannot measure pressure too high due to large length of tube and vacuum pressure due to air entrance inside conduit.

(b) U-tube manometer: A *U*-tube, as shown in figure 2.7, may be used to measure pressure in a conduit or container having a fluid of mass density ρ_A . The manometric tube should contain a liquid having mass density ρ_B , such that $\rho_B > \rho_A$. For water and alcohol, common manometric fluid is mercury, while for gas it may be water.



Fig. 2.6

Now in order to write an equation for pressure equilibrium, following mnemonic, first suggested by *Professor John Foss* of Michigan State University, can be used.

)

$$p_{\downarrow} = p^{\uparrow} + \gamma |\Delta z| \qquad (2.31)$$

Following this, pressure at *P*

$$p_p = p_1 + \rho_A \cdot g (y + x)$$

Then jump across along the same line at Q. As level is not changing, then

$$p_{Q} = p_{p} = p_{1} + \rho_{A} \cdot g(y + x)$$
 (2.33)

Now, pressure at open end of the tube

$$p_{atm} = p_Q - \rho_B \cdot g \cdot x$$

= $p_1 + \rho_A \cdot g (y + x) - \rho_B \cdot g \cdot x$ (2.34a)
$$\Rightarrow p_1 = p_{atm} + (\rho_B - \rho_A) g \cdot x - \rho_A \cdot g \cdot y$$
 (2.34b)

(c) Inclined tube manometer: It is just a modification of an ordinary *U*-tube manometer, for improvement of sensitivity, shown in figure 2.8. The c/s area of the reservoir is made much higher



Fig. 2.7

than that of the inclined tube. This arrangement enables the manometric liquid remain practically constant in the reservoir due to a small change in pressure. The inclined tube facilitates the movement 'l' of the meniscus to be recorded with greater sensitivity, and satisfies the relation

$$h = l\sin\theta \tag{2.35}$$

Fluid Statics

Here smaller the inclination, higher is the sensitivity. But an inclination less than 5° makes the reading of the exact position of meniscus troublesome and hence should not be used.

2.6.2 **Differential Manometer**

It is used to determine a small or moderately high pressure difference between any two points in two conduits or two





Fig. 2.9



Normally U-tube differential manometer is employed to ascertain high pressure difference (Fig. 2.9). Inverted U-tube manometer, shown in figure 2.10, generally measures low or moderate pressure difference. The procedure for obtaining equilibrium equation for pressure difference is same as before. Referring figure 2.9.

$$p_{A} + \rho_{1}gz_{1} - \rho_{2}gz_{2} - \rho_{3}gz_{3} = p_{B}$$

$$\Rightarrow \qquad p_{A} - p_{B}$$

$$= \rho_{2}gz_{2} + \rho_{3}gz_{3} - \rho_{1}gz_{1} \qquad (2.36)$$

Again for inverted U-tube manometer in figure 2.10,

$$p_{A} - \rho_{1}gz_{1} + \rho_{2}gz_{2} + \rho_{3}gz_{3} = p_{B}$$

$$\Rightarrow \quad p_{A} - p_{B} = \rho_{1}gz_{1} - \rho_{2}gz_{2} - \rho_{3}gz_{3}$$
(2.37)

2.6.2.1 Micromanometer

Sometimes it is necessary to measure small pressure difference accurately and micromanometer is employed. As shown in figure 2.11, it is seen that other than manometric fluid and working fluid, another gauge fluid is used. The equation of hydrostatic equilibrium will be



Fig. 2.8

$$p_{1} + \rho_{w} \cdot g \left(h + \Delta z\right) + \rho_{G} \cdot g \left(z + \frac{y}{2} - \Delta z\right)$$
$$- \rho_{m} \cdot g \cdot y - \rho_{G} \cdot g \left(z - \frac{y}{2} + \Delta z\right) - \rho_{w} \cdot g \left(h - \Delta z\right) = p_{2}$$
(2.38)

Here $\rho_m > \rho_G > \rho_w$ From continuity of gauge liquid

$$A \cdot \Delta z = a \cdot y/2 \tag{2.39}$$

h





Clubbing (2.38) and (2.39), we find

$$p_{1} + \rho_{w} \cdot g\left(h + \frac{a}{A} \cdot \frac{y}{2}\right) + \rho_{G} \cdot g\left(z + \frac{y}{2} - \frac{a}{A} \cdot \frac{y}{2}\right) - \rho_{m} \cdot g \cdot y$$
$$-\rho_{G} \cdot g\left(z - \frac{y}{2} + \frac{a}{A} \cdot \frac{y}{2}\right) - \rho_{w} \cdot g\left(h - \frac{a}{A} \cdot \frac{y}{2}\right) = p_{2}$$
$$\Rightarrow \quad p_{1} + \rho_{w} \cdot g \cdot \frac{a}{A} \cdot y + \rho_{G} \cdot g \cdot y - \rho_{G} \cdot g \frac{a}{A} \cdot y - \rho_{m} \cdot g \cdot y = p_{2}$$
$$\Rightarrow \quad p_{1} - p_{2} = gy\left[\rho_{m} - \rho_{G}\left(1 - \frac{a}{A}\right) - \rho_{w} \frac{a}{A}\right]$$
(2.40)
If
$$A \gg a, \frac{a}{A} \approx 0$$

Therefore, $p_1 - p_2 \approx gy (\rho_m - \rho_G)$

2.7 HYDROSTATIC FORCE ON SUBMERGED PLANE SURFACE

Let us consider a plane surface of arbitrary shape and of area A, completely submerged below the free surface of liquid MN. The plane makes an angle θ with the horizontal. The reference system of axes is shown in figure 2.12. The force acting on the elemental area dA at a vertical depth hcan be written

 $dF = \rho g h \, dA \tag{2.42}$

On integration over the entire area *A*, total force acts as



 \Rightarrow

$$F = \int_{A} dF = \int_{A} \rho g h \, dA = \int_{A} \rho g y \sin \theta \, dA \qquad [\because h = y \sin \theta]$$
$$= \rho g \sin \theta \int_{A} y \, dA \qquad (2.43)$$

As we know, centroidal coordinate of the plane is $G(\bar{x}, \bar{y})$

$$\overline{y} = \frac{\int_{A} y \, dA}{\int_{A} dA} \qquad \text{So, } \int_{A} y \, dA = A \overline{y}$$

Substituting in equation (2.43),

$$F = \rho g \sin \theta \left(A \overline{y} \right)$$

= $\rho g A \overline{h}$ [$\because \overline{h} = \overline{y} \sin \theta$] (2.44)

This is the resultant hydrostatic force acting on the plane. Let us consider the point $C(x_c, y_c)$ as centre of pressure, where the above force acts. Considering moment equilibrium about *x*-axis, we can write,

$$y_{C} \cdot F = \int_{A} y \, dF$$

$$y_{C} = \frac{1}{F} \int_{A} y \cdot (\rho g y \sin \theta \, dA)$$

$$= \frac{\int_{A} (\rho g y^{2} \sin \theta \, dA)}{\int_{A} (\rho g y \sin \theta \, dA)} = \frac{\int_{A} y^{2} dA}{\int_{A} y \, dA}$$
(2.45)

Similarly, from moment equilibrium about *y*-axis,

$$x_{C} = \frac{1}{F} \int_{A} x \, dF$$
$$= \frac{\int_{A} (\rho g \, xy \sin \theta \, dA)}{\int_{A} \rho gy \sin \theta \, dA} = \frac{\int_{A} xy \, dA}{\int_{A} y \, dA}$$
(2.46)

Again, second moment of area about x-axis, applying parallel axis theorem,

$$\int_{A} y^{2} dA = I_{x} = I_{\overline{x}} + A\overline{y}^{2}$$
(2.47 a)

and product moment of area, applying parallel axis theorem

$$\int_{A} xy \, dA = I_{xy} = I_{\overline{x} \ \overline{y}} + A\overline{x} \ \overline{y} \tag{2.47 b}$$

Substituting in equation (2.45) and (2.46),

$$y_C = \frac{I_{\overline{x}} + A\overline{y}^2}{A\overline{y}} = \frac{I_{\overline{x}}}{A\overline{y}} + \overline{y}$$
(2.48 a)

$$x_{C} = \frac{I_{\overline{x}\,\overline{y}} + A\overline{x}\,\overline{y}}{A\overline{y}} = \frac{I_{\overline{x}\,\overline{y}}}{A\overline{y}} + \overline{x}$$
(2.48 b)

It is of special mention that, while the plane area becomes symmetrical about any of the centroidal axes $x = \overline{x}$ or $y = \overline{y}$, then $I_{\overline{x}\overline{y}}$ becomes zero and

$$d = y_C - \overline{y} = \frac{I_{\overline{x}}}{A\overline{y}} \tag{2.49}$$

2.8 HYDROSTATIC FORCE ON SUBMERGED CURVED SURFACE

A curved surface of total area *A* in 3-*D* Cartesian coordinate system is shown in figure 2.13. Its outer surface is subjected to atmospheric pressure, while the inner surface under liquid pressure. Free surface of the liquid coincides with x - y plane. An elemental area *dA* is considered, while an elemental force *dF* acts along the outer normal to this area. If *l*, *m*, and *n* be the direction cosines of the outer normal to *dA*, then



$$dA_{y} = dA \cdot m = dA \cos\beta$$

$$dA_{z} = dA \cdot n = dA \cos\gamma$$
(2.50)

Here dA_x , dA_y , dA_z are the projections of dA on x-plane, y-plane and z-plane, respectively. So, along x, y and z direction, components of forces are

$$dF_{x} = \rho gz dA_{x}$$
$$dF_{y} = \rho gz dA_{y}$$
$$dF_{z} = \rho gz dA_{z}$$

 $dA = dA \cdot l = dA \cos \alpha$

On integration total forces on those three planes are

$$F_x = \int \rho g z dA_x = \rho g \overline{z}_x A_x \tag{2.51 a}$$

$$F_{y} = \int \rho g z dA_{y} = \rho g \overline{z}_{y} A_{y}$$
(2.51 b)

$$F_z = \int \rho g z dA_z = \rho g z A_z = \rho g \forall$$
 (2.51 c)

 \Rightarrow

Here \overline{z}_x and \overline{z}_y are centroid of the projected area A_x and A_y respectively and \forall is the volume of liquid supported vertically by the surface A. To find the coordinate of centre of pressure C of the horizontal force F_x ,

$$z_{C} = \frac{\int z \cdot dF_{x}}{F_{x}} \quad \text{and} \quad y_{C} = \frac{\int y \cdot dF_{x}}{F_{x}}$$
$$= \frac{\int z^{2} \cdot dA_{x}}{A_{x} \cdot \overline{z}_{x}} \quad = \frac{\int yz \cdot dA_{x}}{A_{x} \cdot \overline{z}_{x}}$$
$$= \frac{I_{y}}{A_{x} \cdot \overline{z}_{x}} \quad = \frac{I_{yz}}{A_{x} \cdot \overline{z}_{x}} \quad (2.52)$$

2.9 BUOYANCY & ARCHIMEDES LAW

While a body is submerged *entirely* or *partially* in a *non-reacting*, *insoluble*, *static* fluid, it experiences an upward vertical lift. This phenomenon is called 'buoyancy' and the upward thrust generated is 'buoyant force'.

Referring to figure 2.14, the resultant horizontal force in any direction for such a closed surface is zero. To evaluate the vertical resolution of resultant hydrostatic force, an elemental vertical prism of c/s are dA_z is assumed, the top surface which is at z_1 height and bottom surface at z_2 height. Considering the condition of vertical equilibrium,

 $dF_{B} = dF_{2} - dF_{1}$



$$= \rho g (z_2 - z_1) dA_z$$

= $\rho g d \forall$ (2.53)

Where $d\forall$ is the volume of the elemental prism. Integrating the equation we obtain,

 $= p_2 \cdot dA_z - p_1 \cdot dA_z$

$$F_{R} = \int_{\forall} \rho g d \forall = \rho g \forall \tag{2.54}$$

Here \forall is total volume of the body. To obtain the point of action of buoyant force,

$$x_{B} \cdot F_{B} = \int x \cdot dF_{B}$$

$$x_{B} = \frac{\int x \cdot \rho g d \forall}{\rho g \forall} = \frac{\int x d \forall}{\forall}$$
(2.55)

And x_{B} is the centroid of the displaced volume of liquid. Explaining this, *Archimedes* (285–212 B.C.) propounded his famous law, which states,

If a body is immersed partially or entirely in a non-reacting, insoluble, static fluid, it experiences a reduction of its actual weight and this reduction equals with weight of the displaced volume of fluid.

2.10 EQUILIBRIUM AND METACENTRE

A body, submerged or floating in a fluid may experiences three types of hydrostatic equilibrium, depending on the relative position of centre of gravity (G) and centre of buoyancy (B) of the fluid. They are (i) stable equilibrium (ii) unstable equilibrium and (iii) neutral equilibrium.

A body is said to be in stable equilibrium, if being tilted by a small angle and then released, returns back to its initial position of equilibrium, by keeping the original vertical axis as vertical.

A body is said to be in unstable equilibrium, if on being tilted and released, not returns back to its original position but moves further away from it.

A body is said to be in neutral equilibrium, which when tilted by a small angle, neither returns back to its original position nor moves away from it, and simply adapts equilibrium in its new position.

A pictorial illustration is provided in figure 2.15 to show that in (i) stable equilibrium, G is below B (ii) unstable equilibrium, G is above 'B' and (iii) neutral equilibrium, G coincides with B.

In vertically stable condition G and B lies on the same vertical line. On tilt, G lies on the same old line while Bshifts away. A vertical line passing through new position of B intersects the old line holding G at M. This point is defined as metacentre, and GM as metacentric height. From the adjacent figure 2.16, it can be said that,



2.10.1 Determination of Metacentric Height

Figure 2.17 shows c/s of a floating boat in equilibrium condition, with respective position of G and B. In tilted position, elemental prism (dotted line) mn of c/s dA changes to m'n'. The coordinates of B and B' are (x, y, z) and (x', y', z')



302

Initial volume $\forall = \int_{A} z dA$ (2.57 a) New immersed volume $\forall' = \int_{A} (z + x\alpha) dA$ (2.57 b) If the rotation α is so provided that both the above $\frac{1}{77}$ volumes remaining the same,

$$\int_{A} z dA = \int_{A} (z + x\alpha) dA$$
$$\Rightarrow \qquad \int_{A} x \cdot dA = 0$$



This implies that *y*-axis must pass through centroid of area *A*. Now,

$$\forall \cdot x = \int_{A} x \cdot z dA, \qquad \forall \cdot y = \int_{A} y \cdot z dA, \qquad \forall \cdot z = \frac{1}{2} \int_{A} z^{2} dA \qquad (2.59)$$

$$\forall' \cdot x' = \int_{A} x(z + x\alpha) dA, \qquad \forall' \cdot y' = \int_{A} y(z + x\alpha) dA, \qquad \forall' \cdot z' = \frac{1}{2} \int_{A} (z^{2} - x^{2} \alpha^{2}) dA$$

So, considering,
$$\forall = \forall'$$

$$BB' = x' - x = \frac{\alpha}{\forall} \int_{A} x^{2} dA = \frac{\alpha}{\forall} \cdot I_{y}$$

(2.58)

$$\Rightarrow BM \cdot \alpha = \frac{\alpha}{\forall} \cdot I_{y}$$

$$\Rightarrow BM = \frac{I_{y}}{\forall}$$
(2.60)

Hence, $\overline{GM} = \overline{BM} - \overline{BG} = \frac{I_y}{\forall} - \overline{BG}$

Multiple Choice Questions

- 1. Pressure at a point in a static mass of liquid depends upon
 - (a) shape and size of container
 - (b) depth below free liquid surface
 - (c) specific weight of liquid and depth below free liquid surface
 - (d) all the above
- 2. Differential form of pressure variation in static fluid is (z measured vertically upward)

(a)
$$dp = -\gamma dz$$
 (b) $dp = -\rho dz$

(c)
$$d\rho = -\gamma dz$$
 (d) $dp = -zd\rho$

3. Atmospheric pressure is

- (a) gauge pr. Abs. pr. (b) abs. pr. Gauge pr.
- (c) abs. pr. Vacuum pr. (d) gauge pr. + Vacuum pr.

(2.61)

304					Mechanical Science-II						
4	. Pre	essure in metre of oil (sp. gr. = 0.78), equiva	lent to	80 m of water is							
	(a)	64.32	(b)	80.64							
	(c)	102.56	(d)	88.18							
5	. In	isothermal atmosphere, the pressure									
	(a)	varies in same way as in density									
	(b)	linearly decrease with elevation									
	(c)	exponentially increase with elevation (d)	rema	ins constant							
6	б. То	avoid capillary correction in manometer, d	iamet	er of tube should be							
	(a)	=5 mm	(b)	<3 mm							
	(c)	<5 mm,>3 mm	(d)	≥6mm							
7	. Ce	ntre of pressure for a plane, vertically imme	ersed	in static mass of fluid is							
	(a)	always above centroid of area									
	(b)	always coinciding with centroid of area									
	(c)	sometimes above, sometimes below centre	roid o	farea							
	(d)	always below centroid of area									
8	В. Но	Horizontal force component on a submerged curved surface is									
	(a)	weight of liquid vertically above the surface	ace								
	(b)	weight of liquid retained in the surface									
	(c)	product of area and pr. at centroid	C								
	(d)	force on a vertical projection of curved si	urface								
9	. Ve	rtical force component on a submerged cur	ved s	urface is							
	(a)	force on horizontal projection of curved s	surfac	e							
	(\mathbf{D})	weight of liquid vertically above curved	urfoo	9							
	(d)	its horizontal component	surrac	C							
10		small boat corrying 5 passengers in a pond	If al	I passangars jump out o	fboat to pond the water						
10	lev	rel in pond will	. 11 ai	i passengers jump out o	t boat to polid, the water						
	(a)	rise (b)	fall	(c)	not change						
11	. Bu	ovant force is			-						
	(a)	resultant force on floating body from flui	d sur	ounding							
	(b)	resultant force on any body									
	(c)	volume of displaced liquid									
	(d)	force necessary to maintain equilibrium of	of a su	bmerged body							
12	. Lii	ne of action of buoyant force acts through									
	(a)	C.G. of any submerged body									
	(b)	centroid of volume of any floating body									
	(c)	centroid of displaced volume of fluid									

- (d) centroid of horizontal projection of body

- 13. A body will float in stable equilibrium, if
 - (a) metacentric height zero
 - (b) C.G. is below centre of buoyancy
 - (c) metacentric height positive and C.G. above centre of buoyancy
 - (d) metacentre above C.G.
- 14. Normal stress is same in all direction at a point in a fluid
 - (a) when the fluid is at rest
 - (b) when fluid is inviscid
 - (c) when fluid is inviscid and incompressible
- 15. Hydrostatic force on one side of a circular surface of unit area, with centroid 10 m below free water level (density ρ) is
 - (a) $< 10 \, \rho g$
 - (b) = $10 \rho g$
 - (c) $\rho g \times \text{vert.dist.}$ from free surface to centre of pressure
 - (d) none of the above.
- 16. Differential manometer is used to measure
 - (a) velocity at a point in fluid
 - (b) pressure at a point in fluid
 - (c) difference of pressure between two points in fluid
 - (d) none of the above
- 17. Water column equivalent to 760 mm of mercury is
 - (a) 1.033 m (b) 10.33 m
 - (c) 0.033 Nmm (d) 10.33 N/mm^2

18. Hydrostatic law states that rate of increase of pressure in a vertical direction is equal to

- (a) density of fluid (b) specific gravity of fluid
- (c) specific weight of fluid (d) none of the above
- 19. Sensitiveness of an ordinary *U*-tube manometer can be increased without changing manometer fluid by
 - (a) setting it horizontal
 - (c) setting it inclined
- 20. A manometer may be (a) single limbed

(c) multilimbed

- (b) double limbed
 - (d) all of the above

(b) keeping it vertical

(d) none of the above

- 21. Mercury is used in barometer on account of
 - (a) negligible capillary effect
 - (c) very low vapour pressure
- (b) high density
- (d) low compressibility

22.	The	e problems of hydrostatics are influenced b	oy fol	lowing forces
	(a)	gravity and viscous force	(b)	gravity and pressure force
	(c)	viscous and surface tension force	(d)	gravity and surface tension force
23.	Ind	licate the variation of hydrostatic pressure	with	depth below free surface
	(a)	it decreases with increase in depth		
	(b)	there is no change with depth		
	(c)	it increases with increase in depth		
	(d)	after a certain depth, there is no change i	n pre	essure
24.	Ab	solute pressure in flow system		
	(a)	always above local atm. pr.		
	(b)	is a vacuum pressure		
	(c)	may be above, below or equal to the local	atm.	pr.
	(d)	is also called negative pr.		
25.	Loo	cal atmospheric pressure is measured by		
	(a)	barometer	(b)	bourdon gauge
	(c)	vacuum gauge	(d)	manometer
26.	Ga	uge pressure in flow system are measured	by	
	(a)	manometer	(b)	aneroid barometer
	(c)	vacuum gauge	(d)	bourdon gauge
27.	Ab	solute pressure is measured by		
	(a)	Bourdon gauge	(b)	aneroid barometer
	(c)	differential manometer	(d)	hook gauge
28.	Poi	int through which resultant hydrostatic for	ce ac	ts is called
	(a)	metacentre	(b)	centre of pressure
	(c)	centre of buoyancy	(d)	centre of gravity
29.	Me	ercury is used in manometer for measuring		
	(a)	low pressure accurately	(b)	large pressure only
	(c)	all pressure except smaller ones	(d)	very low pressure
30.	A f	loating body displaces a volume of liquid	equal	to
	(a)	its own volume	(b)	its own weight
	(c)	its submerged weight	(d)	none of the above
31.	Cei	ntre of buoyancy is		
	(a)	the point through submerged weight of l	oody	acts
	(b)	the point through which buoyancy force	acts	
	(c)	the point through which resultant hydros	static	force acts
	(d)	centre of gravity of liquid displaced by the	ne bo	dy

32.	Met	tacentr	ric he	ight	is the	dista	nce b	etwee	en									
	(a)	C.G.a	and m	netace	entre				(b)	C.G. a	nd ce	ntre of l	buoya	ncy			
	(c)	metac	centre	e and	centr	re of l	buoya	ncy	(d)	none o	of the	above					
33.	Met	tacentr	ric he	ight	of a f	loatin	ig bod	ly dep	bends	on								
	(a)	the sh	ape o	of its	wate	r-line	area		(b)	volum	e of l	iquid di	isplac	ed by	the bo	ody	
	(c)	distar	ice be	etwee	en me	tacent	tre and	dC.G	. ((d)	none o	of the	above					
34. A floating body will remain in stable equilibrium so long as																		
(a) metacentre lies below C.G.					((b)	metacentre above C.G.											
	(c)	metac	centre	e and	C.G.	at the	same	posit	ion (d)	none o	of the	above					
35.	Ad	ifferen	itial n	nanoi	neter	is use	ed fore	emea	surem	ent	of							
	(a)	press	ure a	t a po	oint				((b)	veloci	ty at a	a point					
	(c)	differ	ence	of pi	essur	e at t	wo po	ints	(d)	discha	rge						
								-										
								Α	nswe	ers								
1. (c	;) 2.	(a)	3.	(b)	4.	(c)	5.	(c)	6.	(d)	7.	(d)	8.	(d)	9.	(c)	10.	(b)
11. (c	d) 12.	(c)	13.	(d)	14.	(a)	15.	(c)	16.	(c)	17.	(b)	18.	(c)	19.	(c)	20.	(d)
21. (c	;) 22.	(b)	23.	(c)	24.	(c)	25.	(a)	26.	(a)	27.	(b)	28.	(b)	29.	(c)	30.	(b)

31. (b) 32. (a,b) 33. (a) 34. (b) 35. (c)

True or False

- 1. Law of Archimedes is valid for a vessel in free fall.
- 2. A submarine can cruise smoothly under water only when C.G. is above centre of buoyancy.
- 3. To decrease frequency of rolling, the metacentric height of a passenger ship is kept lower than that of naval or cargoship.
- 4. Centre of pressure is independent of orientation of the area.
- 5. A piezometer is used to measure gauge pressure of liquid.
- 6. Pascal's Law is not valid for incompressible fluid.
- 7. Different diameters in two limbs of a U-tube manometer will not affect operation of manometer.
- 8. Pressure decreases as depth of liquid increases.
- 9. Find pressure on a surface rarely acts normal to the surface.
- 10 Study of pressure exerted by fluid in motion is known as hydrostatics.

Answers

1.	True	2.	False	3.	True	4.	False	5.	True
6.	False	7.	False	8.	False	9.	False	10.	False
Fill in the Blanks

1	Body and Surface force acting on a fluid element in static equilibrium are and												
2	Pressure at all points on a horizontal plane in a continuous fluid at rest is												
3	Pressure in a fluid at rest remains on a plane normal to the body force.												
4	Buoyant force in case of a satellite moving in a circular orbit is												
5	Once it starts to rise, an inflated balloon attains a definite height, when the buoyant force the weight.												
6	Centre of gravity of the displaced fluid is the centre of												
7	. The resultant force exerted on a body by a surrounding static fluid is force.												
8	3. Centre of pressure is the on the submerged area at which resultant force is supposed to act.												
9	. Atmospheric pressure is also known as pressure.												
10	. In measuring gauge pressure with an instrument pressure is taken as datum.												
11	1. Pressure atmospheric pressure is known as vacuum pressure.												
12	Any pressure measured taking datum as zero is termed as pressure.												
13	3. Name of S.I. unit of pressure is												
14	A pressure of absolute zero can exist only in complete												
15	is a simplest form of single –limb manometer.												
16	An inclined single limbed manometer is useful for measurement of pressure.												
17	manometer is used to measure pressure difference between two points in closed conduit.												
18	Inverted U-tube differential manometer is used for measuring difference of pressure.												
19	Generally for measurement of high pressure gauges are used.												
20	The most common type of mechanical pressure gauge is pressure gauge.												
	Answers												
1.	weight, pressure 2. same 3. constant 4. zero 5. equals												
6. 11.	buoyancy 7. buoyant 8. point, hydrostatic 9. barometric 10. atmospheric below 12. absolute 13. Pascal 14. vacuum 15. Piezometer												
16.	small 17. Differential 18. low 19. mechanical 20. Bourdon tube												
	NUMERICAL EXAMPLES												

EXAMPLE 1

A barometer is made by dipping a graduated inverted tube into the measuring liquid open in atmosphere. Estimate the height of liquid column in the barometer where atmospheric pressure is 100 kN/m^2 . (a) when the liquid is mercury and (b) when the liquid is water. The measuring temperature is 50° C where vapour

Fluid Statics

pressure of mercury and water 0.015×10^4 N/m² and 1.23×10^4 N/m² and densities are 13500 kg/m³ and 980 kg/m³ respectively. What would be % error if effect of vapour is neglected?

SOLUTION

(a) Equation of pressure equilibrium including effect of vapour pressure is $100 \times 10^3 = 0.015 \times 10^4 + 13500 \times 9.807 \times h_m$ $h_m = 0.754 \text{ m}$ \Rightarrow without considering effect of vapour pressure, $100 \times 10^3 = 13500 \times 9.807 \times h'_m$ $h'_{m} = 0.755 \text{ m}$ \Rightarrow Percentage error $=\frac{0.755 - 0.754}{0.754} \times 100 = 0.1326$ *:*.. (b) Equation of pressure equilibrium including effect of vapour pressure is $100 \times 10^3 = 1.23 \times 10^4 + 980 \times 9.807 \times h_{\rm w}$ $h_{w} = 9.125 \text{ m}$ \Rightarrow without considering effect of vapour pressure, $100 \times 10^3 = 980 \times 9.807 \times h'_{w}$ $h'_{w} = 10.405 \text{ m}$ \Rightarrow Percentage error $=\frac{10.405 - 9.125}{9.125} \times 100 = 14.027$. *:*..

EXAMPLE 2

Density (kg/m³) of a fluid mixture varies with z as
$$\rho = 10.1 \left[1 - \frac{z}{500} + \left(\frac{z}{1000} \right)^2 \right]$$
. Assuming the

mixture to be stationary, determine the pressure difference between bottom and top of a 60 m tall reactor. **SOLUTION**

As we know
$$\frac{dp}{dz} = -\rho g$$

$$dp = -\rho g dz = 10.1g \left[1 - \frac{z}{500} + \left(\frac{z}{1000}\right)^2 \right] dz$$

Integrating,

$$\int_{p_0}^{p} dp = -10.1g \int_{0}^{60} \left[1 - \frac{z}{500} + \left(\frac{z}{1000}\right)^2 \right] dz$$
$$p - p_0 = -10.1g \left[z - \frac{z^2}{2 \times 500} + \frac{z^3}{3 \times 10^6} \right]_{0}^{60}$$
$$= -10.1 \times 9.807 \left(60 - \frac{60^2}{1000} + \frac{60^3}{3 \times 10^6} \right)$$

 \Rightarrow

:..

(a)

$$= -5.593 \times 10^3 \text{ N/m}^2$$

So, the difference of pressure is 5.593 kN/m^2

EXAMPLE 3

Find the atmospheric pressure just at the end of troposphere which extends upto a height of 11.02 km from sea level. Consider a temperature variation in the tropospheric as $T = 288.16 - 6.49 \times 10^{-3}z$, where z in metre and T in Kelvin. Atmospheric pressure at sea level is 101.32 kN/m^2 .

SOLUTION

As we know

$$\frac{dp}{dz} = -\rho g$$
 and $p = \rho RT$

:.

$$dp = -\rho g dz = \frac{-p}{RT} g dz$$

 \Rightarrow

$$\frac{dp}{p} = -\frac{gdz}{RT} = -\frac{g}{R} \cdot \frac{dz}{\left(288.16 - 6.49 \times 10^{-3} z\right)}$$
(a)

Integrating equation (a)

$$\ln\left(\frac{p}{p_0}\right) = \frac{g}{R} \cdot \ln\left[\frac{(288.16 - 6.49 \times 10^{-3} z)}{(288.16 - 6.49 \times 10^{-3} z_0)}\right] \times \frac{1}{6.49 \times 10^{-3}}$$
(b)

Now, $p_0 = 101.32 \times 10^3$ N/m², $z_0 = 0$, $z = 11.02 \times 10^3$ m, R = 287 J/kg K Substituting in equation (b)

$$\ln p - \ln(101.32 \times 10^{3})$$

$$= \frac{9.807}{287} \ln \left[\frac{\left(288.16 - 6.49 \times 10^{-3} \times 11.02 \times 10^{3} \right)}{288.16} \right] \times \frac{1}{6.49 \times 10^{-3}}$$

$$\ln p = 10.024$$

$$p = (e)^{10.024} = 22561.76 \text{ N/m}^{2} = 22.562 \text{ kN/m}^{2}.$$

⇒

 \Rightarrow

Find the pressure at an elevation of 300 m above the sea level by assuming (a) an isothermal condition of air and (b) an isothermoic condition of air. Pressure and temperature at sea level are 101.32 kN/m² and 293.15 K. Consider air to be an ideal gas with R = 287 J/kgK and $\gamma = 1.4$

SOLUTION

(a) As we know
$$\frac{dp}{\rho} = -g \cdot dz$$
 and $\frac{p}{\rho} = \frac{p_0}{\rho_0}$

 $\therefore \qquad \frac{dp}{p} \cdot \frac{p_0}{\rho_0} = -g \cdot dz$

Integrating (a),

$$\frac{p_0}{p_0} \int_{p_0}^{p} \frac{dp}{p} = -g \int_{z_0}^{z} dz$$
 (b)

Finally from equation (2.17), we obtain,

$$p = p_0 \cdot \exp\left[\frac{-g(z - z_0)}{RT_0}\right]$$
(c)

[Students will have to complete the integration obtained from (b) to (c). Otherwise marks will be deducted in University Exam.]

Here,

$$p_0 = 101.32 \text{ kN/m}^2$$
, $T_0 = 293.15 \text{ K}$, $R = 287 \text{ J/kgK}$.
 $z_0 = 0, z = 3000 \text{ m}$.

Substituting in equation (c)

$$p = 71.42 \text{ kN/m}^2$$

(b) Go through article 2.3.3 and deduce the final equation as,

$$\frac{p}{p_0} = \left[1 - \left(z - z_0\right) \frac{g}{RT_0} \left(\frac{\gamma - 1}{\gamma}\right)\right]^{\frac{1}{\gamma - 1}} \tag{d}$$

Substituting the values and $\gamma = 1.4$

 $p = 70.07 \text{ kN/m}^2$.

EXAMPLE 5

What is the pressure intensity in the ocean at a depth of 1500 m, assuming (a) salt water incompressible with specific weight of 10050 N/m³ and (b) salt water compressible and weight 10050 N/m³ at the free surface and $E = 2070 \text{ MN/m}^2$ (constant).

SOLUTION

(a) Here pressure

 $p = \rho g h = 10050 \times 1500 \text{ N/m}^2 = 15.075 \text{ MPa}$

(b) As we know

$$\frac{1}{\rho}\frac{dp}{dz} = g \text{ and } E = \rho \frac{dp}{d\rho},$$

Clubbing both equations we obtain,

$$E\int_{\rho_0}^{\rho}\frac{d\rho}{\rho^2} = g\int_{z_0}^{z}dz$$

 $\frac{E \cdot d\rho}{\rho} = g\rho dz$

$$\Rightarrow \qquad \frac{E}{g} \left[-\frac{1}{\rho} \right]_{\rho_0}^{\rho} = (z - z_0)$$

$$\Rightarrow \qquad \frac{E}{g}\left(\frac{1}{\rho_0} - \frac{1}{\rho}\right) = \left(z - z_0\right)$$

(a)

$$\frac{\rho}{\rho_0} = \frac{E}{E - \rho_0 g \left(z - z_0\right)} \tag{b}$$

Again,

$$dp = E \frac{d\rho}{\rho} \tag{c}$$

Integrating,

$$\int_{p_0}^{p} dp = E \int_{\rho_0}^{\rho} \frac{d\rho}{\rho}$$
(d)

$$\Rightarrow$$

 $p - p_0 = E \cdot \ln\left(\frac{\rho}{\rho_0}\right)$ (e)

From equations (b) and (e)

$$p - p_0 = E \cdot \ln\left(\frac{E}{E - \rho_0 g\left(z - z_0\right)}\right) \tag{f}$$

Substituting the values,

$$p = 15.13$$
 MPa.

EXAMPLE 6

A hydraulic press has a ram of 28 cm diameter and plunger of 5 cm diameter. How much weight can be lifted by the press when 438 N force is applied on the piston?

SOLUTION

c/s area of ram =
$$\frac{\pi}{4} \times (0.28)^2 \text{ m}^2 = 0.06157 \text{ m}^2$$

c/s area of plunger = $\frac{\pi}{4} \times (0.05)^2 \text{ m}^2 = 1.9634 \times 10^{-3} \text{ m}^2$

Stress or pressure intensity on plunger

$$=\frac{438}{1.9634\times10^{-3}}$$
 N/m² = 223082.408N/m²

Same pressure intensity will act under ram, by virtue of Pascal's law \therefore Force generated over ram = (223082.408 × 0.06157) N

$$= 13.735 \text{ kN}.$$



EXAMPLE 7

In the adjacent figure, left tube is filled with water, right tube with CCl_4 (sp.gr. = 1.59) and middle bent part with compressed air at 300 kN/m². Barometer reads 76 cm of mercury. Determine

(a) pressure difference in kN/m² between A and B, if z = 0.45 m

(b) absolute pressure in *B* in mm of mercury.

SOLUTION

Following the mnemonic in (2.31), Absolute pressure at *A*

$$p_A = (300 \times 10^3 + 1.8 \times 10^3 \times 9.807) \text{ N/m}^2$$

= (317.6526 × 10³) N/m²

 \Rightarrow

Fluid Statics

Absolute pressure at B

$$p_{B} = [300 \times 10^{3} + (1.8 - 0.45) \times 10^{3} \times 9.807 \times 1.59] \text{ N/m}^{2}$$

= (321.0507 × 10³) N/m²

:. Pressure difference between A and B $p_{p} - p_{d} = (3.3981 \times 10^{3}) \text{ N/m}^{2}$

$$p_B - p_A = (3.3981 \times 10^{\circ})$$

= 3.391 kN/m²

Equivalent height of mercury column for absolute pressure at B

$$=\frac{321.0507\times10^3}{10^3\times9.807\times13.569}$$
 m = 2.4078m.

EXAMPLE 8

A multi-tube using water and mercury is used to measure pressure of air in a vessel, as shown in the figure. For the given values of height, calculate gauge pressure in the vessel. $h_1 = 0.5 \text{ m}, h_2 = 0.5 \text{ m}, h_3 = 0.3 \text{ m}, h_4 = 0.7 \text{ m}, h_5 = 0.1 \text{ m}, h_6 = 0.5 \text{ m}$

SOLUTION

Specific weight of air (γ_a) =12 N/m³ water (γ_w) =10³ × 9.807 = 9807 N/m³

mercury (γ_m) =13.596 × 9.807 = 133335.97 N/m³

As the pressure is to be expressed in gauge pressure, $p_{\text{atm}} = 0$ at open end of tube. Following the principle in (2.31), equation of pressure equilibrium is

$$p_{atm} + h_{6} \cdot \gamma_{m} + h_{5} \cdot \gamma_{m} - h_{5} \cdot \gamma_{w} - h_{4} \cdot \gamma_{w} + h_{4} \cdot \gamma_{m} \\ - h_{3} \cdot \gamma_{m} - (h_{2} - h_{3})\gamma_{w} + h_{2} \cdot \gamma_{m} - h_{1} \cdot \gamma_{a} = p \\ \Rightarrow \qquad p = 0 + 0.5\gamma_{m} + 0.1\gamma_{m} - 0.1\gamma_{w} - 0.7\gamma_{w} + 0.7\gamma_{m} - 0.3\gamma_{m} - (0.5 - 0.3)\gamma_{w} + 0.5\gamma_{m} - 0.4\gamma_{a} \\ = (0.5 + 0.1 + 0.7 - 0.3 + 0.5)\gamma_{m} - (0.1 + 0.7 - 0.2)\gamma_{w} - 0.4\gamma_{a} = 1.5\gamma_{m} - 0.6\gamma_{w} - 0.4\gamma_{a} \\ = 1.5 \times 133335.97 - 0.6 \times 9807 - 0.4 \times 12 \\ = 194114.955 \text{ N/m}^{2} \\ = 194.11 \text{ kN/m}^{2}.$$

EXAMPLE 9

A typical differential manometer is attached to two sections at *A* and *B*, as in the figure. Calculate pressure difference between sections *A* and *B*. Densities of water and mercury 1000 kg/m^3 and 13570 kg/m^3 , respectively.

SOLUTION

Here, $p_G = p_C = p_D$ and $p_E = p_F$ So, following the principle in (2.31), $p_A - \rho_w \cdot g \cdot z - \rho_w \cdot g \cdot 0.6 + \rho_w \cdot g \cdot (0.6 + z) = p_B$ $\Rightarrow p_A - p_B = 0.6 \cdot \rho_m \cdot g - 0.6 \cdot \rho_w \cdot g$





 $= 0.6 (13570 - 1000) \times 9.807 \text{ N/m}^2$ $= 73964.39 \text{ N/m}^2 = 73.96 \text{ kN/m}^2$

EXAMPLE 10

Diameters of reservoir and tube in an inclined tube manometer are 50 mm and 5 mm respectively. What will be % error in measuring p_s , if reservoir deflection is neglected?

SOLUTION

Let the level of gange fluid deflects from b - b to c - c, by a height 'h'. From continuity of fluid,

$$\frac{\pi}{4} \times (50)^2 \times h = \frac{\pi}{4} \times (5)^2 \times R$$

 $\Rightarrow h = 0.01 R$

Now $p_{C-C} = p_D$

$$\Rightarrow \qquad p_s = \rho_g \cdot g (h + R \sin \theta) \\ = \rho_g \cdot g (0.01R + R/2) \\ = \rho_g \times g \times 0.51R$$

Again without considering deflection of gauge fluid

$$p'_{s} = p_{b-b} = \rho_{g} \times g \times R \sin\theta = \rho_{g} \times g \times 0.5 R$$

From (c) and (d) error in per cent

$$= \frac{p_s \times p'_s}{p_s} \times 100 = \frac{0.01}{0.51} \times 100 = 1.96$$

EXAMPLE 11

Calculate the pressure difference between pipes *A* and *B* for the inverted *U*-tube manometers. SOLUTION

From equilibrium of pressure,

$$p_{A} - 0.165 \rho_{w} \cdot g + 0.05\rho_{oil} \cdot g + (0.165 - 0.05 - 0.05) \rho_{w} \cdot g = p_{B}$$

$$\Rightarrow \qquad p_{A} - p_{B} = 0.165 \cdot \rho_{w} \cdot g - 0.065 \cdot \rho_{w} \cdot g - 0.05 \cdot \rho_{oil} \cdot g$$

$$= 0.1 \times 10^{3} \times 9.807 - 0.05 \times 0.9 \times 10^{3} \times 9.807$$

$$= 539.385 \text{ N/m}^{2}$$





(c)

(d)

Fluid Statics

EXAMPLE 12

A manometer attached to a tank containing three different fluids is shown in the figure. What will be the difference in the elevation of the mercury column in the manometer?

SOLUTION

The equation of pressure equilibrium will be,

 $p_{air} + p_{oil} + p_{water} + \rho_{w} \times g \times 1.0 - \rho_{m} \times g \times y = 0$ $\Rightarrow 30 \times 10^{3} + 0.82 \times 10^{3} \times 9.807 \times (5 - 2) + 10^{3} \times 9.807 \times (2 - 0) + 10^{3} \times 9.807 \times 1.0 = 13.596 \times 10^{3} \times 9.807 \times y$ $\Rightarrow (13.596 \times 9.807)y = 30 + 0.82 \times 9.807 \times 3 + 9.807 \times 3$ $\Rightarrow y = 0.6265 \text{ m.}$



EXAMPLE 13

A differential manometer is attached to two tanks, as in the figure. Calculate pressure difference between chambers A and B. Specific gravities of oil, carbon-tetrachloride and mercury are 0.89,1.59 and 13.596 respectively.

SOLUTION

From the equilibrium of pressure, $p_A + 0.89 \times 10^3 \times 9.807 \times 1.1 + 13.596 \times 10^3 \times 9.807 \times 0.3$ $-1.59 \times 10^3 \times 9.807 \times 0.8 - p_B = 0$ $p_A - p_B = 12474.504 - 9601.053 - 40000.792$ $= -37127.341 \text{ N/m}^2 = -37.127 \text{ kPa}$



The result shows that p_B is higher than p_A and the difference of pressure is 37.127 kPa.

EXAMPLE 14

Determine the pressure difference between points A and B. Specific gravities of benzene, kerosene and air are 0.88, 0.82 and 1.2×10^{-3} respectively.



SOLUTION

Considering pressure equilibrium

 $p_{A} + \gamma_{b} \times 0.2 - \gamma_{m} \times 0.08 - \gamma_{k} \times (0.4 - 0.08) + \gamma_{w} \times (0.4 - 0.14) - \gamma_{a} \times 0.09 = p_{B}$ $p_{A} - p_{B} = \gamma_{m} \times 0.08 + \gamma_{k} \times 0.32 + \gamma_{a} \times 0.09 - \gamma_{b} \times 0.2 - \gamma_{w} \times 0.26$ $= 13.596 \times 10^{3} \times 9.807 \times 0.08 + 0.82 \times 10^{3} \times 9.807 \times 0.32$ $+ 1.2 \times 10^{-3} \times 10^{3} \times 9.807 \times 0.09 - 0.88 \times 10^{3} \times 9.807 \times 0.20 - 10^{3} \times 9.807 \times 0.26$ $= 10^{3} \times 9.807 (1.08768 + 0.2624 + 1.08 \times 10^{-4} - 0.176 - 0.26)$ $= 8965.44 \text{ N/m}^{2} = 8.965 \text{ kPa}$

EXAMPLE 15

Referring to the figure, if $p_B - p_A = 99.0$ kPa, what must the height *H* be? Sp. gravity of Meriam red oil = 0.827.

SOLUTION



EXAMPLE 16

Neglecting air pressure, compute the difference in pressure in tanks A and B, while $d_1 = 300$ mm, $d_2 = 150$ mm, $d_3 = 460$ mm, $d_4 = 200$ mm.

SOLUTION

Equation of pressure equilibrium will be $p_{A} + \gamma_{w} \times d_{1} - \gamma_{m} \times d_{3} - \gamma_{m} \times d_{4} \sin 45^{\circ} = p_{B}$ $\Rightarrow \qquad p_{A} - p_{B} = \gamma_{m} \times 0.46 + \gamma_{m} \times 0.2 \times 1/\sqrt{2} - \gamma_{w} \times 0.3$ $= 10^{3} \times 9.807 (13.596 \times 0.46 + 13.596 \times 0.2 \times 1/\sqrt{2} - 0.3)$ $= 77249 \text{ N/m}^{2}$ $= 77.249 \text{ kN/m}^{2}$

EXAMPLE 17

For the setup shown in the figure, what is the absolute pressure in drum A at position a? Assume an atmospheric pressure of 101.3 kPa.

Fluid Statics



SOLUTION

From the equation of pressure equilibrium,

$$\Rightarrow$$

$$p_{\text{atm}} + \gamma_{w} \times (0.6 - 0.3) - \gamma_{m} \times 0.15 + \gamma_{0} \times (0.15 + 0.1) = p_{A}$$

$$p_{A} = 101.3 \times 10^{3} + 10^{3} \times 9.807 \times 0.3 - 13.596 \times 10^{3} \times 9.807 \times 0.15$$

$$+ 0.8 \times 10^{3} \times 9.807 \times 0.25$$

$$= 86.203 \text{ kPa}$$

EXAMPLE 18

The pressure and temperature of air at MSL are 101.3 kN/m^2 and 15°C respectively. Calculate the pressure and density at an altitude of 3 km, if the lapse rate is 7°C per km, and R = 287 J/kg K.

(0.1.5

SOLUTION

For any polytropic state,

$$\frac{p}{\rho^{n}} = \frac{p_{0}}{\rho_{0}^{n}}$$

$$\Rightarrow \qquad \frac{\rho RT}{\rho^{n}} = \frac{\rho_{0}RT_{0}}{\rho_{0}^{n}} \qquad [\because p = \rho RT \text{ and } p_{0} = \rho_{0}RT_{0}]$$

$$\Rightarrow \qquad \frac{T}{\rho^{n-1}} = \frac{T_{0}}{\rho_{0}^{n-1}}$$

$$\Rightarrow \qquad \frac{T}{T_{0}} = \left(\frac{\rho}{\rho_{0}}\right)^{n-1} \qquad (a)$$
Again,
$$\frac{p}{p_{0}} = \frac{\rho RT}{\rho_{0}RT_{0}} = \left(\frac{\rho}{\rho_{0}}\right) \left(\frac{T}{T_{0}}\right)$$

 $= \left(\frac{T}{T_0}\right)^{\frac{1}{n-1}} \cdot \left(\frac{T}{T_0}\right)$ $= \left(\frac{T}{T_0}\right)^{\frac{n}{n-1}}$

=

=

(b)

317

(d)

Change of pressure in a polytropic state, as given in equation (2.26)

$$\frac{p}{p_0} = \left[1 - \left(z - z_0\right) \frac{g}{RT_0} \left(\frac{n-1}{n}\right)\right]^{\frac{n}{n-1}}$$
(c)

n

$$\Rightarrow \qquad \left(\frac{T}{T_0}\right)^{\frac{n}{n-1}} = \left[1 - \left(z - z_0\right)\frac{g}{RT_0}\frac{n-1}{n}\right]^{\frac{n}{n-1}}$$
$$\Rightarrow \qquad \frac{T}{T_0} = \left[1 - \left(\frac{n-1}{n}\right)\frac{gz}{RT_0}\right]$$

 \Rightarrow

 \Rightarrow

 \Rightarrow

differentiating (d) with respect to 'z'

$$\frac{dT}{dz} = -\frac{g}{R} \left(\frac{n-1}{n} \right)$$
$$n = \left[1 - \frac{R}{g} \left(\frac{dT}{dz} \right) \right]^{-1}$$
(e)

So, the polytropic power index n for the present case will be

$$n = \left[1 - \frac{287}{9.807} \times \left(\frac{7}{1000}\right)\right]^{-1}$$
$$= 1.2576$$

Applying equation (c),

$$\frac{p}{p_0} = \left[1 - (3000 - 0) \frac{9.807}{287 \times (273 + 15)} \left(\frac{1.2576 - 1}{1.2576}\right)\right]^{\frac{1.2576}{1.2576 - 1}}$$
$$= 0.6910$$
$$p = 101.3 \times 0.6910$$

 $= 70.00 \text{ kN/m}^2$

Applying equation (2.27), temperature at 3000 m will be $T = T_{o} - \alpha z$

$$= \left[\left(15 - \frac{7}{1000} \times 3000 \right) + 273 \right]$$

From the equation of state, the density at 3000 m

$$\rho = \frac{p}{RT} = \frac{70 \times 10^3}{287 \times 267} = 0.913 \text{ kg/m}^3$$

Fluid Statics

EXAMPLE 19

The diameters of a small piston and a large piston of a hydraulic jack are 4 cm and 11 cm respectively. A force of 85 N is applied on small piston. Find the load lifted by large piston while the small one is 40 cm above the large one. Density of liquid 1000 kg/m³. F_a

SOLUTION

c/s area of small piston (a)
$$= \frac{\pi}{4} \times (0.04)^2 \text{ m}^2$$

large piston (A) $= \frac{\pi}{4} \times (0.11)^2 \text{ m}^2$

Force on small piston
$$(F_a) = 85 \text{ N}$$

So, pressure intensity at X - X

$$= \frac{F_a}{a} + \rho \times g \times 0.4$$

$$= \frac{85}{\frac{\pi}{4} \times (0.04)^2} + 1000 \times 9.807 \times 04 \text{ N/m}^2$$

$$= 71563.65 \text{ N/m}^2$$

... Force generated on large piston

$$= 71563.65 \times \frac{\pi}{4} \times (0.11)^2 N$$
$$= 680.09 N$$

EXAMPLE 20

Find out the value of h of an inverted U-tube manometer, as shown.

SOLUTION

Equation of pressure equilibrium

$$p_A - \gamma_A \times 0.3 - \gamma_m \times h + \gamma_B \times (h + 0.3) - p_B = 0 \qquad (a)$$

As both the pipe are in same datum, $p_A = p_B$

From (a),
$$\gamma_B \times (h+0.3) - \gamma_m \times h - \gamma_A \times 0.3 = 0$$

 $\Rightarrow \qquad (\gamma_B - \gamma_m) h = (\gamma_A - \gamma_B) \times 0.3$

h = 0.2 m

- \Rightarrow
- $(1.0 0.7) h = (1.2 1.0) \times 0.3$ \Rightarrow
- \Rightarrow



1

40 cm

W

EXAMPLE 21

A circular plate 3 m diameter is submerged in water as shown in the figure. It's greatest and least depths are below the surface being 2 m and 1 m respectively. Find total pressure on front face of the plate and position of centre of pressure.

SOLUTION

Area of the plate (A) $= \frac{\pi}{4} \times 3^2 = 7.0685 \text{ m}^2$ $\sin \theta = \frac{2 - 1}{3} = \frac{1}{3}$

Distance of centre of gravity from the free surface of water

$$\overline{h} = 1 + 1.5 \sin\theta = 1.5 \text{ m}$$

 \therefore Total pressure force on the front of the plate.

$$F = \gamma h \cdot A$$

= $(10^3 \times 9.807 \times 1.5 \times 7.0685)$ N
= 103981.17 N

If the distance of centre of pressure from free surface of water is h^* ,

$$h^* = \frac{I\sin^2\theta}{A\overline{h}} + \overline{h} = \frac{\frac{\pi}{64} \times 3^4 \times \left(\frac{1}{3}\right)^2}{\frac{\pi}{4} \times 3^2 \times 1.5} + 1.5$$
$$= 1.5416 \text{ m}$$

EXAMPLE 22

An inclined rectangular sluice gate AB, $1.2 \text{ m} \times 5 \text{ m}$, shown in the figure is installed to control the discharge of water. A is hinged. Determine the force normal to the gate applied at B to open it. Solution

Area of gate $(A) = (1.2 \times 5) \text{ m}^2 = 6 \text{ m}^2$ Depth of C.G. of the gate from free surface of water

$$h = 5 - BG \sin 45^{\circ}$$
$$= \left(5 - \frac{0.6}{\sqrt{2}}\right) \mathrm{m} = 4.576 \mathrm{m}$$

:. Total pressure force acting on the plate

$$F = \gamma A \overline{h}$$

= (10³ × 9.807 × 6 × 4.576) N = 269260.99 N





Depth of centre of pressure

/

$$h^* = \left(\frac{5.0 \times 1.2^3}{12} \times \frac{\sin^2 45^\circ}{6 \times 4.576}\right) + 4.576 = 4.589 \text{ m}$$

Here, $OH = \frac{h^*}{\sin 45^\circ} = 6.489 \text{ m}$, and $OB = \frac{5}{\sin 45^\circ} = 7.071 \text{ m}$ BH = OB - OH = (7.071 - 6.489) m = 0.582 m

:.

$$AH = AB - BH = (1.2 - 0.582) \text{ m} = 0.618 \text{ m}$$

Take moment of force about A,

$$P \times AB = F \times AH$$

$$P = F \times \frac{AH}{AB}$$

$$= \left(269260.99 \times \frac{0.618}{1.2}\right) N = 138669.41 N$$

EXAMPLE 23

A cubical tank has sides of 1.5 m. It contains water for the lower 0.6 m depth. Upper remaining part is filled with oil of sp. gr. 0.9. Calculate total pressure and position of centre of pressure for one vertical side of the tank.

SOLUTION

The pressure intensity variation diagram is shown in the figure



Pressure intensity

at
$$D(p_D) = \rho_{oil} \times g \times h_{oil}$$

 $= (0.9 \times 10^3 \times 9.807 \times 0.9) \text{ N/m}^2$
 $= 7943.67 \text{ N/m}^2$
at $B(p_B) = 7943.67 + \rho_w \times g \times h_w$
 $= (7943.67 + 10^3 \times 9.807 \times 0.6) \text{ N/m}^2$
 $= 13827.87 \text{ N/m}^2$
 $\therefore DE = 7493.67 \text{ N/m}^2 \text{ and } BC = 13,827.87 \text{ N/m}^2$
Now, $F_1 = (1/2 \times 0.9 \times 7943.67) \times 1.5 = 5361.98 \text{ N}$
 $h_1 = (2/3 \times 0.9) \text{ m} = 0.6 \text{ m below } A$
 $F_2 = (0.6 \times 7943.67) \times 1.5 = 7149.303 \text{ N}$
 $h_2 = (0.9 + 0.6/2) \text{ m} = 1.2 \text{ m below } A$
 $F_3 = 1/2 \times 0.6 \times (13827.87 - 7943.67) \times 1.5 = 2647.89 \text{ N}$

F E

$$h_3 = (0.9 + 2/3 \times 0.6) \text{ m} = 1.3 \text{ m below } A$$

Total Force $(F) = F_1 + F_2 + F_3 = 15,159.17 \text{ N}$

If F acts at a depth h^* from free surface of liquid, taking moment of all forces about A,

 $F \times h^* = F_1 \times h_1 + F_2 \times h_2 + F_3 \times h_3$

 \Rightarrow

$$h^* = \frac{(5361.98 \times 0.6) + (7149.303 \times 1.2) + (2647.89 \times 1.3)}{15,159.17} \text{ m}$$

=1.0052 m

EXAMPLE 24

The figure shows the c/s of a tank full of water under pressure. Length of the tank is 2 m. An empty cylinder lies along the length of the tank on one of its corner as shown. Find horizontal and vertical components of the force acting on the curved surface ABC of the cylinder.

SOLUTION

Gauge pressure indicated on the gauge

$$= 0.2 \text{ kgf/cm}^2 = (0.2 \times 9.807) \text{ N/cm}^2$$
$$= (0.2 \times 9.807) \times 10^4 \text{ N/m}^2$$

.: Equivalent pressure head

$$=\frac{0.2\times9.807\times10^4}{10^3\times9.807}$$
 m = 2 m

- : Equivalent free surface of water will be 2 m above tank top.
- :. Horizontal component of force
 - = $\gamma \bar{h} \times (\text{area projected on vertical plane})$

$$= 10^3 \times 9.807 \times (2 + 1.5/2) \times (1.5 \times 2)$$

= 80907.75 N

Vertical component of force

- = Weight of water supported by curved surface *ABC*
- = Weight of water in CODEABC
- = Weight of water in (*CODFBC AEFB*)

Weight of water in $CODFBC(W_1)$

= Weight of water in (COB + ODFBO)



(a)

$$= \left(\frac{\pi}{4} \times 1^2 \times \gamma\right) \times 2 + (\gamma \times 1 \times 2.5) \times N$$
$$= 64439.799 N$$
(b)

Weight *AEFB* $(W_2) = [\gamma \times \text{area of } (AEFG + AGBH - ABH)] \times 2 \text{ N}$ But area *ABH* = Area *AOB* - ΔAOH

$$= \pi R^{2} \times \frac{\theta}{360^{\circ}} - \frac{1}{2} \times 1 \sin \theta \times 1 \cos \theta$$
$$= \pi R^{2} \times \frac{30}{360^{\circ}} - \frac{1}{2} \times \sin 30^{\circ} \times \cos 30^{\circ}$$
$$\left[\because \theta = \sin^{-1} \left(\frac{0.5}{1.0} \right) = 0.5 \right]$$
$$= 0.04529 \text{ m}^{2}$$

From (c),

$$= 10^{3} \times 9.807[2.0 \times (1 - 1\cos 30^{\circ}) + 0.5 \times (1 - 1\cos 30^{\circ}) - 0.04529] \times 2 \text{ N}$$

= 5681.126 N

Substituting (b) and (c) in (a)

Vertical component of force

 W_{2}

= (64439.799 – 5681.126) N = 58758.673 N.

EXAMPLE 25

A dam has a parabolic shape $y = y_0 \left(\frac{x}{x_0}\right)^2$ as shown in figure, having $x_0 = 6$ m and $y_0 = 9$ m. The

fluid is water with density 1000 kg/m³. Compute horizontal, vertical and resultant thrust exerted by water per metre length of the dam.

SOLUTION

Substituting values, $y = y_0 \left(\frac{x}{x_0}\right)^2$

$$=9\left(\frac{x}{6}\right)^2 = \frac{x^2}{4}$$

or,
$$x^2 = 4y$$

Horizontal thrust by water ($\gamma A \overline{h}$)

$$F_x = \text{force on vertical surface } OB$$
$$= [10^3 \times 9.807 \times (9 \times 1) \times 9/2]$$
$$= 397183.5 \text{ N}$$



(c)

Vertical thrust by water

$$F_{y} = \left[10^{3} \times 9.807 \times \int_{0}^{9} x \, dy \times 1 \right] \text{N}$$
$$= \left[9807 \times \int_{0}^{9} 2\sqrt{y} \, dy \right] \text{N}$$
$$= 19614 \left[\frac{2}{3} y^{\frac{3}{2}} \right]_{0}^{9} \text{N}$$
$$= 353052 \text{ N}$$

Resultant thrust exerted by water

$$F = \sqrt{F_x^2 + F_y^2}$$

= $\sqrt{397183.5^2 + 353052^2}$ N = 531413.63 N
 $\theta = \tan^{-1}\left(\frac{F_y}{F_x}\right) = 41.63^\circ$

and

EXAMPLE 26

A float valve regulates the flow of oil of sp. gr. 0.8 into a cistern. The spherical float is 15 cm in diameter. *AOB* is a weightless link carrying the float at one end, and a valve at the other end which closes the pipe through which oil flows into cistern. The link is mounted in a frictionless hinge at *O* and $\angle AOB = 135^\circ$. *OA* = 20 cm and the distance between centre of the float and the hinge is 50 cm. When the flow is stopped, *AO* will be vertical. The valve is to be pressed on to the seat with a force of 9.81 N to completely stop the flow of oil into the cistern. It was observed that the flow of oil is stopped when the free surface of oil in the cistern is 35 cm below hinge. Determine weight of the float.

SOLUTION

Let weight of the float = W

When the oil flow stopped, position of float is shown in the figure. Let the centre of float is h below level of oil.

Now,
$$\sin 45^{\circ} = \frac{0.35 + h}{0.50}$$

⇒ $h = 0.00355 \text{ m}$
Volume of oil displaced $= 2/3\pi r^3 + \pi r^2 \times h$
 $= 2/3\pi (0.075)^3 + \pi (0.075)^2 \times 0.00355 \text{ m}^3$
 $= 9.463 \times 10^4 \text{ m}^3$
 \therefore Buoyant force $= (0.8 \times 10^3 \times 9.807 \times 9.463 \times 10^4) \text{ N}$
 $= 7.424 \text{ N}$

Fluid Statics

 $\therefore \text{ Net vertical force on float} = (7.424 - W) \text{ N}$ Take moment about hinge O $P \times 0.2 = (7.424 - W) \times 0.5 \cos 45^{\circ}$ $\Rightarrow \qquad 9.81 \times 0.2 = (7.424 - W) \times 0.5 \times \frac{1}{\sqrt{2}}$ $\Rightarrow \qquad W = 1.875 \text{ N}$ EXAMPLE 27

A rectangular pontoon is 5 m long, 3 m wide and 1.20 m high. The depth of immersion of the pontoon is 0.8 m in sea-water. If C.G. is 0.6 m above bottom of pontoon, determine metacentric height. The density of sea water = 1025 kg/m^3 .

SOLUTION

Here, $AG = 0.6 \text{ m}, AB = 1/2 \times 0.8 = 0.4 \text{ m}$

:. BG = AG - AB = (0.6 - 0.4) m = 0.2 m

Volume of the submerged part

$$\forall = 3 \times 0.8 \times 5 \text{ m}^3 = 12 \text{ m}^3$$

Moment of inertia of plan of the pontoon

$$I = \frac{5 \times 3^3}{12} \,\mathrm{m}^4 = 11.25 \,\mathrm{m}^4$$

:. Metacentric height

$$GM = \frac{I}{\forall} - BG$$
$$= \left(\frac{11.25}{12}\right) - 0.2$$
$$= 0.7375 \text{ m}$$

G 0.4m B 0.6m G B 0.6m C 1.2m 0.8m A C 5.0m Plan at Water Surface

EXERCISE

1. State and prove Pascal's Law of pressure at a point.

m

- 2. Prove that pressure varies exponentially with elevation for isothermal condition.
- 3. Derive Euler's equation of motion along a streamline for an ideal fluid.
- 4. Differentiate between (i) absolute and gauge pressure (ii) simple manometer and differential manometer (iii) manometer and mechanical gauge.
- 5. Explain the working principle of a micromanometer.
- 6. What is temperature lapse rate? Derive its expression.

- 7. Find out expression for the force acting and its point of action on submerged plane surface due to hydrostatic effect.
- 8. Explain how will you find the resultant pressure on a curved surface immersed in a liquid.
- 9. Define the term buoyancy and centre of buoyancy. Obtain the expression for buoyant force and its centre of action.
- 10. What is metacentre? Derive an expression for metacentric height in terms of centre of gravity and centre of buoyancy.
- 11. An open tank contains water up to a depth of 1.5 m and above it an oil of sp.gr. = 0.8 for a depth of 2 m. Find the pressure intensity at the interface of the two liquids and at bottom of the tank.

Ans. [15.7 kPa, 30.4 kPa]

12. Diameters of a small piston and a large piston of a hydraulic jack are 2 cm and 10 cm respectively. A force of 60 N is applied on the small piston. Find the load lifted by large piston when (i) pistons are at same level (ii) small piston is 20 cm above large piston. Liquid inside the jack is water.

Ans. [1500 N, 1520.5 N]

13. Determine gauge and absolute pressure at 2 m below free surface of water, while atmospheric pressure is 101.043 kPa.

Ans. [0.01962 MPa, 0.12066 MPa]

14. A simple manometer used to measure pressure of oil (sp. gr. = 0.8) flowing in a pipeline. Its right limb is open to atmosphere and left limb is connected to the pipe. Centre of the pipe is 9 cm below the level of mercury (sp. gr. =13.6) in the right limb. If difference of mercury level in the two limbs is 15 cm, determine absolute pressure of oil in the pipe in N/cm².

Ans. [12.058 N/cm²]

15. A *U*-tube differential manometer connects two pressure pipe *A* and *B*. Pipe *A* contains carbon tetrachloride having a specific gravity 1.594 under a pressure of 11.772 N/cm²and pipe *B* contains oil (sp. gr. = 0.8) under a pressure of 11.772 N/cm². The pipe *A* lies 2.5m above pipe *B*. Find the difference of pressure measured by mercury as fluid filling *U*-tube.

Ans. [31.36 cm of mercury]

16. An inverted differential manometer containing an oil of sp. gr. 0.9 is connected to find difference of pressures at two points of a pipe containing water. If manometer reading is 40 cm, find the difference of pressures.

Ans. [392.4 N/m²]

 Calculate the pressure at a height of 8000 m above MSL if atmosphere pressure is 101.3 kN/m² and temperature 15°C at MSL assuming (a) air is incompressible (b) pressure variation follows isothermal law (c) pressure variation follows adiabatic law.

Ans. [607.5 Pa, 31.5 Pa, 37.45 Pa]

[Note: Take density of air at $MSL = 1.285 \text{ kg/m}^3$ and neglect variation of gravitational acceleration with altitude.]

18. A micromanometer consist of two cylindrical bulbs A and B each 10 cm² c/s area, which are connected by a U-tube, with vertical limbs each of 0.25 cm² c/s area. A liquid of sp. gr. 1.2 is filled in A and another liquid of sp.gr. 0.9 is filled in B, the surface of separation being in the limb attached to B. Find the displacement of the surface of separation, when the pressure on the surface in B is greater than that in A by an amount equal to 1.5 cm head of water.

Ans. [4.26 cm]

19. What is the position of the centre of pressure for a vertical semi-circular plane submerged in homogeneous liquid with its diameter *d* at the free surface?

Ans. [on centreline at depth]

20. How thick is the layer of liquid mud of sp.gr. 1.6 at the bottom of the tank with water 7.5 m deep above it, if there is a pressure of 5 kg/cm² against the bottom of the tank?

Ans. [26.56 m]

21 A solid cylinder of diameter 5 m has a height of 5 m. Find the metacentric height of the cylinder, if sp.gr. of the material is 0.7 and floating in water with its axis vertical. State, whether the equilibrium is stable or unstable.

Ans. [-0.304 m, unstable]

22. A rectangular pontoon 8 m long, 7 m wide and 3 m deep weighs 588.6 kN. It carries on its upper deck an empty boiler of 4 m diameter weighing 392.4 kN. The C.G. of the boiler and the pontoon are at their respective centres along a vertical line. Find the metacentric height. Weight density of sea-water is 10104 N/m³.

Ans. [0.325 m]

This page intentionally left blank



KINEMATICS OF FLUID FLOW

3.1 INTRODUCTION

The kinematics of fluid flow deals with only the space-time behaviour and associated velocity and acceleration, but without taking the effect of different forces associated in it. This chapter provides some definitions and types of fluid flow and the ways to characterise them in space-time frame and also the equation of continuity. Understanding of fluid kinematics is a pre-requisite to achieve a better understanding of dynamics of flow governed by the force criterion.

3.2 SCALAR AND VECTOR FIELD

From the name itself scalar is a parameter, which can be *scaled* or measured. Scalar quantity has only magnitude, no direction or sense. Mass, density, temperature, etc., are scalar quantities. If at every point in a region, scalar quantities, a scalar function has a defined value, the region is defined as a *scalar field*.

Vector means a *carrier*, which carries, in the present context, some physical quantity from one place to another. So a vector quantity is specified by its magnitude, direction and sense. Force, velocity and displacement are such vector quantities. If a vector function at every point in a region has a definite value, the region is called a *vector field*.

3.3 DESCRIPTION OF FLUID FLOW

A flow field is region, where at every instant of time, flow is defined at each and every point. Usually a flow field is described by velocity, a vector quantity, defined at each point, bearing same or different magnitude, direction and sense. Normally, the entire picture of fluid motion in any flow field are described by two approaches, one after *Lagrange* and another after *Leonhard Euler*.

3.3.1 Lagrangian Approach

Here motion of a fluid is described by the kinematic behaviour of individual fluid particle. A fluid particle can be described, at any instant of time, by their spatial coordinates. So in a 3-D rectangular Cartesian frame, it can be expressed as,

$$x = x (x_0, y_{0,}z_0, t) y = y (x_0, y_{0,}z_0, t) z = z (x_0, y_{0,}z_0, t)$$
(3.1)

Here x_0, y_0, z_0 are initial coordinates and *x*, *y*, *z* are coordinates at *t*-th instant of time. Hence a position vector of particle can be defined as,

$$\vec{S} = \vec{i} \ x + \vec{j} \ y + \vec{k} \ z \tag{3.2}$$

Here \vec{S} is the position vector and $\vec{i}, \vec{j}, \vec{k}$ are unit vectors along *x*-axis, *y*-axis and *z*-axis respectively. The velocity and acceleration of the particle along three axes, at any instant of time will be:

$$u = \left(\frac{dx}{dt}\right)_{x_0, y_0, z_0} v = \left(\frac{dy}{dt}\right)_{x_0, y_0, z_0} w = \left(\frac{dz}{dt}\right)_{x_0, y_0, z_0}$$
 (3.3*a*)
$$a_x = \left(\frac{d^2 x}{dt^2}\right)_{x_0, y_0, z_0} a_z = \left(\frac{d^2 z}{dt^2}\right)_{x_0, y_0, z_0}$$
 (3.3*b*)

3.3.2 Eulerian Approach

This approach considers velocity vector \vec{V} and its variation at every location (\vec{S}) of fluid flow, instead of considering movement of individual fluid particle. The flow field in this approach is expressed as,

$$\vec{V} = V\left(\vec{S}, t\right) \tag{3.4}$$

where, $\vec{V} = \vec{i} u + \vec{j} v + \vec{k} w$ and $\vec{S} = \vec{i} x + \vec{j} y + \vec{k} z$

$$\begin{array}{l} u = u\left(x, y, z, t\right) \\ \text{So,} \qquad v = v\left(x, y, z, t\right) \\ w = w\left(x, y, z, t\right) \end{array}$$
(3.5)

3.4 CLASSIFICATION OF FLOW

The velocity of flow and other hydrodynamic parameters like pressure, density, etc., are variable from one point to other point or from one instant of time to other instant of time. Depending on these variations, fluid flow can broadly be classified as:

- (i) Laminar and Turbulent flow
- (ii) Steady and Unsteady flow
- (iii) Uniform and Non-uniform flow
- (iv) Rotational and Irrotational flow
- (v) One, Two and Three dimensional flow

3.4.1 Laminar and Turbulent Flow

Laminar flow is one where fluid particles move along a smooth path in layers, with one layer (lamina) sliding smoothly over an adjacent layer. Laminar flow generally occurs when velocity of flow is small or viscosity is high. The velocity distribution in closed conduit is parabolic.

To interpret physically, a *turbulent flow* is a laminar flow on which small but random fluctuating motion is superimposed, as shown in figure 3.1. These random fluctuations, called turbulence, may vary from point to point and from time to time. Time average velocity of turbulent flow is called *temporal mean velocity*. Here shear stress depends on both viscosity and distribution of turbulence.



Fig. 3.1

Experimentations on Reynold's apparatus show (Fig. 3.2) that the path of dye in laminar flow is almost a straight line, while in turbulent flow, wavy, broken, discontinuous lines are found. A dimensionless number, *Reynold's Number* characterises whether a flow is laminar or turbulent.



Fig. 3.2

3.4.2 Steady and Unsteady Flow

A flow is said to be *steady*, if physical properties, such as velocity, pressure or density do not change with time, at any point. That is,

$$\frac{\partial V}{\partial t} = 0, \frac{\partial p}{\partial t} = 0, \frac{\partial \rho}{\partial t} = 0$$
(3.6)

If at any point, above characteristics, either one or some or all, change with time, then the flow will be defined as an *unsteady flow*. So,

$$\frac{\partial V}{\partial t} \neq 0, \frac{\partial p}{\partial t} \neq 0, \frac{\partial \rho}{\partial t} \neq 0$$
(3.7)

Tidal bore in a river is an example of natural unsteady flow.

To determine whether a flow is steady or unsteady, influence of random fluctuations associated with turbulence is generally neglected and only general fluid motion is taken into consideration. The steady and unsteady flow at a point for laminar and turbulent flow is shown in figure 3.3.





3.4.3 Uniform and Non-uniform Flow

Uniform flow is one, where at every point of the flow field, velocity is identical in magnitude and direction, *i.e.*, spatially invariant. So,

$$\frac{\partial V}{\partial S} = 0 \tag{3.8}$$

However uniform flow is not possible in a real fluid, in the strictness of definitions.

If the velocity of flow of the fluid changes from point to point at any instant, it is *non-uniform flow*. Mathematically,

$$\frac{\partial V}{\partial S} \neq 0 \tag{3.9}$$

All these four types of flow discussed till now can also combine to each other to generate (i) steady uniform flow (ii) steady non-uniform flow (iii) unsteady uniform flow (iv) unsteady non-uniform flow.

3.4.4 Rotational and Irrotational Flow

If the fluid particles, while moving in the direction of flow, rotate about their mass centre, the flow is defined as a *rotational flow*.

A flow will be considered *irrotational*, if fluid particles, during moving in the direction of flow, do not rotate about their mass centre. No irrotational flow can exist for a real fluid. For the irrotational flow,

$$\frac{\partial v}{\partial x} = \frac{\partial u}{\partial y} \tag{3.10}$$

3.4.5 Compressible and Incompressible Flow

A *compressible flow* is one, where density of the fluid is not constant and varies from point to point.

A flow will be considered *incompressible*, if the density of fluid is spatially invariant, *i.e.*, does not change from point to point.

3.4.6 One, Two and Three Dimensional Flow

One dimensional flow is defined as the flow, where flow properties like velocity, pressure and density vary in one direction only, *i.e.*, in longitudinal directions. So,

$$V = f(x, t)$$
 or $V = f(y, t)$ or $V = f(z, t)$ (3.11a)

In *two dimensional* flow, above flow parameters vary in two mutually perpendicular directions, *i.e.*,

$$V = f(x, y, t)$$
 or $V = f(y, z, t)$ or $V = f(x, z, t)$ (3.11b)

In *three dimensional* flow, variation of flow properties occur in all the three directions, generating a most complex type of motion to analyse.

3.5 DESCRIPTION OF FLOW PATTERNS

Fluid mechanics is a highly visual subject. The flow patterns can be visualised in many a different way, and sketches or photographs of these facilitates a lot to understand the flow, qualitatively and often quantitatively also. Some of the basic line patterns are described here.

3.5.1 Streamline, Stream Function and Stream Tube

A *streamline* is an imaginary line through the flow field, such that at every point of it, the tangent indicates the direction of the velocity vector at the point, as shown in figure 3.4. If u and v are the components of V along x and y-direction respectively, then,

$$\frac{v}{u} = \tan \theta = \frac{dy}{dx}$$

Here dy and dx are y and x-component of the differential displacement dS long a streamline, in the immediate vicinity of P. Therefore, the differential equation of the streamline will be

$$\frac{u}{dx} = \frac{v}{dy}$$
(3.12)





In three dimensional flow system, it will be,

$$\frac{u}{dx} = \frac{v}{dy} = \frac{w}{dz}.$$
(3.13)

The *stream function* ψ (psi) is defined as a scalar function of space and time, such that its partial derivative with respect to any direction gives the velocity component at right angles to this direction. So, $\psi = (x, y, t)$ can be defined such that,

$$u = \frac{\partial \Psi}{\partial y} \text{ and } v = -\frac{\partial \Psi}{\partial x}$$
 (3.14)

Since ψ is a point function, at every point of the flow field,

$$d\psi = \frac{\partial \psi}{\partial x} \cdot dx + \frac{\partial \psi}{\partial y} \cdot dy = -v \cdot dx = 0$$

 \Rightarrow

 ψ = Constant

So on a streamline, the value of stream function is constant.

The *streamtube* is an imaginary tube formed by a group of streamlines passing through a small closed curve, may or may not be circular. The cross-section of stream tube is finite and its surface across which there cannot be any flow is called a *stream surface*.

3.5.2 Pathline

A *pathline* is the trajectory of an individual fluid particle over a period of time. In a steady laminar flow, a pathline coincides with a streamline. In an unsteady flow, as a particle falls on a new streamline at every instant of time, pathline and streamline are not the same. Two pathlines can intersect each other or can form a loop. In capillary surface waves, pathlines are circular but streamlines are sinusoidal. The equation of pathline is:



Fig 3.5

$$x = \int u dt, \quad y = \int v dt, \quad z = \int w dt \tag{3.15}$$

3.5.3 Streakline

In flow visualisation, it is common practice to inject smoke into gas and dye into liquids. A *streakline* is a locus line at a given instant, connecting temporary locations of all fluid particles passing through a given point of injection. It is also called a *filament line*. In a steady laminar flow, the streamline, pathline and streakline are identical.





A timeline is a set of fluid particles that form a line at a given instant.

3.5.5 Equipotential Lines and Flow Net

A velocity potential ϕ (phi) is defined as scalar function of space and time, so that its negative partial derivative with respect to any direction gives fluid velocity in that direction.

So, velocity potential $\phi = f(x, y, t)$ is such that,

$$u = -\frac{\partial \phi}{\partial x}, \quad v = -\frac{\partial \phi}{\partial y}, \quad w = -\frac{\partial \phi}{\partial z}$$
 (3.16)

Equipotential line is a line for which velocity potential is constant at every point.

A *grid* obtained by drawing a series of streamlines and equipotential lines is defined a orthogonal flow net or simply the *flow net* (Fig. 3.7).

3.6 CONSERVATION OF MASS

The law of conservation of mass states that, *mass* can neither be created nor destroyed. Mathematically, for a closed system of mass m,

this law can be expressed as, $\frac{dm}{dt} = 0$. Hence

for a control volume, the law of conservation of mass can be re-instated as (Fig. 3.8),

Rate of influx – Rate of efflux = Rate of accumulation

This statement can be expressed analytically in terms of velocity and density field of a flow and resulting expression is known as *equation of continuity*.

3.6.1 Continuity Equation in Three Dimension

Let us consider a parallelepiped defining a differential control volume of mass $\rho dxdydz$. Components of velocity along x-axis, y-axis and z-axis are u, v and w respectively. Referring to figure 3.9,

Rate of mass influx through ABCD face = $\rho u dy dz$









dz

х

νZ

!w

Rate of mass efflux through EFGH face

$$=\rho \cdot u \cdot dy \cdot dz + \frac{\partial}{\partial x} \left(\rho \cdot u \cdot dy \cdot dz\right) dx$$

Rate of mass accumulation along x-direction

Similarly rate of mass accumulation along y-direction

$$= -\frac{\partial}{\partial y} (\rho u dx dy dz) \tag{3.17b}$$

and along z-direction = $-\frac{\partial}{\partial z}(\rho u dx dy dz)$ (3.17c)

So, total mass accumulation rate

$$= -\left[\frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) + \frac{\partial}{\partial z}(\rho w)\right] dx dy dz$$
(3.18)

Again rate of mass increase in control volume

$$=\frac{\partial}{\partial t}(\rho dx dy dz) = \frac{\partial \rho}{\partial t}(dx dy dz)$$
(3.19)

Equating (3.18) and (3.19) we have,

$$-\left[\frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) + \frac{\partial}{\partial z}(\rho w)\right] dxdydz = \frac{\partial \rho}{\partial t}(dxdydz)$$
$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) + \frac{\partial}{\partial z}(\rho w) = 0$$
(3.20)

 \Rightarrow

This is the most general form of equation of continuity in three dimensional Cartesian coordinate system. It is applicable for steady and unsteady flow, uniform and non-uniform flow, compressible and incompressible flow.

(i) For steady flow:
$$\frac{\partial \rho}{\partial t} = 0$$

Continuity equation: $\frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) + \frac{\partial}{\partial z}(\rho w) = 0$
(3.21)

(ii) For incompressible flow: $\rho = \text{constant}$

Continuity equation:
$$\frac{\partial u}{\partial x} + \frac{\partial u}{\partial x} + \frac{\partial u}{\partial x} = 0$$
 (3.22)

3.6.2 Continuity Equation in Two Dimension

In this case, velocity component along z-direction, w = 0. So generalised expression will be

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) = 0$$
(3.23)

(i) For steady flow:
$$\frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) = 0$$
 (3.24)

(ii) For incompressible flow:
$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$
 (3.25)

3.6.3 Continuity Equation in One Dimension

Here the velocity components both along *y*-direction and *z*-direction are zero, *i.e.*, v = 0 = w. So generalised expression is

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) = 0 \tag{3.26}$$

(i) For steady flow:
$$\frac{\partial}{\partial x}(\rho u) = 0$$
 (3.27)

(ii) For incompressible flow:
$$\frac{\partial u}{\partial x} = 0$$

 $\Rightarrow \qquad u = \text{constant}$
 $\Rightarrow \qquad A \times u = \text{constant}$
 $\Rightarrow \qquad Q = Au = \text{constant}$ (3.28)
Here A is c/s area of a conduit. So, through a

Here, A is c/s area of a conduit. So, through a closed conduit of uniform c/s, rate of flow is constant for a steady incompressible flow of fluid. This is equally true for conduit of variable c/s.



Fig. 3.10

A streamtube as shown in Fig. 3.10 is considered with c/s area of influx end is dA_1 and of efflux end is dA_2 . The principle of conservation of fluid mass can be expressed mathematically as, $\rho_1 dA_1 V_1 = \rho_2 dA_2 V_2$.

For incompressible fluid, $\rho_1 = \rho_2$. So, $dA_1 = dA_2 = volume \ rate \ of flow = dQ$. Hence, total volume rate of flow

$$Q = \int dQ = \int dA_1 V_1 = \int dA_2 V_2 = A_1 V_1 = A_2 V_2$$
(3.29)

3.6.4 Continuity Equation in Cylindrical Polar Coordinate

In this system of coordinate, a point in space is described by (r, θ, z) . Generalised continuity equation in space takes the shape as,

$$\frac{\partial \rho}{\partial t} + \frac{1}{r} \frac{\partial}{\partial r} (\rho V_r \cdot r) + \frac{1}{r} \frac{\partial}{\partial \theta} (\rho \cdot V_\theta) + \frac{\partial}{\partial z} (\rho \cdot V_z) = 0$$
(3.30)

(i) For steady flow:
$$\frac{1}{r}\frac{\partial}{\partial r}(\rho \cdot V_r \cdot r) + \frac{1}{r}\frac{\partial}{\partial \theta}(\rho \cdot V_{\theta}) + \frac{\partial}{\partial z}(\rho \cdot V_z) = 0$$
 (3.31)

(ii) For incompressible flow:
$$\frac{1}{r}\frac{\partial}{\partial r}(V_r \cdot r) + \frac{1}{r}\frac{\partial}{\partial \theta}(V_{\theta}) + \frac{\partial}{\partial z}(V_z) = 0$$
 (3.32)

3.6.5 Continuity Equation in Spherical Polar Coordinate

In this system, any point in space is described by (R, θ, ϕ) . Continuity equation in its generalised form can be expressed as,

$$\frac{\partial \rho}{\partial t} + \frac{1}{R^2} \frac{\partial}{\partial R} \left(R^2 \cdot \rho \cdot V_R \right) + \frac{1}{R \sin \phi} \frac{\partial}{\partial \theta} \left(\rho \cdot V_\theta \right) + \frac{1}{R \sin \phi} \frac{\partial}{\partial \theta} \left(\rho \cdot V_\theta \cdot \sin \phi \right) = 0$$
(3.33)

For a steady incompressible flow, it reduces to,

$$\frac{1}{R}\frac{\partial}{\partial R}\left(R^2 \cdot V_R\right) + \frac{1}{\sin\phi}\frac{\partial}{\partial\theta}\left(V_{\theta}\right) + \frac{1}{\sin\phi}\frac{\partial}{\partial\theta}\left(V_{\phi} \cdot \sin\phi\right) = 0$$
(3.34)

3.6.6 Velocity and Acceleration

Let us consider V as the resultant velocity at any point in a Eulerian flow field and u, v, w are its components along x-axis, y-axis and z-axis respectively. So,

$$u = f_{1}(x, y, z, t)$$

$$v = f_{2}(x, y, z, t)$$

$$w = f_{3}(x, y, z, t)$$

$$V = \sqrt{u^{2} + v^{2} + w^{2}} = \sqrt{\left(\frac{dx}{dt}\right)^{2} + \left(\frac{dy}{dt}\right)^{2} + \left(\frac{dz}{dt}\right)^{2}}$$
(3.35)

So,

Total or substantial acceleration *a* has components along *x*-axis, *y*-axis and *z*-axis as a_x , a_y , a_z respectively. Applying chain rule of differentiation, we find,

$$a_x = \frac{du}{dt} = \frac{\partial u}{\partial x}\frac{dx}{dt} + \frac{\partial u}{\partial y}\frac{dy}{dt} + \frac{\partial u}{\partial z}\frac{dz}{dt} + \frac{\partial u}{\partial t}$$

$$= u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} + \frac{\partial u}{\partial t}$$
$$= \frac{\partial u}{\partial t} + \left[u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right]$$
$$= Local Acceleration + Convective Acceleration (3.36)$$

Local acceleration is the rate of increase of velocity with respect to time at a given point in the flow field. It is also termed as *temporal acceleration*. For the x-axis, $\frac{\partial u}{\partial t}$ is local acceleration. Convective acceleration is the rate of change of velocity due to change of position of fluid

particles in a fluid flow. For the x-axis, $\left(u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} + w\frac{\partial u}{\partial z}\right)$ is the convective acceleration.

Following the same principle,

$$a_{y} = \frac{\partial v}{\partial t} + \left[u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} \right]$$
$$a_{z} = \frac{\partial w}{\partial t} + \left[u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} \right]$$

For steady flow, there exists no local acceleration, i.e.,

$$\frac{\partial u}{\partial t} = \frac{\partial v}{\partial t} = \frac{\partial w}{\partial t} = 0$$
(3.37)

In two dimensional flow, $a_x = \frac{\partial u}{\partial t} + \left[u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right]$ $a_{y} = \frac{\partial v}{\partial t} + \left[u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} \right] \left[\frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} \right]$

In one dimensional flow,

$$a_{x} = \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x}$$

$$a_{s} = \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial S}$$
(3.39)

 \Rightarrow

This is tangential total acceleration. The expression for normal acceleration,

$$a_n = \frac{\partial u_n}{\partial t} + \frac{u^2}{r}$$
 where, $u_n = u \cdot \frac{dS}{r}$

(3.38)

Features of acceleration of different types of flow are summarised as following:

- (i) Steady uniform flow: No acceleration
 - Steady non-uniform flow: Convective acceleration
 - Unsteady uniform flow: Local or temporal acceleration

(iv) Unsteady non-uniform flow: Both acceleration

In polar coordinate system,

$$\left. \begin{array}{l} a_r = V_r \cdot \frac{\partial V_r}{\partial r} + \frac{V_{\theta}}{r} \frac{\partial V_r}{\partial \theta} - \frac{V_{\theta}^2}{r} \\ a_{\theta} = V_r \cdot \frac{\partial V_{\theta}}{\partial r} + \frac{V_{\theta}}{r} \frac{\partial V_{\theta}}{\partial \theta} - \frac{V_r \cdot V_{\theta}}{r} \end{array} \right\}$$

where, V_r and V_{θ} are components of velocity along radial direction and azimuthal direction respectively.

Multiple Choice Questions

- 1. Steady flow is one in which
 - (a) velocity does not change from place to place
 - (b) velocity may change in direction but magnitude remains unchanged
 - (c) velocity at given point does not change with time
 - (d) mass flow rate remains constant
- 2. A flow is said to be uniform when
 - (a) at any given instant, velocity vector remains unchanged in magnitude and direction
 - (b) at any instant, velocity vector has the same directions at all points in the flow field
 - (c) at any instant, velocity vector may change its direction, but magnitude remains unchanged
 - (d) none of the above
- 3. The flow of river during the heavy rainfall is
 - (a) steady, uniform, 2-D flow
 - (c) unsteady, non-uniform, 3-D flow
- (b) unsteady, uniform, 3-D flow
- (d) none of the above
- 4. Continuity equation is based on the principle of
 - (a) conservation of energy(c) conservation of momentum
- (b) conservation of mass(d) Newton's law of viscosity
- 5. Laminar flow is characterised by
 - (a) irregular motion of fluid particle
 - (b) fluid particles moving in layers parallel to the boundary
 - (c) high Reynold's number of flow
 - (d) none of the above

(ii)

(iii)

- 6. A turbulent flow is considered steady, when
 - (a) discharge remains constant
 - (b) velocities at a point are time invariant
 - (c) temporal mean velocity at a point remains time invariant
 - (d) none of the above
- 7. One obtains a steady, uniform flow
 - (a) when ideal fluid passes through a long pipe at constant rate
 - (b) when real fluid passes through a long pipe at constant rate
 - (c) when ideal fluid passes through a converging duct at constant rate
 - (d) when real fluid passes through a converging duct at constant rate
- 8. Flow of water passing through a pipe is laminar, when
 - (a) fluid is ideal
 - (b) fluid is real
 - (c) velocity of flow is more than critical velocity
 - (d) Reynold's number is less than 2000
- 9. An ideal fluid flowing through a tapering pipe is an example of
 - (b) steady and non-uniform flow
 - (c) unsteady and uniform flow

(a) steady and uniform flow

(d) unsteady and non-uniform flow

- 10. A streamline is a line
 - (a) drawn normal to velocity vector at any point
 - (b) such that, streamline divides the passage into equal number of parts
 - (c) which is along the path of a particle
 - (d) tangent to which is in the direction of velocity vector at every point
- 11. Streamline, pathline and streakline are identical, when
 - (a) flow is uniform
 - (b) flow is steady
 - (c) flow velocities do not change steadily with time
 - (d) flow is neither steady nor uniform
- 12. The material acceleration zero for a $\frac{1}{2}$
 - (a) steady flow
 - (c) unsteady and uniform flow
- (b) steady and uniform flow
- (d) unsteady and non-uniform flow
- 13. Continuity equation in generalised form is

(a)
$$A_1V_1 = A_2V_2$$

- (c) $\rho_1 A_1 V_1 = \rho_2 A_2 V_2$
- 14. Irrotational flow means
 - (a) fluid does not rotate while moving
 - (b) fluid moves in a straight line
 - (c) net rotation of fluid particles about their mass center is zero
 - (d) none of the above

(b) $\rho_1 A_1 = \rho_2 A_2$ (d) $p_1 A_1 V_1 = p_2 A_2 V_2$

42			Mechanica
15.	If the fluid particles move in a zigzag way, th	e flow	is called
	(a) unsteady flow	(b)	turbulent flow
	(c) non-uniform flow	(d)	incompressible flow
16.	If the density of fluid changes from point to	point i	n a flow region, it is called
	(a) steady flow	(b)	unsteady flow
	(c) non-uniform flow	(d)	compressible flow
17.	If velocity in a fluid flow changes with respe	ect to le	ength of direction of flow, it is called
	(a) unsteady flow	(b)	compressible flow
	(c) irrotational flow	(d)	none of the above
18.	The acceleration of a fluid particle in the dire	ection	of x is given by
	$\partial u = \partial u = \partial v = \partial w = \partial u$	(b)	$u = \frac{1}{2} \partial v = \frac{1}{2} \partial w = \frac{1}{2} \partial u$
	(a) $u \frac{\partial x}{\partial x} + v \frac{\partial y}{\partial y} + w \frac{\partial z}{\partial z} + \frac{\partial t}{\partial t}$	(0)	$u\frac{\partial}{\partial x} + u\frac{\partial}{\partial y} + u\frac{\partial}{\partial z} + \frac{\partial}{\partial t}$
	(c) $u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} + w\frac{\partial u}{\partial z} + \frac{\partial u}{\partial t}$	(d)	none of the above
19.	The local acceleration of a fluid particle in the	ne dire	ction of y is given by
	(a) $\frac{\partial v}{\partial t}$	(b)	$u\frac{\partial u}{\partial x} + v\frac{\partial w}{\partial y}$
	(c) $v \frac{\partial v}{\partial x}$	(d)	none of the above
20.	The convective acceleration of a fluid partic	ele in tl	he direction of z is given by
	(a) $u\frac{\partial w}{\partial x} + v\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial z}$	(b)	$w\frac{\partial u}{\partial x} + w\frac{\partial v}{\partial y} + w\frac{\partial w}{\partial z}$
	(c) $u\frac{\partial w}{\partial z} + v\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial x}$	(d)	$w\frac{\partial w}{\partial x} + w\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial z}$
21.	A control volume refers to		
	(a) a fixed region in space	(b)	a specified mass
	(c) a closed system	(d)	a reversible process only
22.	One dimensional flow is		
	(a) steady uniform flow	(b)	uniform flow
	(c) flow which neglect changes in a transve	erse di	rection
	(d) restricted to flow in a straight line		
23.	A practical example of steady non-uniform f	low is	given by
	(a) motion of river around bridge piers		
	(b) steadily increasing flow through a pipe		

(c) constant discharge through a long, straight tapering pipe

(d) none of the above

	24. Specify, which of the following must be fulfilled by the flow of any fluid, real or ideal, lamin turbulent										nar or									
		(a)	Ne	wtor	n's lav	w of v	viscos	ity			(b)) .	velocity	at bo	undar	y mu	st be z	ero		
		(c) continuity equation							(d)) .	velocity	norm	al to a	a solic	l bour	dary	is zero)		
	25.	. The science dealing with the geometry of motion of fluids without reference to the motion is										orces	causi	ng the						
		(a)	hy	drost	tatics			(b) hydrokinetics												
		(c) hydrokinematics							(d) none of the above											
	Answers																			
											3401	3								
1.	(c)	, (d)	2.	(a)	3.	(c)	4.	(b)	5.	(b)	6.	(0	c) 7.	(a)	8.	(d)	9.	(b)	10.	(d)
11.	(b)		12.	(b)	13.	(c)	14.	(c)	15.	(b)	16.	(0	d) 17.	(d)	18.	(C)	19.	(a)	20.	(a)
21.	(a)		22.	(c)	23.	(c)	24.	(c)	25.	(c)										
Fill in the Blanks																				
	1.	Th	e flo	w of	wate	r in tł	ne ho	se of	a fire	-figh	ting pu	um	p is a				,			flow.
	2.	The flow of a water from a domestic tap is a														f	ow.			
	3.	Th	e flo	w fro	om a	rotati	ng lav	wn sp	rinkl	er is	a			,				f	low.	
	4.	Th	e flo	w of	wate	r ove	r a wi	de sp	oillwa	y in a	a river	is	a			,				flow.
	5.	. The flow of gas through the nozzle of a jet engine is a												flow.						
	6.	Th	e flo	w in	a riv	er dur	ing ti	dal b	ore is	an ez	xample	e o	f			_,				flow.
	_		~												~					

7. The flow of fluid through a capillary tube is a _____ flow.

8. In ______ flow, transfer of momentum is on a microscopic level.

9. In ______ flow, transfer of momentum is on molecular level.

10. A 2-D unsteady pressure field is represented by p=_____

Answers

1. uniform, steady2. non-uniform, steady3. non-uniform, steady4. non-uniform, steady5. non-uniform, steady6. non-uniform, unsteady7. laminar8. turbulent

9. laminar 10. p(x, y, t)

True or False

- 1. Pathlines cannot intersect either in steady or unsteady flow.
- 2. For a steady flow, mass flow rate remains constant between two streamlines.
- 3. Streamlines cannot start or end anywhere except at the interface or infinity.
- 4. Streamlines do not intersect, but converge at stagnation point, where velocity is zero.
- 5. Flow of water through a river bed may be uniform or may be non-uniform.
- 6. In turbulent flow, shear stress depends on viscosity and turbulence distribution.
- For a steady flow, tangential velocities do not vary in magnitude with the variation of gap between 7. streamline.
- 8. In a turbulent flow, a pathline coincide with a streamline.
- 9. The equation of continuity is valid both for incompressible and compressible flow.
- A flow net can be obtained associating a set of streamlines and a set of streaklines. 10.

					AllSweis				
1.	False	2.	True	3.	True	4.	True	5.	False
6.	True	7.	False	8.	False	9.	True	10.	False

NUMERICAL EXAMPLES

EXAMPLE 1

In a 1-D flow field, the velocity at a point may be given in the Eulerian system as u = x + t. Determine the displacement of fluid particle whose initial position is x_0 at initial time t_0 in Lagrangian system. SOLUTION

u = x + tHere, $\frac{dx}{dt} = x + t$ \Rightarrow

 $\frac{dx}{dt} - x = t$

The solution of this 1st degree differential equation $x = Ae^{t} - t - 1$

The value of constant A can be found out by applying boundary condition.

$$x_0 = Ae^{t_0} - t_0 - 1$$
$$A = \frac{x_0 + t_0 + 1}{e^{t_0}}$$

 \Rightarrow

So,

 \Rightarrow

$$x = \left(\frac{x_0 + t_0 + 1}{e^{t_0}}\right)e^t - t - 1 = (x_0 + t_0 + 1)e^{(t - t_0)} - t - 1.$$

EXAMPLE 2

A 2-D flow is described in Lagrangian system as $x = x_0 e^{-kt} + y_0 (1 - e^{-2kt})$, $y = y_0 e^{kt}$. Find (a) equation of pathline of the particle (b) velocity components in Eulerian system.

SOLUTION

 \Rightarrow

 \Rightarrow

(a) Pathline will be found out by eliminating t from these two equations. Substituting the value of y we have,

$$x = x_0 e^{-kt} + y_0 \left(1 - e^{-2kt}\right) = x_0 \left(\frac{y_0}{y}\right) + y_0 \left[1 - \left(\frac{y_0}{y}\right)^2\right]$$
$$y^2 x - x_0 y_0 y - y_0 y^2 + y_0^3 = 0$$
$$y^2 \left(x - y_0\right) - x_0 y_0 y + y_0^3 = 0$$

This is the equation of pathline.

и

(b) x-component of velocity

$$= \frac{dx}{dt} = \frac{d}{dt} \left[x_0 e^{-kt} + y_0 \left(1 - e^{-2kt} \right) \right]$$

$$= -kx_0 e^{-kt} + 2ky_0 e^{-2kt}$$

$$= -k \left[x - y_0 \left(1 - e^{-2kt} \right) \right] + 2ky_0 e^{-2kt}$$

$$= -kx + ky_0 - ky_0 e^{-2kt} + 2ky_0 e^{-2kt}$$

$$= -kx + ky_0 + ky_0 e^{-2kt}$$

$$= -kx + ky_0 \left(1 + e^{-2kt} \right)$$

$$= -kx + ky e^{-kt} \left(1 + e^{-2kt} \right)$$

$$= -kx + ky \left(e^{-kt} + e^{-3kt} \right)$$

$$dy = d \left(-kt \right) = -kt$$

y-component of velocity, $v = \frac{dy}{dt} = \frac{d}{dt} (y_0 e^{kt}) = y_0 k e^{kt} = ky.$

EXAMPLE 3

Given a velocity field $\vec{V} = (4 + xy + 2t)\vec{i} + 6x^3\vec{j} + (3xt^2 + z)\vec{k}$. Find the acceleration of a fluid particle (2, 4, -4) and time t = 3.

SOLUTION

For the given velocity field, components along *x*, *y* and *z* axes are:

$$u = 4 + xy + 2t$$
$$v = 6x^{3}$$
$$w = 3xt^{2} + z$$

So,

So,

$$a_{x} = \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z}$$

$$= 2 + 4y + xy^{2} + 2ty + 6x^{4}$$

$$a_{y} = \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z}$$

$$= 72x^{2} + 18x^{3}y + 36x^{2}t$$

$$a_{z} = \frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z}$$

$$= 6xt + 12t^{2} + 3xyt^{2} + 6t^{3} + 3xt^{2} + z$$
Substituting

$$x = 2, y = 4, z = -4, t = 3 \text{ we have}$$

$$a_{x} = 170 \text{ units}, a_{y} = 572 \text{ units}, a_{z} = 1296 \text{ units}$$
and

$$a = \sqrt{a_{x}^{2} + a_{y}^{2} + a_{z}^{2}} = 1426.78 \text{ units}$$

EXAMPLE 4

and

Find the acceleration component at a point (1,1,1) for the following flow field $u = 2x^2 + 3y$,

$$v = -2xy + 3y^2 + 3zy$$
 and $w = \frac{3}{2}z^2 + 2xz - 9y^2z$.

SOLUTION

Components of acceleration

$$a_{x}|_{(1,1,1)} = \left(\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} + w\frac{\partial u}{\partial z}\right)_{(1,1,1)} = 32 \text{ units}$$

$$a_{y}|_{(1,1,1)} = \left(\frac{\partial v}{\partial t} + u\frac{\partial v}{\partial x} + v\frac{\partial v}{\partial y} + w\frac{\partial v}{\partial z}\right)_{(1,1,1)} = -7.5 \text{ units}$$

$$a_{z}|_{(1,1,1)} = \left(\frac{\partial w}{\partial t} + u\frac{\partial w}{\partial x} + v\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial z}\right)_{(1,1,1)} = 23 \text{ units}$$

EXAMPLE 5

Find the acceleration components along radial and azimuthal direction at a point r = 2a, $\theta = \pi/2$

for a 2-*D* flow field $V_r = -u\left(1 - \frac{a^2}{r^2}\right)\cos\theta$ and $V_{\theta} = u\left(1 + \frac{a^2}{r^2}\right)\sin\theta$.

SOLUTION

$$\frac{\partial V_r}{\partial \theta} = u \sin \theta \left(1 - \frac{a^2}{r^2} \right), \quad \frac{\partial V_{\theta}}{\partial \theta} = u \cos \theta \left(1 + \frac{a^2}{r^2} \right)$$
$$\frac{\partial V_r}{\partial r} = -\frac{2ua^2}{r^3} \cos \theta, \qquad \frac{\partial V_{\theta}}{\partial r} = -\frac{2ua^2}{r^3} \sin \theta$$

Hence, acceleration in radial direction

$$a_r \Big|_{(2a,\pi/2)} = \left(V_r \cdot \frac{\partial V_r}{\partial r} + \frac{V_{\theta}}{r} \frac{\partial V_r}{\partial \theta} - \frac{V_{\theta}^2}{r} \right)_{(2a,\pi/2)} = -\frac{5u^2}{16} \text{ units}$$

and that along azimuthal direction

$$a_{\theta}\big|_{(2a,\pi/2)} = \left(V_r \cdot \frac{\partial V_{\theta}}{\partial r} + \frac{V_{\theta}}{r} \frac{\partial V_{\theta}}{\partial \theta} - \frac{V_r \cdot V_{\theta}}{r}\right)_{(2a,\pi/2)} = 0 \text{ unit}$$

EXAMPLE 6

Given the velocity field $\vec{V} = 6x\,\vec{i} + (16y+10)\,\vec{j} + (20t^2)\,\vec{k}$ m/s. What is the pathline of a particle which is at (2,4,6) at time t = 2s? SOLUTION

Component of velocity in x-direction $u = \frac{dx}{dt} = 6x$

Integrating

 \Rightarrow

$$\ln x = 6t - 11.307$$

 $\int_{2}^{x} \frac{dx}{x} = 6 \int_{2}^{t} dt$

Similarly for y-direction and z-direction, we can have

 $\int_{4}^{y} \frac{dy}{16y + 10} = \int_{2}^{t} dt$ $\ln(16y + 10) = 16t - 27.695$

z

 \Rightarrow

 \Rightarrow

Solving for *t*, we have

$$\int_{6} dz = 20 \int_{2} t^{2} dt$$
$$z = \frac{20}{3} t^{3} - 47.33$$

 $t = \left[\frac{3}{20}(z + 47.33)\right]^{1/3}$

t

So, finally the equation of the pathline will be

$$\ln x + \ln (16y + 10) = 22t - 39 = 11.689(z + 47.33)^{1/3} - 39.$$

EXAMPLE 7

A velocity field is given by $u = 3y^2$, v = 2x, w = 0 in arbitrary units. Is this flow steady or unsteady? Is it 2-*D* or 3-*D*? At (2,1,0) compute (a) velocity (b) local acceleration (c) convective acceleration.

SOLUTION

This flow is steady, all equations here are time invariant. The flow is 2-D as everywhere.

(a) The velocity field is $\vec{V} = 3y^2 \vec{i} + 2x\vec{j}$

At (2,1,0), $\vec{V} = 3\vec{i} + 4\vec{j}$ and magnitude $=\sqrt{3^2 + 4^2} = 5$ units

(b) Local acceleration in two directions

$$\frac{\partial u}{\partial t} = 0, \quad \frac{\partial v}{\partial t} = 0$$

(c) Convective acceleration in two directions

$$\left(u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y}\right)_{2,1,0} = 12xy\Big|_{2,1,0} = 24 \text{ units}$$
$$\left(u\frac{\partial v}{\partial x} + v\frac{\partial v}{\partial y}\right)_{2,1,0} = 6y^2\Big|_{2,1,0} = 6 \text{ units}$$

EXAMPLE 8

The velocity vector for a field $\vec{V} = -x\vec{i} + 2y\vec{j} + (3-z)\vec{k}$. Find equation of the streamline passing

(1,1,2).

Here,

SOLUTION

$$u = -x, v = 2y, w = 3 - z$$

From the equation of streamline,

$$\frac{u}{dx} = \frac{v}{dy} = \frac{w}{dz}$$
$$\frac{dx}{-x} = \frac{dy}{2y} = \frac{dz}{3-z}$$

 \Rightarrow

Integrating within limits, $\int_{1}^{x} \frac{dx}{-x} = \int_{1}^{y} \frac{dy}{2y}$

$$\Rightarrow \qquad \qquad x = \frac{1}{\sqrt{y}}$$
$$\Rightarrow \qquad \qquad \int_{1}^{y} \frac{dy}{2y} = \int_{2}^{z} \frac{dz}{3-z}$$

$$\Rightarrow \qquad \frac{1}{\sqrt{y}} = 3 - z$$

Combining these two results, $x = \frac{1}{\sqrt{y}} = 3 - z$

EXAMPLE 9

Two velocity components are given as
$$u = \frac{-2xyz}{(x^2 + y^2)^2}$$
 and $w = \frac{y}{(x^2 + y^2)}$. Find the third

component such that they satisfy the continuity equation.

SOLUTION

For an incompressible flow, the equation of continuity is

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$
$$\frac{\partial u}{\partial x} = \frac{6x^2 yz - 2y^3 z}{\left(x^2 + y^2\right)^3} \text{ and } \frac{\partial w}{\partial z} = 0$$

Now,

 \Rightarrow

Substituting in the governing equation, we have

$$\frac{6x^2yz - 2y^3z}{\left(x^2 + y^2\right)^3} + \frac{\partial v}{\partial y} + 0 = 0$$
$$\partial v = \frac{-6x^2yz + 2y^3z}{\left(x^2 + y^2\right)^3} \partial y$$
$$v = \frac{z\left(x^2 - y^2\right)}{\left(x^2 + y^2\right)^2} + C$$

On integration we have,

EXAMPLE 10

The velocity components in a 2-*D* flow field for an incompressible fluid are expressed as $u = \frac{y^3}{3} + 2x - x^2y \text{ and } v = xy^2 - 2y - \frac{x^3}{3}$ Show that these functions represent a possible case of an irrotational flow.

SOLUTION

Here,

$$\frac{\partial u}{\partial x} = 2 - 2xy$$
 and $\frac{\partial v}{\partial y} = 2xy - 2$

Now,

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 2 - 2xy + 2xy - 2 = 0$$

EXAMPLE 11

A gas, following $p = \rho RT$, flows steadily in a horizontal pipe of constant diameter. Considering the flow isothermal and $p_1: p_2 = 8:9$, find the ratio of V_1 and V_2 .

SOLUTION

From the given condition, we can write $A_1 = A_2$ and $T_1 = T_2$ From the principle of conservation of mass,

 $=\frac{9}{8}$

$$\rho_1 A_1 V_1 = \rho_2 A_2 V_2$$

$$\rho_1 V_1 = \rho_2 V_2$$

[:: $A_1 = A_2$]

 \Rightarrow

 \Rightarrow

$$\Rightarrow \qquad \frac{p_1}{RT_1}V_1 = \frac{p_2}{RT_2}V_2$$
$$\Rightarrow \qquad p_1V_1 = p_2V_2$$

$$\Rightarrow \qquad \qquad \frac{V_1}{V_2} = \frac{p_2}{p_1}$$

[:: flow isothermal]

EXAMPLE 12

Referring to adjoining figure, find the discharge in pipe 1 and velocity of flow in pipe 3.

SOLUTION

Cross-sections of

Solution
Cross-sections of
pipe 1
$$(A_1) = \frac{\pi}{4} \times 0.3^2 = 0.07068 \text{ m}^3$$

pipe 2 $(A_2) = \frac{\pi}{4} \times 0.2^2 = 0.07068 \text{ m}^3$
pipe 3 $(A_3) = \frac{\pi}{4} \times 0.15^2 = 0.07068 \text{ m}^3$
 30 cm
 30 cm
 $V_1 = 2.5 \text{ m/sec}$
 $V_1 = 2.5 \text{ m/sec}$
 0
 15 cm

Discharge of pipe 1 $(Q_1) = A_1V_1 = 0.07068 \times 2.5 = 0.1767 \text{ m}^3/\text{s}$ From the concept of continuity,

$$Q_1 = Q_2 + Q_3$$

$$\Rightarrow \qquad 0.03141 \times 2 + 0.01767V_3 = 0.1767$$

$$\Rightarrow \qquad V_3 = 6.44 \text{ m/s}$$

EXAMPLE 13

In the adjoining figure, CD carries $1/3^{rd}$ of flow in AB. Find volume rate of flow in AB, velocity in BC, velocity in CD and diameter of CE.

SOLUTION

Volume rate of flow in AB

Again,

$$\Rightarrow \qquad \frac{\pi}{4} \times 1.5^2 V_{BC} = 3.393$$
$$\Rightarrow \qquad V_{BC} = 1.92 \text{ m}$$

But,

⇒	$\frac{\pi}{4} \times 0.8^2 V_{CD} = \frac{3.393}{3}$
\Rightarrow	$V_{CD} = 2.25 \text{ m/s}$
Now,	$Q_{BC} = Q_{CD} + Q_{CE}$
⇒	$3.393 = \frac{3.393}{3} + \frac{\pi}{4}Q_{CE}^2 \times 2.5$
\Rightarrow	$Q_{CE} = 1.0733 \mathrm{m}$

 $Q_{BC} = Q_{AB}$

 $V_{BC} = 1.92 \text{ m/s}$

 $Q_{CD} = \frac{1}{3}Q_{AB}$



EXAMPLE 14

A jet of water from a 25mm diameter nozzle is directed vertically upwards. Assuming that jet remains and neglecting any loss of energy what will be the diameter at a point 4.5 m above nozzle, if velocity with which jet leaves the nozzle is 12 m/s.

SOLUTION

If velocity at 4.5 m above the nozzle be V, using equation of motion,

 $Q_{AB} = \frac{\pi}{4} \times 1.2^2 \times 3 = 3.393 \text{ m}^3/\text{s}$

$$V^2 = 12^2 - 2 \times 9.807 \times 4.5$$

V = 7.465 m/s

 \Rightarrow

From equation of continuity,

 $\frac{\pi}{4} \times 0.025^2 \times 12 = \frac{\pi}{4} \times D^2 \times 7.465$ $D = 0.031696 \,\mathrm{m} = 31.696 \,\mathrm{mm}$



EXAMPLE 15

The velocity vector for a 2-*D* incompressible flow field is given by $\vec{V} = \left(\frac{x}{x^2 + y^2}\right)^{\vec{i}} + \left(\frac{y}{x^2 + y^2}\right)^{\vec{j}}$. State whether the flow is continuous or discontinuous.

SOLUTION

Here,

$$u = \frac{x}{x^2 + y^2}, v = \frac{y}{x^2 + y^2}$$

$$\partial u = 1 \qquad 2x^2$$

$$\frac{\partial x}{\partial x} = \frac{1}{(x^2 + y^2)} - \frac{1}{(x^2 + y^2)^2} - \frac{2y^2}{(x^2 + y^2)^2} = \frac{1}{(x^2 + y^2)} - \frac{2y^2}{(x^2 + y^2)^2}$$

Hence,

So, the flow field satisfies continuity. Hence the flow is continuous.

 $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$

EXAMPLE 16

A nurse is withdrawing blood from a patient, as shown in the figure. The piston is being withdrawn at speed of 6.35 mm/s. The piston allows air to move through its peripheral region of clearance with the glass at the cylinder @ 0.02 cm^3 /s. What is the average speed of blood flow in the needle? Choose the region just to the right of the piston to the tip of the needle as a control volume. Inside diameter of needle and tube is 0.5 mm and 5 mm respectively.



SOLUTION

Let the average speed of blood flow in the needle is V m/s. From conservation of mass, it can be said

Total blood collected = blood sucked through needle + blood collected through peripheral clearance.

$$\rho_b \times \frac{\pi}{4} \times 5^2 \times 6.35 = \rho_b \times \frac{\pi}{4} \times 0.5^2 V + \rho_b \times 0.0 \times 10^3$$
$$V = 555 \,\mathrm{mm/s} = 555 \,\mathrm{cm/s}$$

Example 17

Meriam red oil (sp.gr. = 0.827) flows through 750 mm diameter pipeline @40 kL/min. Compute volume flux, average velocity and mass flux.

SOLUTION

Volume flux	$Q = \frac{40,000 \times 10^3}{60 \times 10^6} = 0.667 \mathrm{m^3/s}$
Again,	$0.667 = AV = \frac{\pi}{4} \times 0.75^2 \times V$
\Rightarrow	$V = 1.509 \mathrm{m/s}$
Mass flux	$= \rho Q = 0.827 \times 10^3 \times 0.667 = 551.609 \text{kg/s}$
Example 18	

Water flows steadily at three sections, as shown in the figure, through a box. Diameters at sections 1, 2, 3 are 7.5 cm, 5.0 cm, 2.5 cm respectively. Flow rate in section 1 is 0.03 m^3 /s in and velocity in section 2 is 10 m/s. Compute average velocity and volume flux at section 3.

SOLUTION

Let us assume flow through section 3 is inward.

 $Q_1 + Q_3 = Q_2$

So,

 \Rightarrow

$$Q_3 = \frac{\pi}{4} \times 0.05^2 \times 10 - 0.03 = -0.0103 \text{ m}^3/\text{s}$$

Negative sign indicates that the assumed direction of flow through section 3 is outward and its magnitude 0.0103 m³/s. Now, average velocity at section 3.

$$=\frac{0.0103}{\frac{\pi}{4}\times0.025^2}=20.98\,\mathrm{m/s}$$

EXAMPLE 19

A diffuser consists of two parallel circular plates of 20 cm diameter 0.5 cm apart and connected to a 3 cm pipe as shown in the figure. If the streamline are assumed to be radial in the diffuser, what mean velocity in the pipe will correspond to an exit velocity of 0.5 m/s?

SOLUTION

From continuity equation,

$$Q_{\text{pipe}} = Q_{\text{plates}}$$

 $V = 2.2 \, \text{m/s}$

$$\frac{\pi}{4} \times 0.03^2 V = \pi \times 0.2 \times 0.005 \times 0.5$$

 \Rightarrow





EXAMPLE 20

A liquid is flowing at constant Q through a divergent pipe, as shown in the figure. Assume velocity axial and uniform at any section, find out the acceleration at the inlet and outlet of the pipe.

SOLUTION

Solution
Let the diameter is D_x at a distance x from D_1 .
Now, $D_x = D_1 + \frac{x}{l}(D_2 - D_1)$
The velocity at this section, $u = \frac{4Q}{\pi \left[D_1 + \frac{x}{l} (D_2 - D_1) \right]^2}$
And $\frac{\partial u}{\partial x} = \frac{-8Q}{\pi \left[D_1 + \frac{x}{l} (D_2 - D_1) \right]^3} \times \frac{D_2 - D_1}{l}$
So, convective acceleration $= u \frac{\partial u}{\partial x} = \frac{-32Q^2 (D_2 - D_1)}{\pi^2 l \left[D_1 + \frac{x}{l} (D_2 - D_1) \right]^5}$
Convective acceleration at the inlet (for $x = 0$) = $\frac{-32Q^2(D_2 - D_1)}{\pi^2 l D_1^5}$
Convective acceleration at outlet (for $x = l$) = $\frac{-32Q^2(D_2 - D_1)}{\pi^2 l D_2^5}$.
EXERCISE

- 1. Define (a) streamline (b) pathline (c) timeline (d) streakline.
- 2. Write the differences between (a) steady and unsteady flow (b) uniform and non-uniform flow (c) laminar and turbulent flow (d) rotational and irrotational flow.
- Define velocity potential function and stream function. What are the conditions for the flow to be 3. irrotational?
- 4. Obtain an expression for continuity equation for a two dimensional flow.
- 5. Explain local acceleration and convective acceleration.
- 6. The velocity vector in a fluid flow is given by $\vec{V} = 2x^3 \vec{i} 5x^2 y \vec{j} + 4t \vec{k}$. Find velocity and acceleration of a fluid particle at (1, 2, 3) at time t = 1.

Ans. [10.95 units, 16.12 units]

7. Calculate the unknown velocity components of the following so that they satisfy continuity equations: (a) (b)

(a)
$$u = 4x^2$$
, $v = 4xyz$ (b) $u = 4x^2 + 3xy$, $w = z^3 - 4xy - 2yz$

Ans. [(a)
$$w = -8xz - 2xz^2 + f(x, y)$$

(b)
$$v = -8xy - \frac{y^2}{2} + 3yz^2 + f(x, z)$$
]

8. If for a 2-*D* potential flow, the velocity potential is given by $\phi = 4x(3y - 4)$, determine the velocity at point (2, 3). Determine also the value of stream function Ψ at this point.

Ans. [40 units,
$$\psi = 6x^2 - 4\left(\frac{3}{2}\right)y^2 - 4y$$
, -18]

9. For the steady incompressible flow, are the following values of *u* and *v* are possible? (a) $u = 4xy + y^2$, v = 6xy + 3x (b) $u = 2x^2 + y^2$, v = -4xy.

Ans. [(a) No. (b) Yes]

10. Show that, $u = -\frac{2xyz}{\left(x^2 + y^2\right)^2}$, $v = \frac{\left(x^2 - y^2\right)z}{\left(x^2 + y^2\right)^2}$, $w = \frac{y}{x^2 + y^2}$ are the velocity components of a possible

fluid motion. Is this motion irrotational?

Ans. [Yes]

11. A pipeline 60 cm diameter bifurcates at a Y junction into two branches 40 cm and 30 cm in diameter. If the rate of flow in the main pipe is 1.5 m³/s and the mean velocity of flow in the 30 cm pipe is 7.5 m/s, determine the rate of flow in 40 cm pipe.

Ans. [0.97 m³/s]

This page intentionally left blank



DYNAMICS OF FLUID FLOW

4.1 INTRODUCTION

In the previous chapter, the aspects related to displacement, acceleration and velocity of a fluid flow have been elaborated. It did not take into account of force effect. Here aspect of forces in case of a fluid flow is taken into consideration in dealing flow of a fluid. So this discussion is sometimes designated as *kinetics of fluid flow*. The dynamic behaviour of fluid flow is analysed using *Newton's* 2nd *law of motion*.

Here various types of forces coming into play are forces due to gravity (F_g) , pressure (F_p) , viscosity (F_v) , turbulence effect (F_r) , elasticity (F_F) etc. Hence net force (F) can be written as,

$$F = F_{G} + F_{P} + F_{V} + F_{T} + F_{E}$$
(4.1)

Considering only x-direction, the equation of motion is

$$F_{x} = ma_{x} = (F_{G} + F_{P} + F_{V} + F_{T} + F_{F})_{x}$$
(4.2)

For *y* and *z*-direction also, similar expressions can be found.

Here as change of volume being small, the effects of elasticity are neglected, and hence *Reynold's equation* is obtained as

$$F_{x} = ma_{x} = (F_{G} + F_{P} + F_{V} + F_{T})_{x}$$
(4.3)

At low Reynold's number, effect of turbulence becomes negligible. Considering this in view, following equations in three mutually perpendicular directions are obtained as *Navier-Stokes* equation.

$$F_{x} = ma_{x} = (F_{G} + F_{P} + F_{V})_{x}$$

$$F_{y} = ma_{y} = (F_{G} + F_{P} + F_{V})_{y}$$

$$F_{z} = ma_{z} = (F_{G} + F_{P} + F_{V})_{z}$$
(4.4)

For an ideal fluid flow, the effect of viscosity can be neglected. Hence

$$F_{x} = ma_{x} = (F_{G} + F_{p})_{x}$$
(4.5)

This equation is known as Euler's equation of motion.

dz

Streamline

(p+dp) dA

 $\gamma dA dS = W$

Fig. 4.1

4.2 EQUATION OF STEADY MOTION ALONG STREAMLINE

Let us consider a frictionless steady flow along a streamline. A cylindrical differential element as shown in figure 4.1 of length dS and cross sectional area dA is assumed along streamline. Intensity of pressures on the faces of the element are p and (p + dp). So the pressure force acting on the element tending to accelerate fluid mass will be

$$p \cdot dA - (p + dp)dA = -dp \cdot dA \tag{4.6}$$

If γ be the unit weight of fluid mass, gravity force acting vertically downward will be ($\lambda \cdot dA \cdot dS$). Component of gravity force along direction of motion will be

$$= -(\lambda \cdot dA \cdot dS)\cos\theta = -(\lambda \cdot dA \cdot dS)\frac{dz}{dS} = -(\lambda \cdot dA \cdot dz)$$
(4.7)

Therefore, the net force acting on the element

$$F = -dp \cdot dA - \lambda \cdot da \cdot dZ \tag{4.8}$$

dS

pdA

The velocity being V along the direction of motion, acceleration for the steady flow will $be\left(V \cdot \frac{dV}{dS}\right)$. Applying Newton's law of motion, we have,

$$-dp \cdot dA - \lambda \cdot dA \cdot dZ = \frac{\kappa}{g} \cdot dA \cdot dS \cdot V \frac{dV}{dS}$$
$$\frac{dp}{p} + VdV + gdZ = 0$$
(4.9)

This equation is commonly referred as one-dimensional *Euler's equation for an ideal fluid*, as it was first derived by *Leonhard Euler* in about 1750. And

it is applicable irrespective of any fluid flow, whether compressible or incompressible.

For a real fluid, referring to figure 4.2, an additional term, force due to fluid friction will be effective. If τ be the shear stress and *r* be the radius of the fluid element, the frictional force along the direction of motion will be $(-\tau \cdot 2 \cdot \pi r dS)$. Modifying equation (4.9), we arrive at



 $-dp \cdot dA - \lambda \cdot dA \cdot dZ - 2\pi r \cdot \tau \cdot dS = \frac{\gamma}{g} dA \cdot dS \cdot V \frac{dV}{dS}$ Fig. 4.2 $\frac{dp}{\rho} + V dV + g dZ = -\frac{2\tau \cdot dS}{\rho r}$ (4.1)

This is Euler's equation for one-dimensional real fluid flow.

 \Rightarrow

 \Rightarrow

 \Rightarrow

4.3 BERNOULLI'S EQUATION

For the case of an incompressible *ideal* fluid where $\rho = \text{constant}$, equation (4.9) can be integrated as,

$$\int \frac{dp}{\rho} + \int V dV + \int g dZ = \text{Constant}$$

$$\Rightarrow \qquad \frac{p}{\rho} + \frac{V^2}{2} + gZ = \text{Constant}$$

$$\Rightarrow \qquad \frac{p}{\gamma} + \frac{V^2}{2g} + Z = \text{Constant}$$

$$(4.11)$$

If the fluid becomes *real and incompressible*, integration over two section 1 and 2 of total fluid element length *l* yields,

$$\int_{1}^{2} \frac{dp}{\rho} + \int_{1}^{2} V dV + \int_{1}^{2} g dZ = -\int_{1}^{2} \frac{2\tau \cdot dS}{\rho r}$$

$$\frac{p_{2} - p_{1}}{\rho} + \frac{V_{2}^{2} - V_{1}^{2}}{2} + g(Z_{2} - Z_{1}) = -\frac{2\tau}{\rho r}(l_{2} - l_{1})$$

$$\frac{p_{2} - p_{1}}{\gamma} + \frac{V_{2}^{2} - V_{1}^{2}}{2g} + (Z_{2} - Z_{1}) = -\frac{2\tau}{\gamma r}l \qquad (4.12)$$

$$\Rightarrow \qquad \left(\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + z_1\right) - \left(\frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2\right) = \frac{2\tau}{\gamma r}l$$

Applying equation (4.11) over two section 1 and 2, we have,

$$\left(\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1\right) - \left(\frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2\right) = 0$$
(4.13)

- -

Hence right-hand side of equation (4.12), is designated as

$$h_f = \frac{2\tau \cdot l}{\gamma \cdot r} \tag{4.14}$$

This is the *head-loss* between two pre-selected sections in a conduit carrying *real fluid*. Equations (4.12) and (4.13) are popularly known as *Bernoulli's equation* for real and ideal fluid respectively, as an honour to *Daniel Bernoulli*, who first propounded it, in the year 1738. Now the *assumptions* for Bernoulli's equation for ideal fluid can be listed as:

- (i) Flow is steady
- (ii) Flow is incompressible for Mach number less than 0.3.
- (iii) Flow is frictionless or inviscid.

- (iv) Flow is along a streamline.
- (v) No shaft work or transfer of heat between two pres-selected sections.

4.4 **DIFFERENT HEADS**

In the equation (4.11), each term on the left-hand side has the dimension of length. Thus it can be termed

$$\frac{p}{\gamma}$$
 = Pressure Head, $\frac{V^2}{2g}$ = Velocity Head, Z = Datum Head and sum of these three heads is

denoted by *total head* (H). So for an incompressible ideal fluid, $H_1 = H_2$ and for incompressible real fluid, $H_1 - H_2 = h_f$. Again,

Pressure head
$$= \frac{p}{\gamma} = \frac{p \cdot A \cdot l}{\gamma \cdot A \cdot l} =$$
 pressure energy per unit weight
Velocity head $= \frac{V^2}{2g} = \frac{V^2 \cdot m}{2g \cdot m} = \frac{1/2mV^2}{mg} =$ kinetic energy per unit weight

Datum head = $Z = \frac{mgZ}{mg}$ = potential energy per unit weight

So, total head at any section of a conduit carrying incompressible fluid is H = [pressure energy + K.E. + P.E.] per unit weight.



Fig. 4.3

4.5 HYDRAULIC GRADE LINE (HGL) AND ENERGY GRADE LINE (EGL)

Summation of pressure head and datum head at each section for a conduit of incompressible flow, if joined by an imaginary line is called hydraulic grade line (HGL) or piezometric head line (PHL) or static head (SH).

The locus of the total head along the length of the conduit is *energy grade line* (EGL). Sometimes it is also called total head line (THL). Illustration is given in Fig. 4.3.

4.6 MAJOR AND MINOR HEAD LOSS

The head loss in a conduit carrying incompressible fluid, as shown in equation (4.14) is termed as *major head loss*. This is the frictional head loss resulting from wall shear stress within the conduit. For a circular pipe flowing full, the expression for major head loss is

$$h_f = f \cdot \frac{L}{D} \cdot \frac{V^2}{2g} \tag{4.15}$$

Here L is the length and D is the diameter of the conduit and V is the velocity of flow. The factor f is known as *friction factor*. This equation is popularly known as *Darcy-Weisbach equation*. In a slightly different form, if D is replaced by hydraulic radius R_h , equation (4.15) becomes

$$h_f = f \cdot \frac{L}{R_h} \cdot \frac{V^2}{2g} \tag{4.16}$$

For a circular conduit flowing full, hydraulic radius $R_{h} = D/4$ and hence,

$$h_f = f \cdot \frac{4L}{D} \cdot \frac{V^2}{2g} = 4f \cdot \frac{L}{D} \cdot \frac{V^2}{2g} = C_f \cdot \frac{L}{D} \cdot \frac{V^2}{2g}$$
(4.17)

This is Fanning's equation and C_f is known as coefficient of friction.

Some gas and water pipelines include pipe bends, valves, *T*-joints and regions where cross sectional area of the pipe changes abruptly. Local frictional loss at these points, including entrance and exit arise out of the effect of separation and are known as *minor losses*. It is expressed as

$$h_f = K \cdot \frac{V^2}{2g} \tag{4.18}$$

Here K is a constant bearing different value for commercial fittings, entrance and exit.

4.7 ABSOLUTE AND RELATIVE ROUGHNESS

The internal wall of the conduit or open channel cannot be perfectly smooth. Some protrusions evenly or unevenly spaced and of variable heights are noticed inside the wall. Average height of these protrusions is defined as *absolute roughness*. It is symbolised by ε (epsilon) and expressed in the unit of length.

When absolute roughness is divided by the diameter of the pipe, the obtained quantity is called *relative roughness* (ε_{r}). Relative roughness $\varepsilon_{r} = \varepsilon/D$ is a *pure number*.

4.8 REYNOLD'S NUMBER

In a fluid flow through a completely filled conduit, gravity force and capillarity have very negligible effect on flow pattern. Only the significant forces acting are *inertia force* and *fluid friction* due to viscosity. The ratio of inertia force to viscous force is called *Reynold's number* (R_e), in the honour of *Osborne Reynold*, who presented this through his experimental work in 1882. So,

$$R_e = \frac{\rho V D}{\mu} = \frac{V D}{\upsilon} \tag{4.19}$$

Reynold's number is a quantity which determines whether the flow is laminar or turbulent. If it is less than 2000, the flow is laminar. Turbulent flow starts at 4000. In between laminar and turbulent flow, transition flow exists where 2000 < Re < 4000.

Both the friction factor and coefficient of friction are dependent on Reynold's number (R_e) and relative roughness (ε_r) as $f = \phi_1(R_e, \varepsilon_r)$ and $C_f = \phi_2(R_e, \varepsilon_r)$. For a hydraulically smooth pipe, the effect of relative roughness become negligible and friction factor depends only on Reynold's number.

4.9 APPLICATION OF BERNOULLI'S EQUATION

Bernoulli's equation has wide variety of applications in the solution of a number of fluid flow problems. This equation along with the equation of continuity serves very effectively for measurement of flow through pipes. Some flow measuring devices, otherwise known as flow-meters are discussed here.

4.9.1 Orificemeter

An *orifice* is an opening (usually circular) in the wall of a tank or in a plate normal to the axis of the pipe, the plate being either at the end or at any intermediate location of the pipe. A standard sharp-edged orifice is shown in figure 4.4. Thickness of the orifice will be small enough, so that cannot affect the flow of fluid through it. While the streamlines converge towards approaching an orifice, they continue to converge beyond upstream of the orifice, until reaching at section x-y where they become parallel. This section x-y bears a minimum area or diameter of flow and the section is defined as *vena contracta*. Beyond vena contracta, streamlines diverge.





The performance of an orifice is related with three constants, namely, *coefficient of contraction* (C_{a}) , *coefficient of velocity* (C_{a}) , *coefficient of discharge* (C_{a}) .

Coefficient of contraction is the ratio of the area of a jet at vena contracta (A_c) to the area of orifice (A_o) .

So,
$$C_c = \frac{Area \ at \ vena \ contracta}{Area \ of \ orifice} = \frac{A_C}{A_O}$$
(4.20a)

Ideal velocity is one that would be attained in the jet at frictionless condition. *Coefficient* of velocity is the ratio of actual average velocity at vena contracta (V_c) to ideal velocity (V_i) .

So,
$$C_{v} = \frac{Actual \ average \ velocity}{Ideal \ velocity} = \frac{V_{C}}{V_{i}}$$
(4.20b)

Coefficient of discharge is the ratio of actual rate of discharge (Q) to ideal rate of discharge (Q_i) at no-friction and no-contraction condition.

So,
$$C_d = \frac{Actual \ rate \ of \ discharge}{Ideal \ discharge} = \frac{Q}{Q_i}$$
 (4.20c)

The mutual relationship amongst these three coefficients can easily be established as

$$C_d = \frac{Q}{Q_i} = \frac{A_C \times V_C}{A_O \times V_i} = \frac{A_C}{A_O} \times \frac{V_C}{V_i} = C_c \times C_v \tag{4.21}$$

363

An orifice or an orifice-plate, when used in a pipeline, is called an orificemeter, which provides a simple and cheap arrangement for measurement of flow through a pipe. An orificemeter fitted along a horizontal pipe line is shown in Fig. 4.5, where section 1 is assumed at upstream side and section 2 at vena contracta. Applying Bernoulli's equation at these two section yields,



Fig. 4.5

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} = \frac{p_2}{\gamma} + \frac{V_2^2}{2g}$$
$$\frac{p_1}{\gamma} - \frac{p_2}{\gamma} = \frac{V_2^2}{2g} - \frac{V_1^2}{2g}$$
(4.22)

If h be the piezometric head difference at these two sections, above equation can be written as,

$$h = \frac{V_2^2}{2g} - \frac{V_1^2}{2g}$$

$$V_2 = \sqrt{2gh + V_1^2}$$
(4.23)
at section 2. Now the actual velocity at this section will be

 \Rightarrow

 \Rightarrow

This is the ideal velocity at section 2. Now the actual velocity at this section will be

$$V_{2} = C_{v} \sqrt{2gh + V_{1}^{2}} .$$

$$A_{1}V_{1} = A_{2}V_{2} = C_{c} \cdot A_{O} \cdot V_{2}$$

$$V_{1} = C_{c} \cdot \frac{A_{O}}{A_{1}} \cdot V_{2}$$
(4.24)

Again,

Hence,

 \Rightarrow

 \Rightarrow

 \Rightarrow

 $V_{2}^{2} = C_{v}^{2} \left[2gh + \left(C_{c} \cdot \frac{A_{O}}{A_{1}} \cdot V_{2}\right)^{2} \right]$ $V_{2}^{2} \left[1 - C_{v}^{2}C_{c}^{2} \left(\frac{A_{O}}{A_{1}}\right)^{2} \right] = C_{v}^{2} \cdot 2gh$ $V_{2}^{2} \left[1 - C_{d}^{2} \left(\frac{A_{O}}{A_{1}}\right)^{2} \right] = C_{v}^{2} \cdot 2gh$ $\left[\because C_{d} = C_{v} \times C_{c} \right]$ $V_{2} = C_{v} \sqrt{\frac{2gh}{1 - C_{d}^{2} \left(\frac{A_{O}}{A_{1}}\right)^{2}}}$ (4.25)

$$Q = A_2 V_2 = C_c \cdot A_0 \cdot V_2 = C_c \cdot A_0 \cdot C_v \sqrt{\frac{2gh}{1 - C_d^2 \left(\frac{A_0}{A_1}\right)^2}}$$
$$= C_d \cdot A_0 \cdot \sqrt{\frac{2gh}{1 - C_d^2 \left(\frac{A_0}{A_1}\right)^2}}$$
(4.26)

4.9.2 Venturimeter

A venturimeter is a device which is used to measure both compressible as well as incompressible flow rate through a conduit. $\frac{V^2}{2g}$ *G.B. Venturi*, the famous Italian physicist investigated its principle in about 1791 and it was applied to measurement by *Clemens Herschel* in 1886. As shown in figure 4.6, a venturimeter attached along a horizontal pipe attached along a horizontal pipe consists of a tube having constricted *throat* producing increased velocity and reduced pressure, followed by a diverging part. Applying Bernoulli's equation in section 1 and 2,

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} = \frac{p_2}{\gamma} + \frac{V_2^2}{2g}$$



$$\Rightarrow \qquad \frac{V_2^2}{2g} - \frac{V_1^2}{2g} = \frac{p_1}{\gamma} - \frac{p_2}{\gamma} = h \qquad (4.27)$$

Here h is the piezometric head difference. Now ideal discharge

$$Q_{i} = A_{1}V_{1} = A_{2}V_{2}$$

$$V_{1} = \frac{Q_{i}}{A_{1}}, V_{2} = \frac{Q_{i}}{A_{2}}$$
(4.28)

Hence,

 \Rightarrow

 \Rightarrow

 \Rightarrow

$$h = \frac{V_2^2}{2g} - \frac{V_1^2}{2g} = \frac{Q_i^2}{2g} \left(\frac{1}{A_2^2} - \frac{1}{A_1^2} \right)$$
$$Q_i = A_1 A_2 \sqrt{\frac{2gh}{A_1^2 - A_2^2}}$$
(4.29)

So, actual discharge $Q = C_d \cdot Q_i$

$$=C_d A_1 A_2 \sqrt{\frac{2gh}{A_1^2 - A_2^2}}$$
(4.30)

If a *U*-tube manometer is employed to measure the difference in pressure head at section 1 and 2, then for a definite level difference (say x) of manometric fluid in the two limbs, it is

$$\frac{p_1}{\gamma} - \frac{p_2}{\gamma} = x \left(\frac{\gamma_m}{\gamma} - 1 \right)$$

$$h = x \left(\frac{\gamma_m}{\gamma} - 1 \right)$$
(4.31)

If an inverted *U*-tube manometer is used, then $h = x \left(1 - \frac{\gamma_m}{\gamma} \right)$ (4.32)

4.10 STATIC PRESSURE AND STAGNATION PRESSURE

The static pressure is a parameter necessary for description of the state of a flowing fluid. The hydrostatic pressure created by collision of molecules is known as *static pressure* in a fluid flow. When the fluid is at rest, this pressure is the same in all the directions.

While a fluid is flowing past an immersed body, there exists a particular streamline which strikes a points on the body, almost vertically. As a result, resultant velocity at that point reduces to zero, converting the entire kinetic energy of fluid to the increased pressure. This point on the body

is defined as *stagnation point* or point of stagnation (Fig. 4.7). The values of pressure, density and temperature at stagnation point are defined as *stagnation pressure* (p_s) , *stagnation density* (ρ_s) and *stagnation temperature* (T_s) . Stagnation pressure is possible to occur only in the case of *reversible adiabatic process* or *isentropic process*.

From the definition itself, it is evident that stagnation pressure is the algebraic summation of static pressure (p_{a}) and dynamic pressure, i.e.,

$$\frac{p_s}{\gamma} = \frac{p_o}{\gamma} + \frac{V^2}{2g} \tag{4.33}$$

4.11 PITOT TUBE

It is a device used for measuring the velocity of flow at any point in a pipe or an open channel. The principle of measurement of flow using *pitot-tube* was first adopted by *Henri Pitot*, a French scientist in 1732, during velocity measurement in River Seine.

In its simplest form, a pitot-tube is a right-angled glass tube, long enough to neglect the effect of capillarity, shown in Fig. 4.8. One end of the tube is directed to face the flow while its other end is open to atmosphere. If the rise along fluid be h, then employing the concept of stagnation pressure,

$$\frac{p_s}{\gamma} - \frac{p_o}{\gamma} = \frac{V^2}{2g} = h$$

$$\Rightarrow \qquad V = \sqrt{2gh} \qquad (4.34)$$

This is the theoretical velocity through the pitot-tube. Actual velocity of flow can be obtained by introducing the coefficient of pitot-tube, as

$$V_{ac} = C_p \sqrt{2gh} \tag{4.35}$$

The value of C_p lies between 0.97 and 1.00, depending on the degree of turbulence.

Another arrangement of pitot-tube is shown in Fig. 4.9, where piezometer at section a measures static pressure and pitot-tube at section b measures total pressure or stagnation pressure.



Streamlines







Fig 4.9

The arrangement of measurement of static pressure and total pressure when are combined into one instrument, it is called *pitot-static tube*. A particular type of pitot-tube is *Prandtl's tube* or *Prandtl's static tube*. Schematic views of normal pitot-static tube and Prandtl's tube are given in Fig. 4.10 and 4.11.



Multiple Choice Questions

- 1. Bernoulli's equation is obtained by the integration of
 - (a) Navier-Stokes equation (b) Reynold's equation
 - (c) Euler's equation (d) none of these
- 2. Bernoulli's equation written in the form $\frac{p}{\gamma} + \frac{V^2}{2g} + z = \text{constant}$, represents total energy per unit

(b) mass

- (a) volume
- (c) weight (d) none of these

3.	Pie	zometric head is the summation of		
	(a)	velocity head and datum head	(b)	pressure head and velocity head
	(c)	velocity head in the direction of flow	(d)	none of these
4.	HG	L indicates the variation of		
	(a)	total energy in direction of flow	(b)	piezometric head in direction of flow
	(c)	velocity head in direction of flow	(d)	pressure head in direction of flow
5.	EG	L always lies		
	(a)	above	(b)	below
	(c)	parallel to	(d)	along HGL
6.	EG	L in pipe flow is a graphical representation	ofB	ernoulli's equation and measured above
	(a)	centerline of pipe	(b)	top of the pipe
	(c)	arbitrary horizontal datum	(d)	top of piezometer tube
7.	Tot	al energy represented by Bernoulli's equa	tion	has the unit
	(a)	kgm/m	(b)	Nm/N
	(c)	kgm ² /s	(d)	kgm/s
8.	Sta	gnation point is a point where		
	(a)	pressure is zero	(b)	total energy is zero
	(c)	velocity of flow reduces to zero	(d)	total energy is maximum
9.	Pito	ot-tube is an instrument for measuring		
	(a)	pressure of flow	(b)	velocity of flow
	(c)	discharge of fluid	(d)	total energy
10.	Sta	gnation pressure is the sum of		
	(a)	vacuum pressure and static pressure	(b)	static pressure and dynamic pressure
	(c)	dynamic pressure and vacuum pressure	(d)	none of these
11.	Ver	nturimeter is used to measure		
	(a)	piezometric head	(b)	velocity head
	(c)	discharge through a pipe	(d)	none of these
12.	Av	renturimeter is preferable to orificemeter, be	ecaus	e
	(a)	it is cheaper	(b) i	it is more convenient to use
	(c)	energy loss is less	(d)	it requires smaller length
13.	C_d	of an orificemeter is always		
	(a)	greater than C_c	(b)	equal to C_c
	(c)	equal to C_{v}	(d)	less than C_c
14.	Wh	ich of the following devices is used for me	easuri	ng pipe flow?
	(a)	mouthpiece	(b)	notch
	(c)	weir	(d)	orificemeter

- 15. In orifice flow, vena contracta represents
 - (a) jet has maximum flow area
 - (b) pressure is above atmosphere
 - (c) jet area is minimum and streamline are parallel
 - (d) pressure is below atmosphere
- 16. Reynold's number is the ratio of
 - (a) inertia force and gravity force (b) inertia force and viscous force
 - (c) pressure force and inertia force (d) none of these
- 17. Value of Reynold's number below which flow through closed conduit will be laminar is
 - (a) 5,000 (b) 20,000
 - (c) 2,000 (b) 3,000
- 18. Reynold's number is expressed as

(a)
$$\frac{\rho VD}{\mu}$$
 (b) $\frac{V^2 D}{\rho}$
(c) $\frac{V \rho^2 D}{\upsilon}$ (d) $\frac{V^2 D^2}{\upsilon}$

- 19. Energy loss in pipeline is due to
 - (a) viscous action only
 - (c) pipe wall friction and viscous action
- (b) surface roughness only
- (d) turbulent shear stress alone
- 20. Minor loss in pipe flow are those
 - (a) which have small magnitude
 - (b) which are caused by local disturbances caused by valves, bends etc.
 - (c) caused by friction
 - (d) none of these.
- 21. The unit of relative roughness is
 - (a) m (b) N
 - (d) dimensionless, pure number (c) kg
- 22. The parameters on which the friction factor is dependent are
 - (a) relative roughness and Reynold's number (b) absolute roughness and viscosity

(c) Reynold's number and specific gravity (d) none of these

23.	Coe	efficient of discharge (C_d) in terms of C_c and	dC_v	are
	(a)	$C_d = C_c \times C_v$	(b)	$C_d = C_c / C_v$

(c)	$C_d = C_v / C_c$	(d)	none of these
-----	-------------------	-----	---------------

- 24. Bernoulli's equation can be applied to
 - (a) venturimeter (b) pitot-tube
 - (c) orificemeter (d) all of these

- 25. For hydrodynamiclly smooth boundaries, friction factor for turbulent flow is
 - (a) dependent on Reynold's number
 - (b) constant
 - (c) function of Reynold's number and relative roughness
 - (d) dependent on roughness only
- 26. The losses in fluid flow is more in
 - (a) laminar flow
 - (c) transition flow
- 27. HGL is
 - (a) always above centerline of pipe
 - (c) never above EGL
- 28. Bernoulli's equation is based on
 - (a) principle of conservation of energy
 - (c) principle of conservation of momentum (d) none of these
- 29. Euler's equation of motion can be integrated when it is assumed that
 - (a) continuity equation is satisfied
 - (b) fluid is incompressible
 - (c) velocity potential exists and density is constant
 - (d) flow is irrotational and incompressible
- 30. Specify, which of the following must be fulfilled by the flow of any fluid, real or ideal, laminar or turbulent
 - (a) Newton's law of viscosity (b) continuity equation
 - (c) velocity at boundary must be zero (d) velocity normal to solid boundary is zero
- 31. Head loss in turbulent flow of pipe
 - (a) varies directly as the velocity
 - (b) varies directly as the square of the velocity
 - (c) varies inversely as the square of the velocity
 - (d) varies approximately as the square of the velocity
- 32. In laminar flow through a closed conduit, discharge varies
 - (a) linearly as the viscosity (b) as the square of the radius
 - (c) inversely as the pressure drop (d) inversely as the viscosity
- 33. For laminar flow in a pipe of circular cross section Darcy friction factor is
 - (a) independent of Reynold's number and independent of pipe wall friction
 - (b) inversely proportional to Reynold's number and independent of pipe wall roughness
 - (c) directly proportional to Reynold's number and independent of pipe wall roughness
 - (d) inversely proportional to Reynold's number and directly proportional to pipe wall roughness

(b) always sloping downward in sloping direction

- (b) principle of conservation of mass

(b) turbulent flow

none of these

(d) critical flow

(d)

- (a) equation is valid for steady flow of an incompressible ideal or real fluid along a streamline
- (b) energy equation for the flow of a frictionless fluid of constant density along a streamline with gravity as the only body force
- (c) equation is derived from dynamic consideration involving gravity, viscous and inertia forces
- (d) constant in the equation varies across streamlines if flow is irrotational
- 35. Euler's equation of motion is a
 - (a) statement of energy balance
 - (b) preliminary step to derive Bernoulli's equation
 - (c) statement of conservation of momentum for a real, incompressible fluid
 - (d) statement of conservation of momentum for an inviscid fluid flow
- 36. Study of fluid motion with forces causing flow is known as
 - (a) kinematics of fluid flow (b) kinetics of fluid flow
 - (c) statics of fluid flow (d) none of the above
- 37. The term $V^2/2g$ in Bernoulli's equation is known as
 - (a) kinetic energy (b) pressure energy
 - (c) kinetic energy per unit weight (d) none of these
- 38. Difference of pressure head (h) measured by mercury-oil differential manometer is given by

(a)
$$h = x \left(1 - \frac{S_g}{S_o} \right)$$
 (b) $h = x \left(S_g - S_o \right)$

(c)
$$h = x \left(S_o - S_g \right)$$
 (d) $h = x \left(\frac{S_g}{S_o} - 1 \right)$

- 39. Ratio of area of jet at vena-contracta to area of orifice is known as
 - (a) coefficient of discharge

(a) coefficient of discharge

- (b) coefficient of velocity
- (c) coefficient of contraction (d) coefficient of viscosity
- 40. Ratio of actual velocity at vena-contracta to ideal velocity
 - (b) coefficient of velocity
 - (c) coefficient of contraction (d) coefficient of viscosity

Answers

1.	(c)	2.	(c)	3.	(C)	4.	(b)	5.	(a)	6.	(c)	7.	(b)	8.	(c)	9.	(b)	10.	(b)
11.	(c)	12.	(c)	13.	(d)	14.	(d)	15.	(c)	16.	(b)	17.	(c)	18.	(a)	19.	(c)	20.	(b)
21.	(d)	22.	(a)	23.	(a)	24.	(d)	25.	(a)	26.	(b)	27.	(c)	28.	(a)	29.	(c)	30.	(b)
31.	(d)	32.	(d)	33.	(b)	34.	(a)	35.	(a)	36.	(b)	37.	(c)	38.	(d)	39.	(c)	40.	(b)

120 mm

True or False

- 1. Friction factor in the Darcy-Weisbach equation is given by $f = \tau_0 / 0.5 \rho V^2$
- 2. Minor losses can be expressed either as fraction of kinetic energy or equivalent length.
- 3. HGL cannot have positive slope.
- 4. Concept of hydraulic radius is relevant to both laminar and turbulent flow.
- 5. In ageing of pipes, friction factor increases linearly with tie.
- 6. The term $V^2/2g$ in Bernoulli's equation is known as kinetic energy per unit weight.
- 7. Orificemeter can also be used for measuring flow of an open channel.
- 8. In orificemeter, coefficient of discharge is always less than unity.
- 9. At stagnation point, velocity of flow is maximum.
- 10. In a hydraulically smooth pipe, Reynold's number is not dependent on relative roughness.

				Α	nswers				
1.	False	2.	True	3.	False	4.	False	5.	False
6.	True	7.	False	8.	True	9.	False	10.	True

NUMERICAL EXAMPLES

EXAMPLE 1

A vertical venturimeter shown in figure, has an area ratio 5, with throat diameter 1 cm. When oil of specific gravity 0.8 flows through it, mercury in the differential gauge indicates a difference in height of 12 cm. Find discharge through venturimeter.

100 mm

SOLUTION

Difference in piezometric head

$$\left(\frac{p_1}{\gamma} + Z_1\right) - \left(\frac{p_2}{\gamma} + Z_2\right) = x \left(\frac{\gamma_m}{\gamma_{\text{oil}}} - 1\right)$$
$$h = 0.12 \left(\frac{13.596}{0.8} - 1\right) = 1.919 \text{ m}$$

 \Rightarrow

But, $A_1 = 5A_2$ and

$$A_2 = \frac{\pi}{4} \times 0.01^2 = 7.8539 \times 10^{-5} \text{ m}^2$$

So, actual discharge through venturimeter, from equation (4.30)

$$Q = C_d A_1 A_2 \sqrt{\frac{2gh}{A_1^2 - A_2^2}}$$

= 0.98 × 5 × (7.8539 × 10⁻⁵)² $\sqrt{\frac{2 \times 9.807 \times 1.919}{(5 \times 7.8539 \times 10^{-5})^2 - (7.8539 \times 10^{-5})^2}}$
= 4.819 × 10⁻⁴ m³/s

EXAMPLE 2

A pitot-static tube is used to measure velocity of water in a pipe. The stagnation pressure head is 8 m and static pressure head is 6 m. Calculate the velocity of flow, assuming coefficient of tube 0.98.

SOLUTION

Referring to equation (4.33)

$$\frac{V^2}{2g} = \frac{p_s}{\gamma} - \frac{p_o}{\gamma} = 8 - 6 = 2$$
$$\implies \qquad V = 6.263$$

Introducing discharge coefficient, actual velocity of flow = $0.98 \times 6.263 = 6.137$ m/s

EXAMPLE 3

A vertical pipe of 1 m diameter and 20 m length has a pressure head of 5.5 m of water at the upper end. When water flows through the pipe at an average velocity of 4.5 m/s, the head loss due to friction is 1.2 m. Find the pressure at the lower end of the pipe, when the flow is (a) upward (b) downward.

SOLUTION

(a) For upward flow; upper section is considered section 2-2 and lower section is 1-1

Here,
$$\frac{p_2}{\gamma} = 5.5 \text{ m}, V_1 = V_2 = 4.5 \text{ m/s}, h_f = 1.2 \text{ m}, (Z_2 - Z_1) = 20 \text{ m}$$

Applying Bernoulli's equation

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2 + h_f$$

 \Rightarrow

$$\frac{p_1}{\gamma} = \frac{p_2}{\gamma} + (Z_2 - Z_1) + h_f = 5.5 + 20 + 1.2 = 26.7 \text{ m}$$

(a) For downward flow; upper section is 1-1 and lower section is 2-2.

Here,
$$\frac{p_1}{\gamma} = 5.5 \text{ m}, V_1 = V_2 = 4.5 \text{ m/s}, h_f = 1.2 \text{ m}, (Z_1 - Z_2) = 20 \text{ m}$$

Applying Bernoulli's equation

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2 + h_f$$

⇒

$$\frac{p_2}{\gamma} = \frac{p_1}{\gamma} + (Z_1 - Z_2) - h_f = 5.5 + 20 - 1.2 = 24.3 \text{ m}$$

EXAMPLE 4

A siphon is shown in the figure. A is 1 m above water level, indicates point 1. Bottom of siphon is below A. Assuming friction to be negligible, determine velocity of flow at outlet and the pressure at A.

SOLUTION

Applying Bernoulli's equation at 1 and 2

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2$$

Here, $Z_1 = -1$ m, $Z_2 = -8$ m, atmospheric pressure $= \frac{p_1}{\gamma} = \frac{p_2}{\gamma} = 0$ and cross-sectional area of the tank is

much larger than of the tube, so, $V_1 \ll V_2$.

Hence,

here,
$$Z_1 = \frac{V_2^2}{2g} + Z_2$$

 $V_2 = 11.717 \text{ m/s} \approx V_A$ \Rightarrow Again, applying Bernoulli's equation at 1 and A

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_A}{\gamma} + \frac{V_A^2}{2g} + Z_1$$
$$Z_1 = -1 \text{ m}, Z_A = 0, \frac{p_1}{\gamma} = 0$$

 $Z_1 = \frac{p_A}{\gamma} + \frac{V_A^2}{2g}$

But,

Hence,

$$\Rightarrow \frac{p_A}{\gamma} = -7.999 \text{ m}$$

So, at A, the pressure is negative, *i.e.*, suction pressure.

EXAMPLE 5

Referring to the figure of siphon carrying oil (0.8), calculate (a) velocity of oil through siphon (b) pressure at A, B and C. (c) maximum height of C can be accommodated above the level in the vessel (d) maximum vertical depth of right limb of siphon. Take vapour pressure at working temperature 29.5 kPa and atmospheric pressure 101 kPa.

SOLUTION

(a) Applying Bernoulli's equation at 1 and D, where $Z_1 = 1.5 + 4 = 5.5 \text{ m}, Z_D = 0,$

atmospheric pressure =
$$\frac{p_1}{\gamma} = \frac{p_D}{\gamma} = 0$$
, and $V_1 \ll V_D$,

 $Z_1 = \frac{V_D^2}{2g}$

 $V_D = 10.386 \text{ m/s}$

we have





(b) As cross-section of siphon tube does not change, velocity of oil through siphon will be $V_B = V_D = 10.386$ m/s.

Here
$$\frac{p_A}{\gamma} = \frac{p_1}{\gamma} + 1.5 = 1.5 \text{ m of oil} (\text{gauge}) \text{ and } Z_A = 4 \text{ m} = Z_B \text{ and } V_A = 0$$
.

Applying Bernoulli's equation at A and B

$$\frac{p_A}{\gamma} = \frac{p_B}{\gamma} + \frac{V_B^2}{2g}$$
$$\frac{p_B}{\gamma} = 1.5 - \frac{10.386^2}{2 \times 9.807} = -4.0 \text{ m of oil(gauge)}$$

 \Rightarrow

So, absolute pressure at $B = 101 - 4 \times 0.8 \times 9.807 = 69.62$ kPa

Absolute pressure at $A = 101 + 1.5 \times 0.8 \times 9.807 = 112.768$ kPa

Applying Bernoulli's equation between 1 and C

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_C}{\gamma} + \frac{V_C^2}{2g} + Z_C$$
$$\frac{p_C}{\gamma} = (5.5 - 7) - \frac{10.386^2}{2 \times 9.807} = -7 \text{ m of oil(gauge)}$$

 \Rightarrow

Absolute pressure at $C = 101 - 7 \times 0.8 \times 9.807 = 46.08$ kPa

(c) At maximum height of C, the pressure will be equal with vapour pressure. Considering absolute pressure, Bernoulli's equation yields

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_{Vap}}{\gamma} + \frac{V_C^2}{2g} + Z_C'$$
$$\frac{101 \times 10^3}{0.8 \times 9.807 \times 10^3} + 5.5 = \frac{29.5 \times 10^3}{0.8 \times 9.807 \times 10^3} + \frac{10.386^2}{2 \times 9.807} + (5.5 + h)$$

 \Rightarrow

 \Rightarrow

$$h = 3.614 \text{ m}$$

(d) Let *H* be the value of maximum depth at which pressure at *C* equals to vapour pressure. Now applying Bernoulli's equation at 1 and *D*, $V_D = \sqrt{2gH}$

Again, Bernoulli's equation at 1 and C, considering absolute pressure,

$$\frac{p_1}{\gamma} + H = \frac{29.5 \times 10^3}{0.8 \times 9.807 \times 10^3} + \frac{V_D^2}{2g} + H + 1.5$$

$$H = \frac{101 \times 10^3}{0.8 \times 9.807 \times 10^3} - \frac{29.5 \times 10^3}{0.8 \times 9.807 \times 10^3} - 1.5 = 7.613 \text{ m}.$$

EXAMPLE 6

376

Water flows through a 300×150 venturimeter @0.037 m³/s and differential gauge is deflected 1 m. Specific gravity of the gauge liquid is 1.25. Determine coefficient of discharge of the meter.

SOLUTION

Applying Bernoulli's equation between A and B, where $Z_A = 0 = Z_B$ yields

$$\frac{p_A}{\gamma} + \frac{V_A^2}{2g} = \frac{p_B}{\gamma} + \frac{V_B^2}{2g}$$
(a)

Applying equation of continuity between A and B,

$$V_{A} = \left(\frac{A_{B}}{A_{A}}\right) V_{B}$$
$$V_{A}^{2} = \left(\frac{A_{B}}{A_{A}}\right)^{2} V_{B}^{2}$$
(b)

 \Rightarrow

Substitution (b) into (a), we obtain $V_B = \sqrt{\frac{2g(p_A - p_B)}{\gamma \left(1 - \frac{A_B^2}{A_A^2}\right)}}$

So, actual discharge
$$Q = C_d A_B V_B = C_d A_B \sqrt{\frac{2g(p_A - p_B)}{\gamma \left(1 - \frac{A_B^2}{A_A^2}\right)}}$$
 (c)

Again from differential gauge,

$$\frac{p_A}{\gamma} - Z - (1.25 \times 1) + (Z + 1) = \frac{p_B}{\gamma}$$

 \Rightarrow

$$\frac{p_A}{\gamma} - \frac{p_B}{\gamma} = 0.25$$

 $C_d = 0.9155$

Substituting in (c) we obtain,

$$0.037 = C_d A_B V_B = C_d \cdot \frac{\pi}{4} \ 0.15^2 \ \sqrt{\frac{2 \times 9.807 \times 0.25}{\left(1 - \frac{0.15^4}{0.3^4}\right)}}$$



EXAMPLE 7

A venture section of a pipe flow is shown in the figure. Derive an expression for V_{2^2} just sufficient to cause the reservoir liquid to rise in tube upto section 1.



SOLUTION

For the liquid to rise through the tube, $p_{atm} - p_1 \ge \gamma h$

Using continuity equation, $V_1 = \left(\frac{A_2}{A_1}\right)V_2 = \left(\frac{D_2}{D_1}\right)^2 V_2$

Applying Bernoulli's equation at 1 and 2, while $Z_1 = Z_2$

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} = \frac{p_{\text{atm}}}{\gamma} + \frac{V_2^2}{2g}$$

$$p_{\text{atm}} - p_1 = \gamma \frac{\left(V_1^2 - V_2^2\right)}{2g} = \gamma \frac{V_2^2}{2g} \left[\left(\frac{D_2}{D_1}\right)^4 - 1 \right]$$
(b)

 \Rightarrow

Comparing the two values, given in equations (a) and (b)

$$\gamma \frac{V_2^2}{2g} \left[\left(\frac{D_2}{D_1} \right)^4 - 1 \right] \ge \gamma h$$
$$V_2 = \sqrt{\frac{2gh}{\left[\left(\frac{D_2}{D_1} \right)^4 - 1 \right]}}$$

 \Rightarrow

EXAMPLE 8

Air flows through a duct and Pitot-static tube measuring the velocity is attached to a differential manometer containing water. If deflection of the manometer is 100 mm, calculate air velocity, assuming $\rho_{air} = 1.22 \text{ kg/m}^3$ and coefficient of tube be 0.98.

(a)

SOLUTION

Now,

$$\frac{P_s}{\gamma} - \frac{P_o}{\gamma} = h_{\text{water}} = \frac{0.1 \times 9.807 \times 10^3}{1.22 \times 9.807} = 81.967 \text{ m of air}$$

Applying equation (4.35)

$$V_{ac} = C_p \sqrt{2gh} = 0.98 \sqrt{2 \times 9.807 \times 81.967} = 39.294 \text{ m/s}$$

EXAMPLE 9

Water flows at a velocity of 1.417 m/s. A differential gauge which contains a liquid of specific gravity 1.25 is attached to a Pitot-static tube. What is the deflection of the gauge fluid? Coefficient of the tube is 1.

SOLUTION

Following equation (4.35)

 \Rightarrow \Rightarrow

$$V = C_p \sqrt{2gh}_w$$

1.417 = $1\sqrt{2 \times 9.807 \times h_w}$

 $h_w = 0.1023 \text{ m of water}$

γ

From manometric equation of differential gauge,

$$h_w = \left(\frac{\gamma_m}{\gamma_w} - 1\right) h_m$$
$$0.1023 = \left(\frac{1.25}{1} - 1\right) h_m$$
$$h_m = 0.4092 \text{ m}$$

 \Rightarrow

$$\Rightarrow$$

EXAMPLE 10

A vertical venturimeter carries a liquid of sp.gr. 0.8. Inlet and outlet diameters are 150 mm and 75 mm respectively. Pressure connection at throat is 150 mm above inlet. If actual flow rate is 40 lit/s and $C_d = 0.96$, calculate (a) pressure difference between inlet and throat (b) difference in mercury level in a U-tube manometer connected between inlet and throat.

SOLUTION

(a) Let inlet is section 1 and throat is section 2. Here,

$$A_{1} = \frac{\pi}{4} 0.15^{2} = 0.01767 \text{ m}^{2}, A_{2} = \frac{\pi}{4} 0.075^{2} = 4.4178 \times 10^{-3} \text{ m}^{2},$$
$$Q = 40 \times 10^{3} \times 10^{-6} \text{ m}^{3}/\text{s}, C_{d} = 0.96.$$

Substituting in $Q = C_d A_1 A_2 \sqrt{\frac{2gh}{A_1^2 - A_2^2}}$ we can obtain, h = 4.25 m

So, piezometric head difference

$$h = \left(\frac{p_1}{\gamma} + Z_1\right) - \left(\frac{p_2}{\gamma} + Z_2\right)$$

$$4.25 = \left(\frac{p_1}{\gamma} - \frac{p_2}{\gamma}\right) + \left(0 - 0.15\right)$$

$$\Rightarrow \qquad \frac{p_1}{\gamma} - \frac{p_2}{\gamma} = 4.4$$

 \Rightarrow

 \Rightarrow

$$p_1 - p_2 = 4.4 \times 0.8 \times 9.807 = 34.52$$
 kPa

(b) Again using equation (4.31)

 $h + 0.15 = \left(\frac{\gamma_m}{\gamma} - 1\right) h_m$ $4.4 = \left(\frac{13.596}{0.8} - 1\right) h_m$ $h_m = 0.275 \text{ m}$

 \Rightarrow

 \Rightarrow

EXAMPLE 11

The head loss from the entrance to the throat of a 254×127 venturimeter is $1/6^{\text{th}}$ times the throat velocity head. If mercury in the differential gauge attached to the meter deflects 101.6 mm, what is the flow of water through the venturimeter?

SOLUTION

Difference in piezometric head

$$h = \left(\frac{\gamma_m}{\gamma} - 1\right) h_m = \left(\frac{13.596}{1} - 1\right) 0.1016 = 1.2797 \text{ m}$$

Being horizontal, Bernoulli's equation applied at entrance and throat yields

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + h_f = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + \frac{1}{6}\frac{V_2^2}{2g}$$

 \Rightarrow

 \Rightarrow

$$\frac{p_1}{\gamma} - \frac{p_2}{\gamma} = 1.167 \frac{V_2^2}{2g} - \frac{V_1^2}{2g}$$

$$1.2797 = 1.167 \frac{Q^2}{2g} \left(\frac{1}{A_2^2} - \frac{1}{A_1^2} \right)$$
$$= \frac{Q^2}{2 \times 9.807} \frac{1}{(\pi/4)^2} \left(\frac{1.167}{0.127^4} - \frac{1}{0.254^4} \right) = 0.06 \text{ m}^3/\text{s}$$

[Here C_d is not introduced in the computation, because its effect has already been considered in terms of headloss]

EXAMPLE 12

The air supply to an oil-engine is measured by inducting air directly from the atmosphere into a large reservoir through a sharp edged orifice of 50 mm diameter. Pressure difference across the orifice is measured
by an alcohol manometer set at a slope of $\theta = \sin^{-1}(0.1)$ to the horizontal. Calculate the volume flow rate of air if manometer reading is 271 mm. Consider sp.gr. of alcohol is 0.8, $C_d = 0.62$, $C_c = 0.6$, atm.pr. = 775 mm of Hg, atm. temp. = 15.8 °C.

SOLUTION

Vertical height of manometric fluid

 $h_{al} = h \cdot \sin \theta = 0.271 \times 0.1 = 0.0271 \,\mathrm{m}$

Equivalent height of air column

$$h_{\text{air}} = \left(\frac{\gamma_m}{\gamma} - 1\right) h_{al} = \left(\frac{0.8 \times 10^3}{1.217} - 1\right) 0.0271 = 17.787 \text{ m}$$

Here, $V_1 = 0$, as section 1 is open to atmosphere, and

$$V_{2} = C_{v} \sqrt{2gh_{air} + V_{1}^{2}} = \frac{C_{d}}{C_{c}} \sqrt{2gh_{air}}$$
$$= \frac{0.62}{0.6} \sqrt{2 \times 9.807 \times 17.787}$$
$$= 19.3 \text{ m/s}$$

Applying equation of continuity, flow rate

$$Q = A_2 V_2 = C_c A_0 V_2 = 0.6 \times \frac{\pi}{4} \times 0.05^2 \times 19.3 = 0.02273 \text{ m}^3/\text{s}$$

EXAMPLE 13

What is the size of an orifice required to discharge 0.016 m³/s of water under a head of 8.69 m, while

 $C_d = 1.0.$ Solution

Velocity of flow

$$V = \sqrt{2gh} = \sqrt{2 \times 9.807 \times 8.69} = 13.055 \text{ m/s}$$

Applying equation of continuity, c/s area

$$A = \frac{Q}{V} = \frac{0.016}{13.055} \times 10^6 = 1225.58 \text{ mm}^2$$

EXAMPLE 14

Flow of air at 49°C is measured by a Pitot-static tube. If the velocity of air is 18.29 m/s and the coefficient of tube is 0.95, what differential reading will be shown in a water manometer? Assume density of air to be constant at 1.2 kg/m^3 .

SOLUTION

Applying equation (4.35)

$$V = C_p \sqrt{2gh_{air}}$$

$$18.29 = 0.95 \sqrt{2 \times 9.807 \times h_{air}}$$

$$h_{air} = 18.89792 \text{ m}$$

 \Rightarrow \Rightarrow Again from equation of differential manometer

$$h_{\text{air}} = \left(\frac{\gamma_w}{\gamma_{\text{air}}} - 1\right) h_w$$
$$18.89792 = \left(\frac{10^3}{1.2} - 1\right) h_w$$

 \Rightarrow

 \Rightarrow

$$89792 = \left(\frac{10^3}{1.2} - 1\right)h_w$$
$$h_w = 0.022704 \text{ m} = 22.704 \text{ mm}$$

EXAMPLE 15

A pump is installed in a pipeline, 5 cm dia., carrying oil of *sp.gr*: 0.83. It returns the oil to a 5 cm dia. pipe at the same elevation with a pressure increase of 1.4 kg/cm^2 . Quantity of oil flowing in the pipeline is 10 lit/s. The motor driving the pump delivers 3.8 metric H.P. to the pump shaft. Calculate loss of energy in the pump in kg-m/kg and kg-m/s.

SOLUTION

Applying Bernoulli's equation at sections 1 and 2,

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 + \text{ energy from pump}$$

$$= \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2 + \text{ losses in put}$$

$$Z_1 = Z_2, V_1 = V_2 \text{ and } p_2 > p_1$$

1002

$$Pump$$

Here,

$$\frac{p_1}{\gamma}$$
 + energy from pump = $\frac{p_2}{\gamma}$ + losses in pump

$$\Rightarrow \qquad \text{losses in pump} = \left(\frac{p_1}{\gamma} - \frac{p_2}{\gamma}\right) + \text{ energy from pump}$$

=

So,

$$= -\frac{1.4 \times 100^2}{10^3 \times 0.83} + \text{ energy from pump}$$

$$= -16.87 \text{ m} + \text{energy from pump}$$

Now, energy from pump = $3.8 \times 75 = 285$ kg × m/s Energy from pump per kg of oil

$$=\frac{285}{10\times10^3\times10^{-6}\times0.83\times10^3}=34.3 \text{ kgm/kg}$$

So, losses in pump

=
$$(-16.87 + 34.3) = 17.43$$
 kgm/kg of oil
or, $17.43 \times (10 \times 10^3 \times 10^{-6} \times 0.83 \times 10^3) = 144.7$ kgm/s

EXAMPLE 16

Water moves steadily through turbine @ 0.23 m³/s. Pressure at (1) and (2) are 186.4 kN/m^2 and 19.6 kN/m² respectively. Neglecting transfer of heat, determine the H.P. delivered to the turbine from water.

SOLUTION

Here,

$$A_1 = \frac{\pi}{4} 0.2^2 = 0.03141 \text{ m}^2, A_2 = \frac{\pi}{4} 0.4^2 = 0.12566 \text{ m}^2$$

So,

$$V_1 = \frac{0.23}{0.03141} = 7.322 \text{ m/s}, V_2 = \frac{0.23}{0.12566} = 1.83 \text{ m/s}$$

Applying Bernoulli's equation at sections 1 and 2

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2 + H$$

Head delivered to the turbine

$$H = \frac{p_1 - p_2}{\gamma} + (Z_1 - Z_2) + \frac{V_1^2 - V_2^2}{2g}$$
$$= \frac{(186.4 + 19.6)10^3}{9.807 \times 10^3} + (1.3 - 0) + \frac{7.322^2 - 1.83^2}{2 \times 9.807} = 24.8675 \text{ m}$$

H.P. delivered to turbine

$$=\frac{\gamma QH}{75}=\frac{10^3 \times 0.23 \times 24.8675}{75}=76.26 \text{ metric H.P.}$$

EXAMPLE 17

A venturimeter with 30 cm diameter inlet and 15 cm throat is used for measuring flow of oil of sp.gr. 0.9. The oil-mercury differential gauge shows a reading of 33 cm. Assuming the coefficient of meter 0.98, calculate the discharge.

SOLUTION

He

Here,
$$A_1 = \frac{\pi}{4} 0.3^2 = 0.07068 \text{ m}^2, A_2 = \frac{\pi}{4} 0.15^2 = 0.01767 \text{ m}^2$$

Again,
$$h_{\text{oil}} = \left(\frac{\gamma_m}{\gamma_{\text{oil}}} - 1\right) h_m = \left(\frac{13.596}{0.9} - 1\right) 0.33 = 4.655 \text{ m}$$

To obtain the discharge through venturimeter, substitute values in $Q = C_d A_1 A_2 \sqrt{\frac{2gh}{A_1^2 - A_2^2}}$ which is 0.17089 m³/s.



EXAMPLE 18

A 20 cm waterpipe has in it a venturimeter of throat diameter 12.5 cm as shown in the figure, which is connected to a mercury manometer showing a difference of 86.5 cm. Find the velocity in the throat and the discharge.

SOLUTION

Applying Bernoulli's equation in sections 1 and 2,

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + y_1$$

From the principle of manometer

$$\frac{p_1}{\gamma} + y_2 + 0.865 - (0.865 \times 13.596) - y_2 - y_1 = \frac{p_2}{\gamma}$$

$$\frac{p_1}{\gamma} = 10.8955 + y_1 + \frac{p_2}{\gamma}$$

 \Rightarrow

$$10.8955 + y_{1} + \frac{p_{2}}{\gamma} + \frac{V_{1}^{2}}{2g} = \frac{p_{2}}{\gamma} + \frac{V_{2}^{2}}{2g} + y_{1}$$

$$\Rightarrow \qquad \frac{V_{2}^{2}}{2g} = 10.8955 + \frac{V_{1}^{2}}{2g} = 10.8955 + \frac{V_{2}^{2}}{2g} \left(\frac{A_{2}}{A_{1}}\right)^{2}$$

$$= 10.8955 + \frac{V_{2}^{2}}{2g} \left(\frac{0.125}{0.2}\right)^{4}$$

$$\Rightarrow \qquad 0.8474 \frac{V_{2}^{2}}{2g} = 10.8955$$

$$\Rightarrow \qquad V_{2} = 15.88 \text{ m/s}$$

Discharge

$$Q = A_2 V_2 = \frac{\pi}{4} \times 0.125^2 \times 15.88 = 0.1948 \text{ m}^3\text{/s}$$

EXAMPLE 19

A venturimeter is installed in a pipeline of 30 cm diameter. The throat-pipe diameter ratio is 1:3. The pressure in the pipeline is 137.7 kN/m^2 and the vacuum in the throat 37.5 cm of mercury. If 4% of the differential head is lost between gauges, find the flow in the pipeline.

SOLUTION

Let us consider pipeline as section 1 and throat section 2.

So,
$$\frac{A_2}{A_1} = \left(\frac{1}{3}\right)^2 = \frac{1}{9} = \frac{V_1}{V_2}$$

$$\Rightarrow \qquad A_1 = 9A_2 \text{ and } V_2 = 9V_1$$



Again,

 $p_2 = -0.375 \times 13.596 \times 9.807 = -50 \text{ kN/m}^2$ $\frac{p_1}{\gamma} + \frac{V_1^2}{2g} = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + h_f$

Now,

$$\Rightarrow \qquad 0.96 \left(\frac{p_1}{\gamma} - \frac{p_2}{\gamma}\right) = \frac{V_2^2}{2g} - \frac{V_1^2}{2g}$$

$$\Rightarrow \qquad 0.96 \left(\frac{137.7 + 50}{9.807}\right) = \frac{(9V_1)^2}{2 \times 9.807} - \frac{V_1^2}{2 \times 9.807}$$
$$\Rightarrow \qquad V_1 = 2.122 \text{ m/s}$$

So, discharge

$$Q = A_1 V_1 = \frac{\pi}{4} \times 0.3^2 \times 2.122 = 0.1499 \text{ m}^3/\text{s}$$

EXAMPLE 20

At a point in the pipeline where diameter is 20 cm, velocity of water is 4 m/s and pressure 3.5 kg/cm^2 . At a point 15 m downstream, diameter reduces to 10 cm. Calculate pressure at this point if pipe is vertical with (a) downward flow and (b) upward flow.

SOLUTION

Considering no head loss between 1 and 2, velocity at 2 is

$$V_2 = \frac{A_1}{A_2} V_1 = \left(\frac{0.2}{0.1}\right)^2 \times 4 = 16 \text{ m/s}$$

(a) Here,

$$\Rightarrow \qquad \frac{3.5 \times 10^4}{1000} + \frac{4^2}{2 \times 9.807} + 15 = \frac{p_2}{\gamma} + \frac{16^2}{2 \times 9.807} + 0$$

 $\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2$

$$\Rightarrow \qquad \frac{p_2}{\gamma} = 37.7638 \text{ m of water} = 3.77638 \text{ kg/cm}^2$$

(b) Here,
$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2$$

$$\Rightarrow \qquad \frac{3.5 \times 10^4}{1000} + \frac{4^2}{2 \times 9.807} + 0 = \frac{p_2}{\gamma} + \frac{16^2}{2 \times 9.807} + 15$$

$$\Rightarrow \qquad \frac{p_2}{\gamma} = 7.7638 \text{ m of water} = 0.77638 \text{ kg/cm}^2$$



EXAMPLE 21

The velocity of water in pipeline was measured with Pitot-static tube. The difference of head at the centre of the pipe was 89 mm of water. If mean velocity of water is *three-fourth* the velocity at the centre, determine rate of flow. Diameter of the pipe is 280 mm. Take $C_p = 0.99$.

 $V = C_p \sqrt{2gh} = 0.99 \sqrt{2 \times 9.807 \times 0.089} = 1.308 \text{ m/s}$

SOLUTION

Velocity at the centre of the pipe

Mean velocity

$$=\frac{3}{4}$$
 × 1.308 = 0.981 m/s

Rate of flow

$$Q = AV = \frac{\pi}{4} 0.28^2 \times 0.981 = 0.0604 \text{ m}^3/\text{s}$$

EXAMPLE 22

Adjacent figure shows a siphon discharging oil (sp.gr. = 0.84) from a reservoir into open air. If the velocity of flow in the pipe is V, the head loss from point 1 to 2 is $2V^2/2g$, and that from point 2 to 3 is $3V^2/2g$, determine the volume flow rate in the siphon pipe and absolute pressure at point 2. Assume an atmosphere pressure 101.358 kPa.

SOLUTION

Different head losses between sections are as following:

$$h_{1-2} = \frac{2V^2}{2g}$$
, and $h_{2-3} = \frac{3V^2}{2g}$
 $h_{1-3} = h_{1-2} + h_{2-3} = \frac{5V^2}{2g}$

Applying Bernoulli's equation between sections 1 and 3

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_3}{\gamma} + \frac{V_3^2}{2g} + Z_3 + h_{1-3}$$

 $V_1 = 0, V_2 = V_3 = V$ and being open to atmosphere, $p_1 = p_3 = 0$

Here,

:.

$$(4.6 - 1.5) = \frac{V_2}{2g} + 0 + \frac{5V^2}{2g}$$
$$V = 3.183 \text{ m/s} = V_2 = V_3$$

$$\Rightarrow$$

So,

Volume flow rate $Q = A_3 V_3 = \frac{\pi}{4} \times 0.075^2 \times 3.183 = 0.014 \text{ m}^3/\text{s}$



Applying Bernoulli's equation between sections 1 and 2

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2 + h_{1-2}$$

$$\Rightarrow \qquad (4.6 - 1.5) = \frac{p_2}{\gamma} + \frac{3.183^2}{2 \times 9.807} + 4.6 + 2\frac{3.183^2}{2 \times 9.807}$$

$$\Rightarrow \qquad \frac{p_2}{\gamma} = -3.049 \text{ m of oil}$$

$$\Rightarrow \qquad p_2|_{\text{gauge}} = (-3.409 \times 0.84 \times 9.807) = -25.117 \text{ kPa}$$

$$\Rightarrow \qquad p_2|_{\text{absolute}} = (101.358 - 25.117) = 76.241 \text{ kPa}$$

EXAMPLE 23

Compute the ideal flow rate through the pipe system shown in the figure.

SOLUTION

Applying Bernoulli's equation between 1 and 2 and assuming no head loss,

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2$$

$$\Rightarrow \qquad \frac{p_1}{\gamma} + \frac{V_1^2}{2g} + 0.6 \sin 30^\circ = \frac{p_2}{\gamma} + 0 + 0$$

$$\Rightarrow \qquad \frac{V_1^2}{2g} = \left(\frac{p_2}{\gamma} - \frac{p_1}{\gamma}\right) - 0.3$$

Again from the equation of manometer,

$$\frac{p_1}{\gamma} - 1.2\sin 60^\circ = \frac{p_2}{\gamma}$$

$$\Rightarrow \qquad \left(\frac{p_1}{\gamma} - \frac{p_2}{\gamma}\right) = 1.0392$$

Combining these two results, we obtain

$$\frac{V_1^2}{2g} = 1.0392 - 0.3 = 0.7392$$

 \Rightarrow

$$V_1 = 3.807 \text{ m/s}$$

Ideal flow rate $Q = A_1 V_1 = \frac{\pi}{4} 0.2^2 \times 3.807 = 0.1196 \text{ m}^3/\text{s}$



EXAMPLE 24

The flow from two reservoirs mixes together and flows through a common pipe. Both reservoir contain same liquid and open to atmosphere. Neglecting frictional loss, find flow rate through common pipe.

SOLUTION

Applying Bernoulli's equation between 1 and 2

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2$$

$$\Rightarrow \qquad \frac{p_{\text{atm}}}{\gamma} + 0 + h_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + 0$$

$$\Rightarrow \qquad \frac{p_{\text{atm}}}{\gamma} + h_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g}$$



Applying Bernoulli's equation between 3 and 4, in the similar way,

$$\Rightarrow \qquad \frac{p_{\rm atm}}{\gamma} + h_2 = \frac{p_4}{\gamma} + \frac{V_4^2}{2g}$$

Applying Bernoulli's equation between 5 and 6, in the similar way

$$\frac{p_5}{\gamma} + \frac{V_5^2}{2g} + Z_5 = \frac{p_6}{\gamma} + \frac{V_6^2}{2g} + Z_6$$
$$\frac{p_5}{\gamma} + \frac{V_5^2}{2g} + 0 = \frac{P_{\text{atm}}}{\gamma} + \frac{V_6^2}{2g} + (-h_3)$$
$$\frac{p_5}{\gamma} = \frac{P_{\text{atm}}}{\gamma} - h_2$$
(a)

(a)

 \Rightarrow

 \Rightarrow

$$\frac{p_5}{\gamma} = \frac{P_{\text{atm}}}{\gamma} - h_3 \tag{c}$$

As points 2,4 and 5 lie on the boundary of control volume, as shown in the figure,

$$p_2 = p_4 = p_5 \tag{d}$$

Substituting (c) and (d) into (a),

 $\frac{P_{\text{atm}}}{\gamma} + h_1 = \frac{P_{\text{atm}}}{\gamma} - h_3 + \frac{V_2^2}{2g}$ $V_2 = \sqrt{2g(h_1 + h_3)}$ (e)

Substituting (c) and (d) into (b),

$$\frac{P_{\rm atm}}{\gamma} + h_2 = \frac{P_{\rm atm}}{\gamma} - h_3 + \frac{V_4^2}{2g}$$
$$V_4 = \sqrt{2g(h_2 + h_3)}$$
(f)

 \Rightarrow

 \Rightarrow

(b)

(b)

So, flow through common pipe

$$Q = Q_1 + Q_2 = A_1 V_2 + A_2 V_4$$

= $\frac{\pi}{4} \Big[d_1^2 \sqrt{2g(h_1 + h_3)} + d_2^2 \sqrt{2g(h_2 + h_3)} \Big]$
= $\frac{\pi}{4} \sqrt{2g} \Big[d_1^2 \sqrt{h_1 + h_3} + d_2^2 \sqrt{h_2 + h_3} \Big]$

EXAMPLE 25



 \Rightarrow

Solving (a) and (b), $\frac{V_2^2}{2g} = 53.989$

 \Rightarrow

 $Q = A_2 V_2 = \frac{\pi}{4} 0.06^2 \times 32.541 = 0.092 \text{ m}^3/\text{s}$ Flow rate

 $V_2 = 32.541 \text{ m/s}$

EXAMPLE 26

Inlet and throat diameters of a horizontal venturimeter are 30 cm and 10 cm respectively. The liquid flowing through the venturimeter is water. Pressure intensity at inlet is 13.734 N/cm², while vacuum pressure at throat is 37 cm of mercury. Find flow rate, assuming 4% of the differential head is lost between inlet and throat. Find also value of C_{a^*}

SOLUTION

Now,

$$A_1 = \frac{\pi}{4} 0.3^2 = 0.07068 \text{ m}^2$$
, $A_2 = \frac{\pi}{4} 0.1^2 = 7.854 \times 10^{-3} \text{ m}^2$

$$p_1 = 13.734 \text{ N/cm}^2 = 137.34 \text{ kPa}$$

$$p_2 = -37 \text{ cm of Hg} = -(0.37 \times 13.596 \times 9.807) = -49.33 \text{ kPa}$$

Therefore, differential head

ŀ

$$h = \frac{p_1}{\gamma} - \frac{p_2}{\gamma} = \frac{(137.34 + 49.33) \times 10^3}{9.807 \times 10^3} = 19.034 \text{ m of water}$$

Head loss $h_f = 0.04h = 0.76137$ m of water

Coefficient of discharge $C_d = \sqrt{\frac{h - h_f}{h}} = 0.979$

Flow rate
$$Q = C_d A_1 A_2 \sqrt{\frac{2gh}{A_1^2 - A_2^2}}$$
 and substitute the values to find as 0.14949 m³/s.

EXAMPLE 27

In a vertical pipe conveying oil (sp.gr. = 0.8), two pressure gauges have been installed at *A* and *B*, where the diameters are 16 cm and 8 cm respectively. *A* is 2 m above *B*. The pressure gauge readings show that pressure at *B* is greater than pressure at *A* by 0.981 N/cm^2 . Neglecting all losses, calculate the flow rate. If gauges at *A* and *B* are replaced by tubes filled with same liquid and connected to a *U*-tube containing mercury, calculate difference in mercury level in two limbs of the *U*-tube.

SOLUTION

Here,

$$A_2 = \frac{\pi}{4} 0.08^2 = 5.0265 \times 10^{-3} \text{ m}^2$$

 $A_{\rm l} = \frac{\pi}{4} \, 0.16^2 = 0.0201 \, {\rm m}^2,$

Applying Bernoulli's theorem,

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + Z_1 = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + Z_2$$

$$\Rightarrow \qquad \frac{p_1 - p_2}{\gamma} + (Z_1 - Z_2) = \frac{V_2^2 - V_1^2}{2g}$$

$$\Rightarrow \qquad -\frac{0.981 \times 10^4}{9.807 \times 10^3 \times 0.8} + 2 = \frac{Q_2^2}{2g} \left[\frac{1}{A_2^2} - \frac{1}{A_1^2} \right]$$



$$\Rightarrow \qquad 0.7496 = \frac{Q_2^2}{2 \times 9.807} \left[\frac{1}{\left(5.0265 \times 10^3 \right)^2} - \frac{1}{\left(0.0201 \right)^2} \right]$$

 $\Rightarrow \qquad Q = 0.019906 \text{ m}^3/\text{s}$

Piezometric head difference

$$h_{\text{oil}} = \frac{p_1 - p_2}{\gamma} + (Z_1 - Z_2)$$

= 0.7496 m = $h_m \left(\frac{\gamma_m}{\gamma_{\text{oil}}} - 1\right) = h_m \left(\frac{13.596}{0.8} - 1\right)$
 $h_m = 0.04686$ m

 \Rightarrow

 \Rightarrow

EXAMPLE 28

A submarine moves horizontally in sea and has its axis 15 m below the surface of water. A Pitottube properly placed just in front of the submarine and along its axis is connected to the two limbs of *U*-tube containing mercury. Difference of mercury level is found to be 170 mm. Find speed of the submarine knowing that sp.gr. of mercury is 13.6 and that of sea water is 1.026 with respect to fresh water.

SOLUTION

Now,
$$h_{\text{sea}} = h_m \left(\frac{\gamma_m}{\gamma_{\text{sea}}} - 1\right) - 0.17 \left(\frac{13.6}{1.026} - 1\right) = 2.0834 \,\text{m}$$

Speed of the submarine

$$V = C_p \sqrt{2gh_{\text{sea}}} = 1.0\sqrt{2 \times 9.807 \times 2.0834} = 6.3924 \text{ m/s}$$

EXAMPLE 29

A Pitot tube is inserted in a pipe of 300 mm diameter. Static pressure in a pipe is 100 mm of mercury (vacuum). The stagnation pressure at the centre of the pipe, recorded by a Pitot-tube is 0.981 N/cm². Calculate rate of flow of water through pipe, if mean velocity of flow is 0.85 times the central velocity. Take $C_v = 0.98$. **SOLUTION**

Area of pipe	$A = \frac{\pi}{4} \ 0.3^2 = 0.07068 \ \mathrm{m}^2,$
Static pressure	$p_o = -100 \text{ mm or Hg} = -13.3335 \times 10^3 \text{ N/m}^2$
Stagnation pressure	$p_s = 0.981 \times 10^4 \text{ N/m}^2$
Rise in fluid	$h = \frac{p_s}{\gamma} - \frac{p_o}{\gamma} = \frac{9.81 + 13.3335}{9.807} = 2.359 \text{ m}$
Velocity at centre	$V = 0.98 \sqrt{2 \times 9.807 \times 2.359} = 6.666 \text{ m/s}$

Mean velocity $V_m = 0.85 \times 6.666 = 5.6661 \text{ m/s}$ Rate of flow $Q = V_m \times A = 5.6661 \times 0.07068 = 0.400479 \text{ m}^3/\text{s}$

EXAMPLE 30

Find the discharge of water flowing through a pipe 30 cm diameter placed in an inclined position where a venturimeter is inserted, having throat diameter of 15 cm. Level difference of manometric fluid (sp.gr. = 0.6) is 30 cm. Head loss between inlet and throat is 0.2 times the kinetic head of pipe. SOLUTION





EXERCISE

1. (a) State Bernoulli's theorem for steady flow of an incompressible fluid.

(b) Derive an expression for Bernoulli's equation from the first principle.

- 2. Write down Euler's equation of motion along a streamline and show how Bernoulli's equation can be derived from it for a compressible fluid.
- With the aid of a diagram, mention a device for measuring the velocity of flow at any point in a pipe. 3.
- 4. Define (i) total pressure (ii) static pressure (iii) stagnation pressure (iv) HGL (v) EGL.
- 5. What is venturimeter? Derive an expression for the discharge through a venturimeter.
- 6. What is the difference between Pitot tube and Pitot-static tube?
- 7. Compare relative merits and demerits of an orificemeter and a venturimeter.
- Define an orificemeter. Obtain an expression for flow rate through an orificemeter. 8.

9. Water is pumped at the rate of 300 lit/s through a 30 cm pipe up to a hill top. On the hill top, which has an elevation of 50 m, the diameter of pipeline reduces to 20 cm. If the pump maintains a pressure of 1 kg/cm² at the hill top, what is the pressure at the foot-hill having zero elevation?

Ans. [6.373 kg/cm², 258.6 H.P.]

10. Water flows @75 lit/s in a pipe whose diameter at sections 1-1 and 2-2 are 300 mm and 150 mm respectively, which are at heights of 5 m and 3 m above datum. If pressure at 1-1 be 450 kPa, neglecting loss of energy, calculate pressure at 2-2.

Ans. [461.177 kPa]

11. Oil of specific gravity 0.90 flows in a pipe of 30 cm diameter @120 lit/s and pressure at point A is 0.25 kg/cm^2 gauge. If the point A is 5.2 m above datum, compute total energy at A in terms of metre of oil.

Ans. [8.152 m]

12. A vertical venturimeter of d/D = 0.6 is filled in a 10 cm dia. pipe. Throat is 20 cm above inlet and $C_d = 0.92$. Determine (i) pressure difference at inlet and throat (ii) difference on a vertical differential manometer, when liquid of sp.gr. = 0.8 flows through the venturimeter @50 lit/s.

Ans. [1.33 kg/cm², 1.024 m]

13. A venturimeter is used for measuring petrol (sp.gr. = 0.81) in a pipe line inclined 35° to horizontal. If throat area ratio is 4 and difference in mercury level in the gauge is 5 cm, calculate flow rate in lit/hr. if pipe diameter is 30 cm.

Ans. $[2.52 \times 10^5 \text{ lit/hr.}]$

14. A pipe carrying oil (sp.gr. = 0.877) changes in size from 15 cm at section A to 45 cm at section B. Section A is 3.6 m lower than B and the pressures are 0.92 kg/cm² and 0.61 kg/cm² respectively. If discharge is 0.145 m³/s, determine head loss and direction of flow.

Ans. [3.32 m of oil, A to B]

15. Find the velocity of flow of CCl_4 (sp.gr. = 1.59) through a pipe, when a differential gauge attached to Pitot tube shows a reading of 10 cm. Take the coefficient of Pitot tube as 0.98.

Ans. [3.77 m/s]

16. A Pitot-tube was used to measure the quantity of water flowing in a pipe of 30 cm diameter. The water was raised to a height of 25 cm above centerline of the pipe in a vertical limb of the tube. If mean velocity is 0.78 times the velocity at center and $C_p = 0.98$, find quantity of water in lit/s. Static pressure head at the center of the pipe is 20 cm.

Ans. [53.52 lit/s]

17. Water flows @ 0.147 m^3 /s through a 15 cm diameter orifice inserted in a 30 cm diameter pipe. If pressure gauges fitted *u/s* and *d/s* of the orifice plate have shown readings of 1.8 kg/cm² and 0.9 kg/cm² respectively, find the coefficient of discharge of the orificemeter.

Ans. [0.61]

18. An orificemeter with orifice diameter 15 cm is inserted in a pipe of 30 cm diameter. The pressure gauges fitted u/s and d/s of the orificemeter gives readings of 4.715 N/cm² and 9.81 N/cm² respectively. Find the flow rate while $C_d = 0.6$. If instead of water, oil of sp.gr. = 0.8 is flowing through the orificemeter in which pressure difference is measured by a mercury-oil differential manometer on two sides of the orificemeter, find rate of flow of oil when reading of manometer is 40 cm.

Ans. [108.434 lit/s, 122.68 lit/s]

BIBLIOGRAPHY

FLUID MECHANICS

- Daugherty, R.L., Franzini, J.B., Finnemore, E.J., Fluid Mechanics with Engineering Applications (SI Metric Ed.) McGraw-Hill Book Co., Singapore, 1989.
- 2. White, F.M., Fluid Mechanics (4th. Ed.), McGraw-Hill Companies Inc., Singapore, 1999.
- 3. Evett, J.B. and Liu, C., Schaum's 2500 Solved Problems in Fluid Mechanics and Hydraulics, McGraw-Hill Book Co., Singapore, 1989.
- 4. Ray, D.N., Applied Fluid Mechanics, Affiliated East West Press Pvt. Ltd., New Delhi, 2002.
- 5. Modi, P.N. and Seth, S.M., Hydraulics and Fluid Mechanics, Standard Book House, Delhi, 1987.
- 6. Som, S.K. and Biswas, G., Introcution to Fluid Mechanics and Fluid Machines, Tata McGraw-Hill Publishing Company Limited, New Delhi, 2004.
- 7. Bansal, R.K., A Text Book of Fluid Mechanics and Hydraulic Mechanics, Laxmi Publications Pvt. Ltd., New Delhi.
- 8. Jain, A.K., Fluid Mechanics.

This page intentionally left blank

ENGINEERING & TECHNOLOGY EXAMINATIONS, JUNE-2005 MECHANICAL SCIENCE SEMESTER-2

Time: 3 Hours

Full Marks: 70

Note: Answer Question No. 1 and any *three* from Group-A and any *two* from Group-B.

1.	Choo	se the	correct answer for any ten of the fol	lowing	g: $10 \times 1 = 10$	
	(a)	Worl	k done in a free expansion process is	5		
		(i)	Positive	(ii)	Negative	
		(iii)	Zero	(iv)	Maximum	
	(b)	The	internal energy of system is a function	on of o	nly	
		(i)	Pressure	(ii)	Absolute temperature	
		(iii)	Volume	(iv)	Pressure and temperature	
	(c)	Entr	opy change depends on			
		(i)	Heat transfer	(ii)	Mass transfer	
		(iii)	Change of temperature	(iv)	Thermodynamic state	
	(d)	The	cycle with constant volume heat add	ition a	nd heat rejection is called	
		(i)	Carnot cycle	(ii)	Joule cycle	
		(iii)	Rankine cycle	(iv)	Otto cycle	
	(e)	'Tripple point' of a pure substance is a point at which				
		(i)	Liquid and vapour exist together			
		(ii)	Solid and vapour exist together			
		(iii)	Solid and liquid exist together			
		(iv)	Solid, liquid and vapour phases exi	st toget	ther	
	(f)	The more effective way of increasing the efficiency of a Carnot engine is to				
		(i)	Increase the upper temperature	(ii)	Decrease the upper temperature	
		(iii)	Increase the lower temperature	(iv)	Decrease the lower temperature	
	(g)	An ideal fluid is the one which is				
		(i)	Non-viscous and incompressible	(ii)	Compressible and has low density	
		(iii)	Elastic and viscous	(iv)	Steady and incompressible	
	(h)	Mere	cury is used inbarometers on account	t of its		
		(i)	Negligle capillary effect	(ii)	High density	
		(iii)	Very low vapour pressure	(iv)	Low compressibility	

Velocity of flow reduces to zero

Total energy is maximum

		(i)	Pseudo plastic	(11)	Dilatant
		(iii)	Inviscid	(iv)	Newtonian
				Group-A	
			(Ans	wer any three q	uestions)
2.	(a)	Wha	t do you mean by conti	ol volume?	2
	(b)	Wha	t is Quasi-static proces	ss?	4
	(c)	A ce cons	rtain gas has $C_p = 1.968$ tant.	$B \text{ and } C_v = 1.507$	kJ/kg. Find its molecular weight and the gas 6
		A co trans heat	nstant volume chambe sferred to the gas until transferred, and the ch	r of 0.3 m ³ capa the temperature anges in interna	city contains 2 kg of this gas at 5°C. Heat is e becomes 100°C. Find the work done, the l energy, enthalpy and entropy.
3.	(a)	Wha volu	t is a steady flow pro me and explain the var	cess? Write the ious terms in it.	steady flow energy equation for a control 3
	(b)	At the velocizionta nozz the r exit	ne inlet to a certain noz city is 60 m/s. At the di al and there is negligil ile. (b) If the inlet area i mass flow rate. (c) If th area of the nozzle.	The enthalp excless the enthalp scharge end; the ble heat loss from s 0.1 m^2 and the e specific volum	y of the fluid passing is 3000 kJ/kg and the e enthalpy is 2762 kJ/kg. The nozzle is hori- om it. (a) Find the velocity at exit from the specific volume at inlet is $0.187 \text{ m}^3/\text{kg}$, find he at the nozzle exit is $0.498 \text{ m}^3/\text{kg}$, find the 9
4.	(a)	State	e Clausius inequality an	d explain its sig	nificance. 4
	(b)	Defi	ne COP of a refrigerate	or and a heat pu	mp. Establish the relation
		COP	$P_{\rm hp} = \rm COP_{ref} + 1.$		8
5.	(a)	Wha	t is a pure substance?		2
	(b)	Wha of w	t is the 'critical point'?	State the values	s of critical pressure and critical temperature 3
	(c)	Wha	t arc the four basic com	ponents of a ste	am power plant? Show by a block diagram.3
	(d)	Drav turbi	w the nature of <i>p</i> - <i>v</i> and ne inlet).	d T-S plots of a	Rankine cycle (with superheated steam at 4
6.	(a)	Drav corre	v a block diagram of a esponding <i>p-v</i> and <i>T-S</i>	vapour compre plots.	ession refrigeration cycle and also show the 5

(ii)

(iv)

A stagnation point is a point is fluid flow where

Which fluid does not experience shear stress during flow?

Pressure is zero

Total energy is zero

(i)

(j)

(i)

(iii)

Question Paper

(b) An engine working on the Otto cycle is supplied with air at 0.1 MPa, 35°C. The compression ratio is 8. Heat supplied is 2100 kJ/kg. Calculate the maximum pressure and temperature of the cycle, the cycle efficiency, and the mean effective pressure (For air, $C_p = 1.005$, $C_v = 0.718$ and R = 0.287 kJ/kg K). 7

Group-B

(Answer any *two* questions)

7.	(a)	State and explain Newton's law of viscosity. 3
	(b)	Define non-Newtonian fluid with two examples. 3
	(c)	The space between two large flat and parallel walls 25 mm apart is filled with liquid of viscosity 0.7 N.S/m^2 . Within this space a thin flat plate, 250 mm × 250 mm is towed at a velocity of 150 mm/s at a distance of 19 mm from one wall, the plate and its movement being parallel to the walls. Assuming linear variation in velocity between the plates and the walls. Determine the force exerted by the liquid on the plate. 6
8.	(a)	Define the term "bulk" modulus of elasticity" of a compressible fluid.
	(b)	State Pascals law of hydrostatics. 2
	(c)	Calculate the atmospheric pressure at the end of troposphere which extends up to a height of 11 km from sea level. Consider a temperature variation in the troposphere as $T = 288 - 0.00649 z$ where z is the elevation in metres and T is temperature in kelvin. The atmospheric pressure at sea level is 101 kN/m, for air $R = 287$ J/kg K. Derive the equation you see for your calculation. 8
9.	(a)	Write the difference between steady flow and uniform flow. 2
	(b)	Define (i) Streamline and (ii) Streakline 2
	(c)	Obtain an expression for continuity equation for a two dimensional flow. 5
	(d)	A two dimensional flow is described in the Lagrangian as
		$X = X_0 e^{kt} + Y_0 (1 - e^{-2kt})$
		$Y = Y_0 e^{kt}$
	Find	the equation of path line of the particle and the velocity components in Eulerian system. 3
10.	(a)	Write Bernoulli's equation and describe the various terms in it.2
	(b)	What are the assumptions involved in the derivation of Bernoulli's equation? 3
	(c)	Water flows through a 300 mm \times 150 mm venturimeter at the rate of 0.037 m ³ /3 and differential gauge is deflected 1 m as shown in the figure. Specific gravity of the gauge

liquid is 1.25. Determine the coefficient of discharge of the meter.

ENGINEERING & TECHNOLOGY EXAMINATIONS, JUNE - 2006 MECHANICAL SCIENCE SEMESTER- 2

Time Allotted: 3 Hours

The figures in the margin indicate full marks.

Candidates are required to give their answers in their own words as far as practicable.

Any missing data may be assumed suitably with Justification.

Note: Answer Question No. 1 and any three from Group-A and any two from Group-B.

1. Choose the correct answers for any *ten* of the following:

- (i) An open system is one in which
 - (a) heat and work cross the boundary of the system, but the mass of the working substance does not
 - (b) mass of working substance crosses the boundary of the system but the heat and work do not
 - (c) both the heat and work as well as mass of the working substances cross the boundary of the system
 - (d) neither the heat and work nor the mass of the working substances crosses the boundary of the system.

(d)

(b)

(d)

Energy.

- (ii) Which of the following is an intensive property of a thermodynamic system ?
 - (a) Volume (b) Temperature
 - (c) Mass
- (iii) During throttling process
 - (a) internal energy does not change
 - (c) entropy does not change
 - (e) volume change is negligible.
- (iv) Kelvin-Planck's law deals with
 - (a) conservation of energy
 - (c) conservation of mass
 - (e) conversion of work into heat.
- (v) If a heat engine attains 100% thermal efficiency, it violates
 - (a) zeroth law of thermodynamics
 - (c) second law of thermodynamics

- (b) conservation of heat
- (d) conversion of heat into work

pressure does not change

enthalpy does not change

- (b) first law of thermodynamics
- (d) law of conservation of energy.

 $10 \times 1 = 10$

Full Marks: 70

(a)

(vi) For an irreversible process entropy change is

greater than $\delta Q / T$

	(c)	less than δQ / T	(d)	equal to zero.
(vii)	The c	continuity equation (at two sections 1 and 2	2) for	an incompressible fluid is given as
	(a)	$\rho_1 A_1 V_1^2 = \rho_2 A_2 V_2^2$	(b)	$\rho_1 A_1 V_1 = \rho_2 A_2 V_2$
	(c)	$\mathbf{A}_1 \mathbf{V}_1 = \mathbf{A}_2 \mathbf{V}_2$	(d)	$\rho_1^2 A_1 V_1 = \rho_2^2 A_2 V_2.$
(viii)	Euler	's equation is written as		
	(a)	$\frac{\mathrm{d}p}{\rho} + \upsilon^2 \cdot d\upsilon + \mathbf{g} \cdot \mathbf{dz} = 0$	(b)	$\frac{\mathrm{d}p}{\rho} + \upsilon \cdot d\upsilon + \mathrm{g} \cdot \mathrm{dz} = 0$
	(c)	$\frac{\mathrm{d}p}{\rho} + \upsilon^2 \cdot d\upsilon + \mathrm{g}^2 \cdot \mathrm{d}z = 0$	(d)	$\frac{\mathrm{d}p}{\rho^2} + v^2 \cdot dv + \mathbf{g} \cdot \mathbf{dz} = 0.$
(ix)	Kine	matics viscosity is defined as equal to		
	(a)	dynamic viscosity \times density	(b)	dynamic viscosity / density
	(c)	dynamic viscosity / pressure	(d)	pressure \times density.
(x)	For p	ipes, turbulent flow occurs when Reynolds	numt	per is
	(a)	less than 2000	(b)	between 2000 and 4000
	(c)	more than 4000	(d)	less than 4000.
(xi)	The c	change of entropy when heat is absorbed b	y the	gas is
	(a)	positive	(b)	negative
	(c)	positive or negative	(d)	none of these.
(xii)	Sphe	rical shape of droplets of mercury is due to)	

(b)

equal to $\delta Q / T$

- (a) high density (b) high surface tension
- (c) high adhesion (d) low vapour pressure.

Group-A

(Answer any *three*)

2. (a) Explain intensive property.

(b) Prove that for polytropic process
$$W_{1-2} = \frac{P_2 V_2 - P_1 V_1}{1-n}$$
 4

(c) A gas is enclosed in a piston-cylinder assembly as a system. The gas is initially at a pressure of 500 kPa and occupies a volume of 0.2 m³. The gas is taken to the final state where $P_2 = 100$ kPa by following two different processes. Calculate the work done by the gas in each case.

6

4

3

- The volume of the gas is inversely proportional to the pressure. (i)
- The process follows the path. $PV^{\gamma} = \text{constant}$ where $\gamma = 1.4$. (ii)
- Explain first law of thermodynamics for a closed system undergoing a change of state. 3 3. (a)
 - A turbine operates under steady flow conditions, receiving steam at the following state: (b) Pressure 1.2 MPa, Temperature 188°C, Enthalpy 2785 kJ/kg, Velocity 33.3 m/sec and elevation 3 m. The steam leaves the turbine at the following state: Pressure 20 kPa, Enthalpy 2512 kj/kg, Velocity 100 m/sec and elevation 0 m. Heat is lost

to the surroundings at the rate of 0.29 kj/kg. If the rate of steam flow through the turbine is 0.42 kg/sec. what is the power output of the turbine in kW? 9

- (a) What is cyclic heat engine? 4.
 - (b) A Carnot cycle has an efficiency of 32%. Assuming that the lower temperature is kept constant, determine the' percentage increase of the upper temperature of the cycle if the cycle efficiency is raised to 48%. 8
- Write the principle of entropy. 5. (a)
 - Calculate the entropy change of 10 gm of water at 20°C, when it is converted to ice at (b) -10°C. Assume the specific heat of water to remain constant at 4.2 J/gmK and that of ice to be half of the value and taking the latent heat of fusion of ice at 0°C to be 335 J/gm. 9
- Prove that entropy change for an ideal gas 6. (a)

$$s_2 - s_1 = c_p \ln \frac{T_2}{T_1} - R \ln \frac{p_2}{p_1}$$
 5

A vessel of volume 0.04 m³ contains a mixture of saturated water and saturated steam at (b) a temperature of 250°C. The mass of the liquid present is 9 kg. Find the pressure, the mass, the specific volume, the enthalpy, the entropy and the internal energy. 7

Group-B

(Answer any *two* questions)

- 7. (a) State Newton's law of viscosity.
 - What is Bulk Modulus of Elasticity? (b)
 - The space between two large flat and parallel walls 25 mm apart is filled with a liquid of (c) absolute viscosity 0.7 Ns/m^2 . Within this space a thin flat plate, 250 mm × 250 mm is towed at a velocity of 150 mm/s at a distance of 6 mm from one wall, the plate and its movement being parallel to the walls. Assuming linear variations of velocity between the plate and the walls, determine the force exerted by the liquid on the plate. 8
- The velocity distlibution for a two-dimensional incompressible flow is given by 8. (a)

$$u = -\frac{x}{x^2 + y^2}, u = -\frac{y}{x^2 + y^2}$$

Show that it satisfies continuity.

4

(b) Water is flowing through two different pipes A and B to which an inverted differential manometer having an oil of sp.gr. 0.9 is connected. The pressure in the pipe A is 2.5 m of water. Find the pressure in the pipe B for the manometer readings as shown in **Fig.** 1. 8





9. (a) Derive Euler's equation of motion along a streamline.

(b) A vertical venturimeter shown in Fig. 2 has an area ratio of 5. It has a throat diameter of 1 cm. When oil of sp. gr. 0.8 flows through it the mercury in the differential gauge indicates a difference in height of 12 cm. Find the discharge through the venturimeter. 7



Fig. 2

ENGINEERING & MANAGEMENT EXAMINATIONS, JUNE - 2007 MECHANICAL SCIENCE SEMESTER- 2

Time: 3 Hours

Full Marks: 70

Group–A (Multiple Choice Type Questions)

1.	Cho	ose th	e correct alternatives for the following	ng:		$10 \times 1 = 10$	
	(i)	Which of the following quantities is not a property of a system?					
		(a)	Pressure	(b)	Temperature		
		(c)	Heat	(d)	Specific volume.		
	(ii)	Duri	ng throttling process				
		(a)	internal energy does not change				
		(b)	pressure does not change				
		(c)	entropy does not change				
		(d)	enthalpy does not change				
		(e)	volume does not change.				
	(iii)	Ther	mal power plant works on				
		(a)	Carnot cycle	(b)	Joule cycle		
		(c)	Rankine cycle	(d)	Otto cycle.		
	(iv)	Work done in a free expansion process is					
		(a)	positive	(b)	negative		
		(c)	zero	(d)	maximum.		
	(v)	The more effective way of increasing efficiency of Carnot engine is to					
		(a)	increase higher temperature				
		(b)	decrease higher temperature				
		(c)	increase lower temperature				
		(d)	decrease lower temperature.				
	(vi)	A refrigerator and a heat pump operate between the same temperature line of the refrigerator is 4, the C.O.P. of the pump would be					
		(a)	3	(b)	4		
		(c)	5	(d)	cannot be predicted		
		(e)	None of these.				

- (vii) A pitot tube is used for measuring(a) state of flow(b) density of fluid
 - (c) velocity of fluid (d) none of these.
- (viii) Stream line, path line and streak line are identical when
 - (a) the flow is uniform
 - (b) the flow is steady
 - (c) the flow velocities do not change steadily with time
 - (d) the flow is neither steady nor uniform.
- (ix) Bernoulli's equation deals with the law of conservation of
 - (a) mass (b) momentum
 - (c) energy (d) work.
- (x) Which fluid does not experience shear stress during flow?
 - (a) Pseudoplastic (b) Dilataut
 - (c) Inviscid (d) Newtonian.

force necessary to maintain this velocity and also the power required.

Group-B

(Short Answer Type Questions)

		Answer any three questions	$3 \times 5 = 15$
2.	(a)	What is the basic difference between a process and a cycle?	2
	(b)	Show that work done in isothermal process from state 1 to state 2 is given	ı by:
		$W_{1-2} = p_1 V_1 (lnp_1 - lnp_2).$	3
3.	Wha and a	t is steady flow process? Write the steady flow energy equation for a single strain single stream leaving a control volume and explain the various terms in it.	ream entering $2+3$
4.	(a)	State second law of thermodynamics.	3
	(b)	What is a perpetual motion machine of 2nd kind?	2
5.	(a)	Explain the no-slip condition of viscous fluids.	2
	(b)	State and prove Pascal's law of hydrostatics.	3
6.	A ho dista	t plate of area 0.125 m^2 is pulled at 0.25 m/s with respect to another parallel nt from it the space between the plates containing water of viscosity 0.001 N -	el plate 1 mm s/m^2 . Find the

Group–C (Long Answer Type Questions) Answer any *three* questions.

 $3 \times 5 = 45$

2

3

73

- 7. (a) A mass of 8 kg gas expands within a flexible container so that p-v relationship is of the form $pv^{1\cdot 2} = constant$. The initial pressure is 1000 kPa and the initial volume is 1 m³. The final pressure is 5 kPa. If specific internal energy of the gas decreases by 40 kj/kg. find the heat transfer in magnitude and direction. 7
 - (b) A heat pump working on the Carnot cycle takes in heat from a reservoir at 5°C and delivers heat to a reservoir at 60°C. The heat pump is driven by a reversible heat engine which takes in heat from a reservoir at 840°C and rejects heat to a reservoir at 60°C. The reversible heat engine also drives a machine that absorbs 30 kW. If the heat pump extracts 17 kJ/s from the 5°C reservoir, determine (a) the rate of heat supply from the 840°C source and (b) the rate of heat rejection to the 60°C sink.
- 8. (a) A lump of steel of mass 10 kg at 627°C is dropped in 100 kg of oil at 30°C. The specific heats of steel and oil are 0.5 kJ/kg-K and 3.5 kJ/kg-K respectively. Calculate the entropy change of the steel, the oil and the universe.
 - (b) At the inlet to a certain nozzle, the enthalpy of the flowing fluid is 3000 kJ /kg and the velocity is 60 m/s. At the discharge end the enthalpy is 2762 kJ /kg. The nozzle is horizon-tal and there is negligible heat loss from it.
 - (i) Find the velocity at exit from the nozzle,
 - (ii) If the inlet area is 0.1 m^2 and the specific volume at inlet is $0.187 \text{ m}^3/\text{kg}$, find the mass flow rate,
 - (iii) If the specific volume at the nozzle exit is 0.498 m³/kg, find the exit area of the nozzle.
- 9. (a) What is a pure substance?
 - (b) What is the "critical point"? State the values of critical pressure and critical temperature of water. 3
 - (c) Why is Carnot cycle not practicable for a steam power plant?
 - (d) In an ideal air-standard diesel cycle, the pressure and temperature at intake are 1.03 bar and 27°C respectively. The maximum pressure in the cycle is 47 bar and heat supplied during the cycle is 545 kJ /kg. Determine
 - (i) The compression ratio
 - (ii) The temperature at the end of compression
 - (iii) The temperature at the end of combustion
 - (iv) The air standard efficiency.

Assume $\gamma = 1.4$ and $C_p = 1.004$ kJ/kg-K for air.

- 10. (a) How does a heat pump differ from a refrigerator?
 - (b) A household refrigerator is maintained at a temperature of 2°C. Every time the door is opened, warm material is phased inside, introducting an average of 420 kJ but making

only a small change in the temperature of the refrigerator. The door is opened 20 times a day and the refrigerator operates at 15% of the ideal COP. The cost of work is Rs. 5.00 per kWh. What is the monthly bill for this refrigerator? The atmosphere is at 30° C. 6

- (c) A blower handles 1 kg/s of air at 20°C and consumes a power of 15 kW. The inlet and outlet velocities of air are 100 m/s and 150 m/s respectively. Find the exit air temperature, assuming adiabatic conditions. Take C_p of air as 1.005 kJ/kg-K. 6
- 11. (a) Draw the rheological curve for a class of Newtonian and non-Newtonian fluids.
 - (b) Distinguish between incompressible and compressible flow.
 - (c) What is the stagnation pressure at a point in a fluid flow?
 - (d) Two pipes A and B are in the same elevation. Water 1s contained in A and rises to a level of 1.8 m above it. Carbon tetrachloride(Sp. Gr. 1.6) is contained in B. The inverted Utube is filled with compressed air at 350 kN/m² and 27°C. Barometer reads 760 mm of mercury.

Determine:

- (i) The Pressure difference between A and B if z = 0.4 m, and
- (ii) The absolute pressure in B.



12. (a) A diffuser consists of two circular parallel plates 200 mm in diameter and 5 mm apart and connected to a 30 mm diameter pipe as shown in fig. below. If the stream lines are assumed to be radial in the diffuser, what mean velocity in the pipe will correspond to an exit velocity of 0.5 m/s.



(b) A venturimeter with inlet and throat diameters are 150 mm and 75 mm respectively is mounted in a vertical pipe carrying water, the flow being upwards. The throat section is 250 mm above the inlet of the venturimeter. The discharge of the venturimeter is 40 lit/sec. Find the static pressure difference between the inlet and throat section if the coefficient of discharge for venturimeter is 0.96. 10

7

3

3

ENGINEERING & MANAGEMENT EXAMINATIONS, JUNE - 2008 MECHANICAL SCIENCE SEMESTER - 2

Time: 3 Hours

Full Marks: 70

Group–A

(Multiple	Choice	Type	Ouestions)
ſ.	manupic	Choice	-, -, -, -, -, -, -, -, -, -, -, -, -, -	Zuconono ,	,

1.	Choo	se the	correct alternatives for the foll	owing:		$10 \times 1 = 10$
	(i)	Whi	ch of the following is an intens	ive therm	nodynamic property?	
		(a)	Volume	(b)	Temperature	
		(c)	Mass	(d)	Energy.	
	(ii)	For	an irreversible process, change	in entro	py is	
		(a)	greater than dQ/T	(b)	less than dQ/T	
		(c)	zero	(d)	equal to dg/T .	
	(iii)	Duri	ing throttling which of the follo	wing qua	antity does not change?	
		(a)	Internal energy	(b)	Entropy	
		(c)	Pressure	(d)	Enthalpy.	
	(iv)	Wor	k done in a free expansion is			
		(a)	Positive	(b)	Negative	
		(c)	Zero	(d)	Maximum.	
	(v)	A cycle with constant volume heat addition and constant volume heat rejection is				
		(a)	Otto cycle	(b)	Diesel cycle	
		(c)	Joule cycle	(d)	Rankine cycle.	
	(vi)	Trip	le point of a pure substance is	a point at	t which,	
		(a)	liquid and vapour coexist	(b)	solid and vapour coexist	
		(c)	solid and liquid coexist	(d)	all three phases coexist.	
	(vii)	Berr	noulli's equation deals with the	conserva	ation of	
		(a)	Mass	(b)	Momentum	
		(c)	Energy	(d)	Work.	
((viii)	Con	tinuity equation is based on the	principle	e of conservation of	
		(a)	Mass	(b)	Momentum	
		(c)	Energy	(d)	Entropy.	
		(-)	- 00	()	-rJ-	

(ix)	A Pi	tot tube is used for measuring			
	(a)	State of fluid	(b)	Velocity of fluid	
	(c)	Density of fluid	(d)	Viscosity of fluid.	
(x)	Dyna	amic viscosity has dimensions of			
	(a)	MLT ⁻²	(b)	$ML^{-1} T^{-1}$	
	(c)	ML ⁻¹ T ⁻²	(d)	M ⁻¹ L ⁻¹ T ⁻¹	

Group–B (Short Answer Type Questions)

Answer any *three* of the following.

 $3 \times 5 = 15$

- 2. State the first law of thermodynamics for a closed system undergoing a cycle and a process.
- 3. Explain thermodynamic equilibrium.
- 4. The fluid flow is given by $\overline{V} = x^2 Y_i + Y^2 z (2xyz + yz^2)_k$. Show that this is a case of possible steady incompressible flow. Calculate the velocity and acceleration at (2, 1, 3).
- 5. Draw a block diagram of vapour compression refrigeration cycle and also show the corresponding P-V and T-S plots.
- 6. Derive Bernoulli's equation from first principles, stating the assumptions.
- 7. Explain PMM-1 and PMM-2.

Group-C

(Long Answer Type Questions)

Answer any *three* of the following. $3 \times 15 = 45$

- 8. (a) Which is a more effective way of increasing the efficiency of a Carnot engine to increase source temperatue (T_1) , keeping sink temperature (T_2) constant or to decrease T_2 keeping T_1 constant.
 - (b) State Clausius inequality.
 - (c) A mass of m kg of liquid (specific heat. = C_p) at a temperature T_1 is mixed with an equal mass of the same liquid at a temperature $T_2(T_1 > T_2)$ and the system is thermally insu-

lated. Show that the entropy change of the universe is given by $2 mC_p l\left(\frac{T_1 + T_2}{\sqrt{T_1 T_2}}\right)$ and

prove that this is necessarily positive.

- 9. (a) Derive the expression for efficiency of an Otto cycle and show the process on p-V and T-s planes.
 - (b) For the same compression ratio, explain why the efficiency of Otto cycle is greater than that of Diesel cycle.

- (c) In a diesel engine the compression ratio is 13 : 1 and fuel is cut off at 8% of the stroke. Find the air standard efficiency of the engine. Take γ for air = 1.4. 5 + 3 + 2 + 5
- 10. (a) A gas occupies 0.024 m^3 at 700 kPa and 95°C. It is expanded in the non-flow process according to the law $pv^{l\cdot 2}$ = constant to a pressure of 70 kPa after which it is heated at a constant pressure back to its original temperature.

Sketch the process on, the *p*-*V* and *T*-*s* diagrams and calculate for the whole process the work done and the heat transferred. Take $C_p = 1.047$ and $C_y = 0.775$ kJ/kg K for the gas.

- (b) A rigid closed tank of volume 3 m³ contains 5 kg of wet steam at a pressure of 200 kPa. The tank is heated until the steam becomes dry saturated. Determine the pressure and the heat transfer to the tank. 8+7
- 11. (a) Write the steady flow energy equation for a single steam entering and single steam leaving a control volume and explain the various terms.
 - (b) At the inlet to a nozzle, the enthalpy of the fluid passing is 3000 kJ/kg and velocity is 60 m/s. At the exit, the enthalpy is 2762 kJ/kg. The nozzle is horizontal and there is negligible heat loss.
 - (i) Find the velocity at the nozzle exit'
 - (ii) The inlet area is 0.1 m² and the specific volume at inlet is 0.187 m³/kg. Find the mass flow rate.
 - (iii) If the specific volume at the nozzle exit is $0.498 \text{ m}^3/\text{kg}$, find the exit area of the nozzle. 6+9
- 12. (a) Derive Euler's equation of motion along a streamline.
 - (b) A venturimeter has inlet and throat diameters of 300 mm and 150 mm. Water flows through it at the rate of 0.065m³/s and the differential gauge is deflected by 1.2 m. The specific gravity of the manometric liquid is 1.6. Determine the coefficient of discharge of the venurimeter.
 - (c) A jet of water from a 25 mm diameter nozzle is directed vertically upwards. Assuming that the jet remains circular and neglecting any loss of energy, what will be the diameter of the, jet at a point 4.5 m above the nozzle, if the jet leaves the nozzle with a velocity of 12 m/s? 5+5+5
- 13. (a) A circular disk of diameter *d* is slowly rotated in a liquid of viscosity μ at a small distance *h* from a fixed surface. Derive an expression for the torque *T* necessary to maintain an angular velocity ω .
 - (b) Distinguish between the follow:
 - (i) laminar and turbulent flow
 - (ii) compressible and incompressible fluid
 - (iii) static pressure and stagnation pressure
 - (iv) viscous and inviscid fluid.

14. Write short notes on any *three* of the following:

- (a) Pitot tube (b) Orifice meter
- (c) Point function and path function (d) Streamline streakline and pathline.

7+8 3×5

NDEX



Thermodynamics

A

Absolute pressure 17 Adiabatic expansion 110 Adiabatic index 38, 62 Adiabatic process 47, 71, 114 Air cycle 229 Air standard efficiency 230, 232

B

Boiler 116 Boundary 4 Boyle's Law, Charle's Law and Gay-Lussac's Law 59

С

Carnot cycle 139, 249 Carnot efficiency 140 Carnot theorem 142 Centigrade scale 36 Characteristic equation of gas 60 Characteristic gas constant 38, 61 Chemical equilibrium 9 Classification of properties of a system 7 Clausius inequality 162 Clausius statement 136 Clausius Theorem 158 Clearance volume 230 Closed system 4, 6 Co-efficient of performance 134 Compressibility 76 Compressibility factor 76 Compression ratio 231 Compressor 117 Condenser 117 Conduction 37 Convection 37 Constant enthalpy process 111 Constant internal energy process 110 Constant pressure cycle 235 Constant pressure line 219 Constant pressure process 9, 66, 114, 166 Constant temperature process 69, 114 Constant volume cycle 232 Constant volume gas thermometer 34 Constant volume line 219 Constant volume process 9, 64, 114, 165 Continuum 3 Control mass 6,7 Control surface 6, 7 Control volume 4, 5, 6, 7, 111

Convection 37 Critical compressibility 78 Critical isotherm 188 Critical point 78, 185, 216 Critical pressure 185 Critical temperature 78, 185 Cut-off ratio 231 Cyclic process 103 Cylinder bore 230

D

Derived Units 14 Diesel cycle 235 Diffuser 116 Displacement energy 112 Dry saturated steam 189 Dryness fraction 189 Dryness fraction lines 219

E

Efficiency 22 Efficiency of cycle 231 Electrical resistance thermometer 36 Electrical work 39,48 Enthalpy 63, 157, 189 Entropy 3, 157 Equation 113 Equation of continuity 113 Equation of state 60 Evaporator 118 Exact differentials 7 Expansion ratio 231 Extensive properties 7 External combustion engines 3

F

First law of thermodynamics 103, 105, 133 First law thermodynamics 105 Flow energy 112 Flow process 50, 64, 111 Flow work 49, 112, 113 Four stroke cycle 229 Free expansion process 48 Freezing point 36 Frictionless adiabatic process 71 Fundamental units 13

G

Gas cycle 140 Gas refrigeration cycle 251 Gauge pressure 17

Η

Heat capacity 39 Heat engine 3, 133 Heat exchanger 118 Heat pump (Reversible heat engine) 134 Heat pumps 133 Heterogeneous system 8 High temperature 33 Homogeneous system 8 Hyperbolic process 68

I

Ideal efficiency 232 Ideal gas 35, 76 Indicated thermal efficiency 232 Inexact differentials 43, 44, 45, 104 Intensive properties 7 Intensive property 44 Internal combustion engine 3 Internal energy 20 International practical temperature scale 36 Irreversibility 139, 167 Irreversible 64 Irreversible non-flow process 64, 109 Isentropic lines 216 Isentropic process 71 Isobaric process 66, 166 Isochoric process 64, 165 Isolated system 5, 105 Isothermal lines 216 Isothermal process 68, 110

Index

J

Joule's cycle 234

K

Kelvin-Planck statement 135 Kinetic energy 20

L

Latent heat of vaporisation 185 Latent heating 184 Law of conservation of energy 21, 108, 112 Law of internal energy 157 Law of the conservation of energy 107 Low temperature 33

M

Macroscopic approach 3 Manometer 18 Mass balance equation 113 Mean effective pressure 231 Mechanical equilibrium 9 Mechanical equivalent of heat 39 Mechanical work 39 Melting point 36 Microscopic approach 3 Molar gas constant 60 Molar volume 8 Mollier chart 217 Momentum 16

N

No-flow energy equation 108 Non-ideality of gases 77 Non-flow process 50, 64, 109, 113, 114, 167 Non-flow work 50 Nozzle 115

0

Open system 5, 6 Otto cycle 232

P

Paddle wheel work 45 Path function 13, 44, 105, 106 Path 9 Perfect gas 59 Perpetual motion machine of the first kind (PMM-1) 108 Perpetual motion machine of the second kind 135 Phase 8, 183 Phase equilibrium 183 **PMM-II 136** Point function 13, 44, 106 Poisson's equation 73 Polytropic index 74 Polytropic process 74, 115 Potential energy 20 Principle of increase of entropy 163 Process 9 Properties of a system 7 Property 7 Pump work 250 Pure substance 183

Q

Quasi-equilibrium 41 Quasi-equilibrium process 48 Quasi-static 41, 63 Quasi-static process 43

R

Radiation 37 Rankine cycle 249 Real gas 76 Reciprocating compressor 117 Reduced pressure 77 Reduced specific volume 77 Reduced temperature 77 Refrigerators 133, 134 Relative efficiency 232 Residual volume 77 Reversibility 138 Reversible adiabatic expansion 110 Reversible non-flow process 64, 109 Reversible process 43, 63 Rotary compressor 117

S

Saturated liquid line 185 Saturated steam line 185 Second law of thermodynamics 133 Sensible heating 184 Shaft-work 45, 48 Sink 133 Source 133 Specific air consumption 230 Specific enthalpy 63, 112, 113, 189 Specific heat 38, 61 Specific heat at constant pressure 38, 61 Specific heat at constant volume 38, 61 Specific volume 8, 17 Specific volume of ideal gas 77 Specific volumes of real gas 77 Specific weight 17 Specific work transfer 230 Standard atmospheric pressure 36 State function 106 State of a system 7 State point 7 Steady flow 111 Steady flow energy equation 116, 118 Steady flow energy equation (S.F.E.E) 111, 112, 116, 117, 118 Steady flow process 111, 113, 167 Steam tables 213 Stored energy 20, 107 Stroke length 230 Superheated steam 189 Superheated vapour 187 Supplementary units 13

Surrounding 4 Swept volume 231 system boundary 4, 6

Т

Temperature scale 35 Theoretical thermal efficiency 232 Thermal equilibrium 33 Thermal efficiency (η_{max}) 133 Thermal engineering 3 Thermal equilibrium 9,34 Thermal reservoir 133 Thermodynamic co-ordinates 7 Thermodynamic cycle 103 Thermodynamic equilibrium 8 Thermodynamic property 33, 38, 111 Thermodynamic substance 33 Thermodynamic surface 187 Thermodynamic system 105 Thermodynamic system 4 Thermodynamics 3 Thermometric property 33 Thermometric substance 34 Throttling lines 220 Throttling process 115 Total cylinder volume 231 Total energy 107 Transit energy 20 Transit energy 37 Triple point 34, 187 Triple point line 187 Triple point of water 36 Triple point temperature 34 Turbine 116 Turbine work 250 Two stroke cycle 229

U

Universal gas constant 60 Universe 4 Unresisted expansion process 110

Index

Unsteady flow 111 Unsteady flow process 111

V

Vacuum pressure 18 van der Walls equation 78 Vapour 59 Vapour compression refrigeration cycle 251 Vapour cycle 139

W

Water equivalent 39 Wet steam 189 Wheatstone bridge 36 Work ratio 230 Working substance 183

Z

Zeroth law of thermodynamics 33

Group-B Fluid Mechanics

A

Absolute pressure 294 Absolute roughness 361 Archimedes law 301 Atmospheric pressure 294

B

Barometer 295 Barometric pressure 294 Bernoulli's equation 359 Bingham plastic 267 Body forces 289, 290 Bulk modulus of elasticity 267 Buoyancy 301

С

Capillarity 270 Centre of buoyancy 302 Centre of gravity 302 Characteristic length 264 Co-efficient of compressibility 268 Coefficient of contraction 362 Coefficient of discharge 362 Coefficient of dynamic 266 Coefficient of friction 361 Coefficient of velocity 266, 362 Coefficient of viscosity 266 Compressibility 267 Compressible flow 333 Continuity equation 335 Continuum 263 Convective acceleration 339

D

Darcy-Weisbach equation 361 Datum head 360 Density 264 Differential manometer 297 Dilatant 267

E

Equipotential line 335 Euler's equations 291, 357

F

Fanning's equation 361 Filament line 335 Flow 331 Flow net 335 Flow-meters 362 Fluid dynamics 263 Fluid element 289 Fluid mechanics 263 Fluid statics 263, 289 Friction factor 361

G

Gauge pressure 294, 295

H

Head-loss 359 Hydraulic grade line 360 Hydrostatic force 298

I

Ideal fluid 267 Inclined tube manometer 296 Incompressible flow 333 Irrotational flow 333 Isentropic process 268, 366 Isothermal process 268

K

Kinetic viscosity 266 Kinetics of fluid flow 357

Index

L

Laminar flow 331 Local acceleration 339

Μ

Mach number 269 Major head loss 361 Mass density 264, 290 Mean free path 264 Metacentric height 302 Micromanometer 297 Minor losses 361

Ν

Navier-stokes equation 357 Neutral equilibrium 302 Newton's law of viscosity 265, 266 Newtonian fluid 266, 267 Non-newtonian fluid 266 Non-uniform flow 332 Normal force 289

0

Orifice 362 Orificemeter 363 Orifice-plate 363

P

Pascal's law 289, 290 Pathline 334 Piezometer 296 Piezometric head line 360 Pitot-static 367 Pitot-tube 366 Prandtl's tube 367 Pseudoplastic 267

R

Real Fluid 267

Relative roughness 361 Reynold's equation 357 Reynold's number 331, 361 Rheology 266 Rotational flow 333

S

Scalar field 329 Simple manometer 296 Specific gravity 265 Specific mass 264 Specific volume 264 Specific weight 264 Stable equilibrium 302 Stagnation density 366 Stagnation point 366 Stagnation pressure 366 Stagnation temperature 366 Static head 360 Static pressure 365 Steady 331 Streakline 335 Stream function 334 Stream surface 334 Streamline 333 Streamtube 334 Suction pressure 295 Surface forces 289, 290 Surface tension 269

T

Tangential force 289 Temperature lapse rate 294 Temporal acceleration 339 Temporal mean velocity 331 Timeline 335 Total head 360 Total head line 360 Turbulent flow 331
U Uniform flow 332 Unstable equilibrium 302 Unsteady flow 332 U-tube manometer 296

V

Vacuum pressure 295

Vector field 329 Velocity head 360 Velocity potential 335 Vena contracta 362 Venturimeter 364 Viscosity 265

W

Weight density 264

416